Heat kernel small time asymptotics at the cut locus in sub-Riemannian geometry

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Introduction

Sub-Riemannian geometry:

- generalization of Riemannian geometry under non holonomic constraints
- geometry underlying the theory of hypoelliptic operators

Main motivation:

- understand the interplay between
- \rightarrow the geometry of these spaces (optimality of geodesics, curvature)
- \rightarrow the analysis of the diffusion processes on the manifold (heat equation)

Questions:

- Which is the "right" Laplace operator associated with the geometry?
- How to relate geometric properties and the heat kernel?

Outline



2 Sub-Laplacian and Heat Equation



Outline

Sub-Riemannian geometry

2 Sub-Laplacian and Heat Equation



Sub-Riemannian manifolds

Definition

A sub-Riemannian manifold is a triple $S = (M, \mathcal{D}, \langle \cdot, \cdot \rangle)$, where

- (i) M is a connected smooth manifold of dimension $n \ge 3$;
- (ii) \mathcal{D} is a smooth distribution of (constant) rank k < n, i.e. a smooth map that associates to $q \in M$ a k-dimensional subspace \mathcal{D}_q of T_qM .

(iii) $\langle \cdot, \cdot \rangle_q$ is a Riemannian metric on \mathcal{D}_q , that is smooth as function of q.

• The set of *horizontal* vector fields on *M*, i.e.

$$\overline{\mathcal{D}} = \{X \in \textit{Vec} \ M \mid X(q) \in \mathcal{D}_q, \ orall \ q \in M\}$$
 .

satisfies the Lie bracket generating condition $\operatorname{Lie}_q \overline{\mathcal{D}} = T_q M$, $\forall q$.

Sub-Riemannian distance



• The Carnot-Caratheodory distance induced by the sub-Riemannian structure on *M* is

$$d(q,q') = \inf\{\ell(\gamma) \mid \gamma(0) = q, \gamma(T) = q', \gamma \text{ horizontal}\}.$$

 (M, d) is a metric space and d(·, ·) is finite and continuous with respect to the topology of M (Chow's Theorem)

Orthonormal frame

Locally, the pair $(\mathcal{D}, \langle \cdot, \cdot \rangle)$ can be given by assigning a set of k smooth vector fields, called a *local orthonormal frame*, spanning \mathcal{D} and that are orthonormal

$${\mathcal D}_{m{q}} = {\sf span}\{X_1(m{q}),\ldots,X_k(m{q})\}, \qquad \langle X_i(m{q}),X_j(m{q})
angle = \delta_{ij}.$$

The problem of finding geodesics, i.e. curves that minimize the length between two given points q_0 , q_1 , is equivalent to the optimal control problem

$$\begin{cases} \dot{q} = \sum_{i=1}^{k} u_i X_i(q) & \longleftarrow \text{ linear in control} \\ \int_0^T \sum_{i=1}^k u_i^2 \to \min & \longleftarrow \text{ quadratic cost} \\ q(0) = q_0, \quad q(T) = q_1 \end{cases}$$

Pontryagin Maximum Principle

Theorem (PMP)

If (u(t), q(t)) is optimal, there exists a lift $p(t) \in T^*_{q(t)}M$ and $\nu \leq 0$ such that (q(t), p(t)) is the solution of the Hamiltonian system associated with

$$\mathcal{H}^{\nu}(p,q,u) = \sum_{i=1}^{k} \langle p, X_i(q) \rangle - \frac{\nu}{2} \sum_{i=1}^{k} u_i^2$$
$$\mathcal{H}^{\nu}(p(t),q(t),u(t)) = \max_{w} \mathcal{H}^{\nu}(p(t),q(t),w)$$
(1)

- $\nu = 1$ (extr. normal): the condition (1) implies that in fact (q(t), p(t)) is a solution of $H(p, q) = \frac{1}{2} \sum_{i=1}^{k} \langle p, X_i(q) \rangle^2$
- normal extremals are smooth
- we parametrize extremals starting from q_0 with

$$\Lambda_{q_{0}} = \{p_{0} \in T_{q_{0}}^{*}M, H(p_{0}, q_{0}) = \frac{1}{2}\} \simeq S^{k-1} \times \mathbb{R}^{n-k}$$

• they are local minimizers

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• $\nu = 0$ (extr. abnormal): are abnormal minimizers smooth? open problem \rightarrow see Monti's talk

Remark: In general a trajectory can be normal and abnormal.

