

# Preconditioning for electromagnetic scattering via localization

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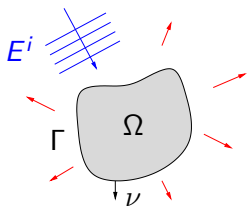


# Summary

- ▶ Time-harmonic Maxwell equations
- ▶ General ideas of Boundary Integral Methods
- ▶ A well-conditioned scheme for electromagnetic scattering
- ▶ Numerical results and future directions

# Electromagnetic scattering

Scattering of time-harmonic waves  $E^i(x)e^{-\omega t}$



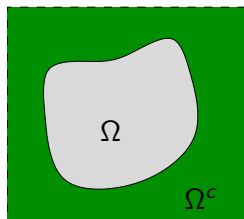
$$\nabla \times \nabla E - k^2 E = 0 \text{ in } \Omega^c$$

$$\gamma_D E = -\nu \times E_0 \text{ on } \Gamma$$

$$\lim_{|x| \rightarrow \infty} ((\nabla \times E) \times x - ikE|x|) = 0$$

“Dirichlet” trace  $\gamma_D = \nu \times$

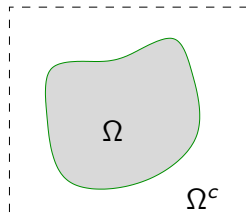
# Boundary integral method for PDE-problems



Direct discretization of

$$\begin{aligned}Lu &= 0 \text{ in } \Omega^c \\ \gamma_D u &= u_0 \text{ on } \Gamma\end{aligned}$$

$\rightsquigarrow$  **sparse** matrix equation for  
 $n$ -dim domain (unbounded!)



Parametrization  $u = V(g)$  with

$$\begin{aligned}LV &= 0 \text{ in } \Omega^c \\ \gamma_D V(g) &= u_0 \text{ on } \Gamma\end{aligned}$$

$\rightsquigarrow$  **dense** matrix equation for  
 $n - 1$ -dim boundary

## Natural choices for the potential

$$\left. \begin{array}{l} \text{Single-layer potential } \mathcal{T} = \frac{1}{ik} \nabla \times (\nabla(\mathcal{G}_k u)) \\ \text{Double-layer potential } \mathcal{K} = \nabla \times (\mathcal{G}_k) \end{array} \right\}$$

where  $\mathcal{G}_k u = \int g_k(x-y)u(y)dy$  and  $g_k = -\frac{1}{4\pi} \frac{e^{ik|x|}}{|x|}$

Suitable coupling

$$E = (\mathcal{T} R_1 - \mathcal{K} R_2)g$$

Classical choices for the coupling

	$R_1$	$R_2$
EFIE	1	0
MFIE	0	1
CFIE	$\lambda$	$1 - \lambda$

lead to **ill-conditioned** and **dense** systems

$\rightsquigarrow$  complexity  $\sim \mathcal{O}(n \times N^2)$

# An intrinsically well-conditioned formulation

... using representation formula

$$E = \mathcal{T}\gamma_N E - \mathcal{K}\gamma_D E$$

... and Dirichlet-to-Neumann operator  $Y : \gamma_D E \rightarrow \gamma_N E$

$$E = \mathcal{T}Y\gamma_D E - \mathcal{K}\gamma_D E$$

... boundary-integral operator reduces to identity

$$\gamma_D E = \underbrace{\gamma_D(\mathcal{T}Y - \mathcal{K})}_{=\text{id}} \gamma_D E$$

$Y$  **not** explicitly known  $\rightsquigarrow$  construct suitable approximations  $\tilde{Y}$

## Crucial observation [Nédélec]

Define

$$\begin{aligned}T &:= \gamma_D \mathcal{T} \\ K - \frac{1}{2} &:= \gamma_D \mathcal{K}\end{aligned}$$

Calderon formula

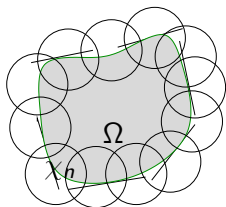
$$\underbrace{T^2}_{\text{order 1}} = \underbrace{K^2}_{\text{order -1}} - \frac{1}{4}$$

= compact perturbation of identity

Formal expansion

$$\begin{aligned}Y &= -2T(1 - 2K)^{-1} \\ &\approx -2T + h.o.t.\end{aligned}$$

# Localization for high frequencies [Borel]



Dirichlet-to-Neumann  
operator on a plane

$$Y = -2T$$

Local approximation using  
partition of unity  $\sum \chi_n^2 = 1$

$$\tilde{Y} = \sum_n \chi_n Y_n \chi_n$$

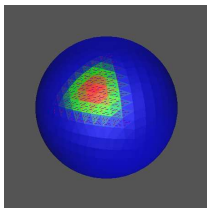


# Library at Onera

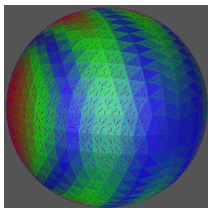
So far: Implementation of

- ▶ Classical methods (EFIE,MFIE)
- ▶ Helmholtz decomposition (smooth surface)
- ▶ Fast Multipole Method (FMM)

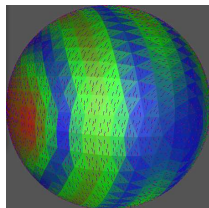
# Results using localization



$\chi_n$



Incident wave



Solution

Observe

diam supp  $\chi_n$  and  $\lambda$  are comparable

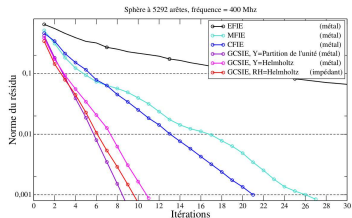


FIGURE 2.5 – Résultats numériques pour une fréquence de 400 MHz, sur une sphère maillée avec 5292 arêtes.

## Future directions

- ▶ Geometric singularities
- ▶ Consider higher order expansion of  $Y = -2T(1 - 2K)^{-1}$
- ▶ Open surfaces