In what follows, we assume no abnormal minimizers that are not normal!

Basic features in SRG - 1

The cut and the conjugate locus starting from a point can be defined analogously to the Riemannian case



- Conjugate Locus: set of point where geodesics lose local optimality
- Cut locus: set of points where geodesics lose global optimality (and the distance is not smooth)
- Ball of radius R: points reached optimally in time $T \leq R$

Basic features in SRG - 2

Consider geodesics from a point $x_0 \in M$

- there are geodesics loosing optimality arbitrarily close to x₀
- $x \mapsto d^2(x_0, x)$ is not smooth in x_0



Theorem (Agrachev, after Trelat-Rifford)

Let $x_0 \in M$ and $E(x) = \frac{1}{2}d^2(x_0, x)$. Then

• E is smooth on the open dense set $\Sigma \subset M$

 $\Sigma(x_0) = \{x \in M \mid \exists! \text{ strictly normal non-conjugate minimizer from } x_0 \text{ to } x\}$

Hausdorff dimension

Define
$$\mathcal{D}^1 := \mathcal{D}$$
, $\mathcal{D}^2 := \mathcal{D} + [\mathcal{D}, \mathcal{D}]$, $\mathcal{D}^{i+1} := \mathcal{D}^i + [\mathcal{D}^i, \mathcal{D}]$.

- Hörmander cond. $\rightarrow \exists m \in \mathbb{N} (step) \text{ s.t. } \mathcal{D}_q^m = T_q M$
- growth vector of the structure is the sequence

$$\mathcal{G}(S) := (\dim_{k} \mathcal{D}, \dim \mathcal{D}^{2}, \dots, \dim_{n} \mathcal{D}^{m})$$

• Mitchell Theorem: the Hausdorff dimension of (M, d) is given by the formula

$$Q=\sum_{i=1}^m ik_i=k_1+2k_2+\ldots+mk_m,$$

The Hausdorff dimension is bigger than the topological one: Q > n

Sub-Riemannian geometry

2 Sub-Laplacian and Heat Equation

3 Examples

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Sub-Laplacian

The sub-Laplacian operator Δ on $(M, \langle \cdot, \cdot \rangle, D)$ is the natural generalization of the Laplace-Beltrami operator and is defined as follows

 $\Delta \phi = div(grad\phi)$

• grad ϕ is the unique *horizontal* vector field dual to $d\phi|_{\mathcal{D}}$ (use the *metric*)

$$grad\phi = \sum_{i=1}^{k} X_i(\phi) X_i$$

- The divergence of a vector field X says how much the flow of X change a volume μ, i.e. L_Xμ = (div X)μ
- The Leibnitz rule for *div* gives

$$\Delta = \sum_{i=1}^k X_i^2 + (\operatorname{div} X_i) X_i$$

 \rightarrow sum of squares + first order terms that depend on the volume

We assume that a smooth volume μ is fixed!

Which volume?

The choice of a canonical volume is related with the definition of a canonical sub-Laplacian.

- in the Riemannian case the metric canonically defines the Riemannian volume
- being a metric space, one can consider the (spherical) Hausdorff volume
- in this case they are proportional

In the sub-Riemannian case the metric does not define a canonical volume in principle

- in the equiregular case we can use the structure of the Lie bracket to define a volume, called Popp volume
- the spherical Hausdorff volume is not proportional and in general, it is not even smooth
- \rightarrow See Gauthier's talk.

Heat equation

Sub-Riemannian heat equation on a *complete* manifold M with smooth measure μ

$$\begin{cases} \frac{\partial \psi}{\partial t}(t,x) = \Delta \psi(t,x), & \text{ in } (0,\infty) \times M, \\ \psi(0,x) = \varphi(x), & x \in M, \quad \varphi \in C_0^\infty(M). \end{cases}$$
(*)

where $\psi(0, x) = \lim_{t \to 0} \psi(t, x)$. Recall that

$$\Delta = \sum_{i=1}^k X_i^2 + (\operatorname{div} X_i) X_i$$

Theorem (Hörmander)

If X_1, \ldots, X_k are bracket generating then $\Delta = \sum_{i=1}^k X_i^2 + X_0$ is hypoelliptic.

Heat kernel

For $f,g \in C_0^\infty(M)$ with compact support, the standard divergence theorem

$$\int_{\mathcal{M}} f(\operatorname{div} X) \, d\mu = - \int_{\mathcal{M}} X f \, d\mu, \qquad \text{(no metric required!)}$$

applied to X = grad g gives

$$\int_{M} f \, \Delta g \, d\mu = - \int_{M} \langle \operatorname{grad} f, \operatorname{grad} g \rangle \, d\mu$$

implies that Δ is symmetric and negative $\rightarrow \Delta$ essentially self-adjoint on $C_0^{\infty}(M)$. The problem (*) has a unique solution for every initial datum $\varphi \in C_0^{\infty}(M)$

$$\psi(t,x) := e^{t\Delta} \varphi(x) = \int_M p_t(x,y) \varphi(y) d\mu(y), \qquad \varphi \in C_0^\infty(M),$$

where $p_t(x, y) \in C^{\infty}$ is the so-called *heat kernel* associated to Δ .

Fix $x, y \in M$, dim M = n:

Theorem (Main term, Leandre)	
$\lim_{t\to 0} 4t \log p_t(x,y) = -d^2(x,y)$	(1)

Theorem (Smooth points, Ben Arous)

Assume $y \in \Sigma(x)$, then

$$p_t(x,y) \sim \frac{1}{t^{n/2}} \exp\left(-\frac{d^2(x,y)}{4t}\right)$$

- 1. In Riemannian geometry $x \in \Sigma(x)$, in sub-Riemannian is not true!
- 2. The on-the-diagonal expansion indeed is different.

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Theorem (On the diagonal, Ben Arous)	

We have the expansion

$$p_t(x, \mathbf{x}) \sim \frac{1}{t^{\mathbf{Q}/2}}$$

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Theorem (Smooth points, Ben Arous)

Assume $y \in \Sigma(x)$, then

$$p_t(x,y) \sim \frac{1}{t^{n/2}} \exp\left(-\frac{d^2(x,y)}{4t}\right)$$
 (2)

Questions

- 1. What happens in (2) if $y \in Cut(x)$?
- 2. Can we relate the expansion of $p_t(x, y)$ with the properties of the geodesics joining x to y?

Example: Heisenberg

In the Heisenberg group the Heat kernel is explicit (here q = (x, y, z))

$$p_t(0,q) = \frac{1}{(4\pi t)^2} \int_{-\infty}^{\infty} \frac{\tau}{\sinh \tau} \exp\left(-\frac{x^2 + y^2}{4t} \frac{\tau}{\tanh \tau}\right) \cos\left(\frac{z\tau}{t}\right) d\tau.$$

and gives the asymptotics for cut-conjugate points $\zeta = (0, 0, z)$

$$p_t(0,\zeta) \sim \frac{1}{t^2} \exp\left(-\frac{\pi z}{t}\right) = \frac{1}{t^2} \exp\left(-\frac{d^2(0,\zeta)}{4t}\right)$$

Remark: The factor $\frac{1}{2}$ that is added confirm the fact that the points $\zeta = (0, 0, z)$ are not smooth points. What is the meaning of the $\frac{1}{2}$?

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$$p_t(0,\zeta) \sim \frac{1}{t^{\frac{3}{2}+\frac{1}{2}}} \exp\left(-\frac{\pi z}{t}\right) = \frac{1}{t^{\frac{3}{2}+\frac{1}{2}}} \exp\left(-\frac{d^2(0,\zeta)}{4t}\right)$$

Remark: The factor $\frac{1}{2}$ that is added confirm the fact that the points $\zeta = (0, 0, z)$ are not smooth points. What is the meaning of the $\frac{1}{2}$?

What happens at non good point?

Let $x, y \in M$ with $y \in Cut(x)$ and write

$$p_t(x,y) = \int_M p_{t/2}(x,z) p_{t/2}(z,y) d\mu(z)$$

Idea: assume $z \in \Sigma(x) \cap \Sigma(y)$ and apply Ben-Arous expansion

$$p_{t/2}(x,z)p_{t/2}(z,y) \sim \frac{1}{t^n} \exp\left(-\frac{d^2(x,z)+d^2(z,y)}{4t}\right)$$

This led to the study of an integral of the kind

$$p_t(x,y) = \frac{1}{t^n} \int_M c_{x,y}(z) \exp\left(-\frac{h_{x,y}(z)}{2t}\right) d\mu(z)$$

where $h_{x,y}$ is the hinged energy function

$$h_{x,y}(z) = \frac{1}{2} \left(d^2(x,z) + d^2(z,y) \right).$$

Properties of $h_{x,y}$ hinged energy function

Lemma

Let Γ be the set of midpoints of the minimal geodesics joining x to y. Then min $h_{x,y} = h_{x,y}(\Gamma) = d^2(x,y)/4$.

• A minimizer is called strongly normal if any piece of it is strictly normal.

Theorem

Let γ be a strongly normal minimizer joining x and y. Let z_0 be its midpoint. Then

- (i) y is conjugate to x along $\gamma \Leftrightarrow \text{Hess}_{z_0}h_{x,y}$ is degenerate.
- (ii) The dimension of the space of perturbations for which γ is conjugate is equal to dim ker $\text{Hess}_{z_0}h_{x,y}$.

Remark: Hess $h_{x,y}$ is never degenerate along the direction of the geodesic!

Laplace integrals in \mathbb{R}^n

Let $K \subset \mathbb{R}^n$ be a compact set and assume g has a minimum in int K

• The asymptotic of the Laplace integral

$$\int_{\mathcal{K}} f(z) e^{-g(z)/t} dx \quad \text{as } t \searrow 0$$

is determined by the behavior of the function g near its minimum.

• Example in \mathbb{R} with $g(z) = z^{2m}$

$$I(t) = \int_{a}^{b} f(z) e^{-z^{2m}/t} dz \sim C_m f(0) t^{1/2m}, \quad t \to 0$$

• Example in \mathbb{R}^n with $g(z) = \sum z_i^{2m_i}$

$$I(t) = \int_{a}^{b} f(z) e^{-g(z)/t} dz \sim C f(0) t^{\sum_{i} 1/2m_{i}}, \quad t \to 0$$

Hinged vs Asymptotics

Theorem

For a sufficiently small neighborhood $N(\Gamma)$ of the set of midpoints from x to y

$$p_t(x,y) = \frac{1}{t^n} \int_{N(\Gamma)} \exp\left(-\frac{h_{x,y}(z)}{2t}\right) (c_{x,y}(z) + O(t)) d\mu(z)$$

Assume that, in a neighborhood of the midpoints of the strongly normal geodesics joining x to y there exists coordinates such that

$$h_{x,y}(z) = rac{1}{4}d^2(x,y) + z_1^{2m_1} + \ldots + z_n^{2m_n} + o(|z_1|^{2m_1} + \ldots + |z_n|^{2m_n})$$

Then for some constant C > 0

$$p_t(x,y) = \frac{1}{t^{n-\sum_i \frac{1}{2m_i}}} \exp\left(-\frac{d^2(x,y)}{4t}\right) (C+o(1)).$$

Note: $h_{x,y}$ non degenerate $(m_i = 2) \rightarrow$ the exponent is n/2

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Example: Heisenberg

In the Heisenberg group we had the asymptotics for cut-conjugate points $\zeta = (0,0,z)$

$$p_t(0,\zeta) \sim \frac{1}{t^2} \exp\left(-\frac{\pi z}{t}\right) = \frac{1}{t^2} \exp\left(-\frac{d^2(0,\zeta)}{4t}\right)$$

Remark: The factor $\frac{1}{2}$ is a consequence of the fact that there exists a one parametric family of optimal trajectories (varying the angle), hence the hinged energy function is actually a function of two variables, being constant on the midpoints.

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Cut/Conjugacy vs Asymptotics

Using the geometric properties of $h_{x,y}$ we can state also the following

Corollary

Let M be an n-dimensional complete SR manifold, μ smooth volume. Let $x \neq y$ and assume that every optimal geodesic joining x to y is strongly normal.

• Then there exist positive constants C_i, such that for small t

$$\frac{C_1}{t^{n/2}}e^{-d^2(x,y)/4t} \le p_t(x,y) \le \frac{C_2}{t^{n-(1/2)}}e^{-d^2(x,y)/4t},$$

If x and y are conjugate along at least one minimal geodesic

$$\frac{C_3}{t^{(n/2)+(1/4)}}e^{-d^2(x,y)/4t} \le p_t(x,y)$$

• If x and y are not conjugate

$$p_t(x,y) = \frac{C_4 + O(t)}{t^{n/2}} e^{-d^2(x,y)/4t},$$

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• If x and y are not conjugate

$$p_t(x,y) = \frac{C_4 + O(t)}{t^{n/2}} e^{-d^2(x,y)/4t},$$

Remarks

- The precise correspondence between *Hess* $h_{x,y}$ and cut/conjugacy works at order 1.
- To have the precise asymptotic one need that the expansion of $h_{x,y}$ is diagonal in some coordinates.

Nevertheless there are at least two cases that simplifies the analysis

- If we have a one parametric family of optimal trajectories then $h_{x,y}$ is constant along the trajectory of midpoints.
- If there is only one degenerate direction then $h_{x,y}$ is always diagonalizable

Lemma (Splitting Lemma - Gromoll, Meyer)

Let $g : \mathbb{R}^n \to \mathbb{R}$ smooth such that g(0) = dg(0) = 0 and that dim ker $d^2g(0) = 1$ with 0 isolated minimum of g.

Then there exists a diffeomorphism ϕ of \mathbb{R}^n and a smooth function $\psi : \mathbb{R} \to \mathbb{R}$ such that $_{n-1}$

$$g(\phi(u)) = \sum_{i=1}^{n-1} u_i^2 + \psi(u_n), \quad \text{where} \quad \psi(u_n) = O(u_n^4).$$

Outline

Sub-Riemannian geometry

2 Sub-Laplacian and Heat Equation



Bi-Heisenberg group

Is the sub-Riemannian structure on \mathbb{R}^5 defined by the o.n. frame

$$\begin{split} X_1 &= \partial_{x_1} - \frac{\alpha_1}{2} y_1 \partial_z, \qquad X_2 &= \partial_{y_1} + \frac{\alpha_1}{2} x_1 \partial_z \\ X_3 &= \partial_{x_2} - \frac{\alpha_2}{2} y_2 \partial_z, \qquad X_4 &= \partial_{y_2} + \frac{\alpha_2}{2} x_2 \partial_z, \end{split}$$

Remark: If $Z = \partial_z$ then $[X_1, X_2] = \alpha_1 Z$ and $[X_3, X_4] = \alpha_2 Z$,

- we restrict to the contact case $\alpha_1, \alpha_2 > 0$.
- one can always assume $0 < \alpha_1 \le \alpha_2$
- the case $\alpha_1 < \alpha_2$ and $\alpha_1 = \alpha_2$ are really different!!!

Exponential map

The exponential map starting from the origin $\operatorname{Exp}_0 : \Lambda_0 \times \mathbb{R}^+ \to M$ where

$$\Lambda_0 = \{ p_0 = (r_1, r_2, \theta_1, \theta_2, w) \in T_0^*M \mid r_1^2 + r_2^2 = 1, \ \theta_1, \theta_2 \in S^1, \ w \in \mathbb{R} \}.$$

The arclength geodesics associated with $p_0 = (r_1, r_2, \theta_1, \theta_2, w) \in \Lambda_0$ and $|w| \neq 0$

$$\begin{aligned} x_i(t) &= \frac{r_i}{\alpha_i w} (\cos(\alpha_i wt + \theta_i) - \cos \theta_i), \\ y_i(t) &= \frac{r_i}{\alpha_i w} (\sin(\alpha_i wt + \theta_i) - \sin \theta_i), \qquad i = 1, 2, \\ z(t) &= \frac{1}{2w^2} (wt - \sum_i \frac{r_i^2}{\alpha_i} \sin \alpha_i wt). \end{aligned}$$
(3)

Lemma

Each of these geodesic is optimal up to its cut time

$$t_{cut} = t_{conj} = \frac{2\pi}{|w| \max\{\alpha_1, \alpha_2\}}$$

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(4)

Bi-Heisenberg group

Cut locus

- $\alpha < 1 \rightarrow$ the cut locus is a three dimensional set containing the z-axis.
 - Each point is reached by a one parametric family of optimal geodesics
- $\alpha = 1 \rightarrow$ the cut locus degenerates to the z-axis only.
- Each point is reached by a three parametric family of optimal geodesics Heat Kernel at a point $\zeta = (0, 0, 0, 0, z)$ on the z-axis
 - $\alpha < 1 \rightarrow$ the expansion

$$p_t(0,\zeta) \sim \frac{1}{t^3} \exp\left(-\frac{d^2(0,\zeta)}{4t}\right)$$

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Application: Grushin plane

 $\bullet\,$ It is the generalized sub-Riemannian structure on \mathbb{R}^2 for which an orthonormal frame of vector fields is given by

$$X = \partial_x, \qquad Y = x \partial_y. \tag{5}$$

• we want to compute the expansion of the heat kernel in the Grushin plane at a cut-conjugate point, starting from a Riemannian point.

$$\Delta = X^2 + Y^2 = \partial_x^2 + x^2 \partial_y^2$$

Remarks:

- the exponential map from the origin has no cut-conjugate points
- from a Riem point the exponential map is integrable in trigonometric functions and provides a generic structure of cut and conjugate locus.
- the canonical (Riemannian) volume is not smooth (actually the heat does not cross the vertical line in this case [Boscain, Laurent])

• The sub-Riemannian Hamiltonian in $\mathcal{T}^*\mathbb{R}^2$ is the smooth function

$$H: T^* \mathbb{R}^2 \to \mathbb{R}, \qquad H(p_x, p_y, x, y) = \frac{1}{2}(p_x^2 + x^2 p_y^2).$$
 (6)

- There are no abormal minimizers
- The arclength geodesic flow $\mathsf{Exp}:\mathbb{R}_+\times S^1\to\mathbb{R}^2$
 - ightarrow starting from the Riemannian point $q_0=(-1,-\pi/4)$
 - \rightarrow initial condition $(p_x(0), p_y(0)) = (\cos \theta, \sin \theta)$

$$\begin{aligned} \mathsf{Exp}(t,\theta) &= (x(t,\theta), y(t,\theta)), \\ x(t,\theta) &= -\frac{\sin(\theta - t\sin\theta)}{\sin\theta}, \\ y(t,\theta) &= -\frac{\pi}{4} + \frac{1}{4\sin\theta} \left(2t - 2\cos\theta + \frac{\sin(2\theta - 2t\sin\theta)}{\sin\theta} \right), \end{aligned} \tag{7}$$

• the point $q_1 = (1, \pi/4)$ is a cut-conjugate point





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The hinged energy function

$$\begin{split} h(x,y) &= \frac{1}{2} \left(d^2(q_0,(x,y)) + d^2(q_1,(x,y)) \right) \\ &= \frac{1}{2} \left(d^2(q_0,(x,y)) + d^2(q_0,(-x,-y)) \right), \end{split}$$

We found an (explicit!) change of coordinates $(x, y) = \phi(u, v)$ near the origin such that

$$h(\phi(u, v)) = u^{2} + v^{4} + O(||(u, v)||^{5})$$

Using that $d(q_0,q_1)=\pi$, this proves

Theorem

The heat kernel $p_t(q_0, q_1)$ satisfies the following asymptotic expansion

$$p_t(q_0, q_1) = \frac{1}{t^{5/4}} exp\left(-\frac{\pi^2}{4t}\right) (C + O(t))$$
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Theorem

The heat kernel $p_t(q_0, q_1)$ satisfies the following asymptotic expansion

$$p_t(q_0, q_1) = \frac{1}{t^{2-\frac{1}{2}-\frac{1}{4}}} exp\left(-\frac{\pi^2}{4t}\right) (C + O(t))$$
(8)

Open problems and related works

- \rightarrow Expansion at the cut/conjugate locus
 - Case of a generic 2D Riemannian and 3D sub-Riemannian
 - in the case of a 3D sub-Riemannian the cut locus is adjacent to the origin
- \rightarrow Heat kernel and geometry
 - we obtained an expansion on the diagonal in the 3D

$$p(t,x,x)\sim rac{1}{16t^2}(1+\kappa(x)t+O(t^2)), \qquad ext{for} \quad t
ightarrow 0.$$

- higher dim? Curvature of general sub-Riemannian spaces
 → see Agrachev's talk
- → References:
 - D.B., U. Boscain, R. Neel, *Small time asymptotics of the SR heat kernel at the cut locus*, accepted on JDG, on ArXiv
 - D.B., U. Boscain, R. *The Bi-Heisenberg group in SR geometry*, in preparation.