The influence of the wall on the motion of micro-swimmer

Laetitia Giraldi¹ Advisor François Alouges

¹CMAP Ecole Polytechnique

Journée de bilan de la chaire MMSN





• Developing self-propulsion at micro-scales?

• Application in human diagnostics and therapy...



< < >> < </p>

The swimmer is described by the vector (ξ, p) such as :

- ξ is a function which defines the shape of the swimmer.
- $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\implies \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

$$\begin{bmatrix} -\nu \Delta u + \nabla p = f, \\ \operatorname{div} u = 0. \end{bmatrix}$$

Self-propulsion constraints $\implies \begin{cases} \sum Forces = 0 \\ Torque = 0 \end{cases}$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0 \\ \int_{\partial\Omega} x_{0} \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0. \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p} = V(p,\xi)\dot{\xi}$.

イロト 不得 とくほと くほとう

The swimmer is described by the vector (ξ, p) such as :

• ξ is a function which defines the shape of the swimmer.

• $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\Longrightarrow \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

> $-\nu\Delta u + \nabla p = f,$ div u = 0.

Self-propulsion constraints $\implies \begin{cases} \sum Forces = 0 \\ Torque = 0 \end{cases}$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0 \\ \int_{\partial\Omega} x_{0} \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0. \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p}=V(p,\xi)\dot{\xi}.$

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ▶ ◆ □ ● ○ ○ ○

The swimmer is described by the vector (ξ, p) such as :

- ξ is a function which defines the shape of the swimmer.
- $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\implies \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

$$\begin{aligned} -\nu \Delta u + \nabla p &= f, \\ \operatorname{div} u &= 0. \end{aligned}$$

Self-propulsion constraints $\implies \begin{cases} \sum Forces = 0 \\ Torque = 0 \end{cases}$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0 \\ \int_{\partial\Omega} x_{0} \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0. \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p}=V(p,\xi)\dot{\xi}.$

イロン 不得 とくほ とくほ とうほ

The swimmer is described by the vector (ξ, p) such as :

- ξ is a function which defines the shape of the swimmer.
- $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\implies \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

 $\begin{aligned} -\nu \Delta u + \nabla p &= f, \\ \operatorname{div} u &= 0. \end{aligned}$

Self-propulsion constraints $\implies \begin{cases} \sum Forces = 0 \\ Torque = 0 \end{cases}$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_0 = 0 \\ \int_{\partial\Omega} x_0 \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_0 = 0. \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p}=V(p,\xi)\dot{\xi}.$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

The swimmer is described by the vector (ξ, p) such as :

- ξ is a function which defines the shape of the swimmer.
- $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\Longrightarrow \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

 $\begin{bmatrix} -\nu\Delta u + \nabla p = f, \\ \operatorname{div} u = 0. \end{bmatrix}$

Self-propulsion constraints $\implies \begin{cases} \sum Forces = 0 \\ Torque = 0 \end{cases}$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_0 = 0 \\ \int_{\partial\Omega} x_0 \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_0 = 0 \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p}=V(p,\xi)\dot{\xi}.$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

The swimmer is described by the vector (ξ, p) such as :

- ξ is a function which defines the shape of the swimmer.
- $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\Longrightarrow \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

$$\begin{bmatrix} -\nu\Delta u + \nabla p = f, \\ \operatorname{div} u = 0. \end{bmatrix}$$

 $\label{eq:Self-propulsion constraints} \Longrightarrow \left\{ \begin{array}{cc} \sum Forces & = 0 \\ Torque & = 0 \end{array} \right.$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0 \\ \int_{\partial\Omega} x_{0} \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{p} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0. \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p} = V(p,\xi)\dot{\xi}$.

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

The swimmer is described by the vector (ξ, p) such as :

- ξ is a function which defines the shape of the swimmer.
- $p = (c, R) \in \mathbb{R}^3 \times SO(3)$ parametrizes the swimmer's position.

The swimmer changes its shape $\implies \xi(t)$ pushes the fluid. The fluid reacts, under the Stokes Equation

$$\begin{bmatrix} -\nu\Delta u + \nabla p = f, \\ \operatorname{div} u = 0. \end{bmatrix}$$

 $\label{eq:Self-propulsion constraints} \Longrightarrow \left\{ \begin{array}{cc} \sum Forces & = 0 \\ Torque & = 0 \end{array} \right.$

$$\iff \begin{cases} \int_{\partial\Omega} DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{\rho} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0 \\ \int_{\partial\Omega} x_{0} \times DN_{\rho,\xi} \left((\partial_{\rho} \Phi) \dot{\rho} + (\partial_{\xi} \Phi) \dot{\xi} \right) dx_{0} = 0. \end{cases}$$

As a result the swimmer moves, under the ODE

 $\dot{p} = V(p,\xi)\dot{\xi}.$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

Controllability issues

$$\left\{ \begin{array}{l} \dot{p} = V(p,\xi) \dot{\xi} \\ p_0 \end{array} \right.$$

Questions

- Is it possible to control the state of the system (ξ and p) using as controls only the rate of shape changes ^d/_{dt}ξ?
- Does the boundary have an effect on the controllability of the swimmer?



The swimmer that we consider consists of *n* spheres connected by the swimmer's arm.

The change of the swimmer's shape consists in changing the length of its arms $(\xi_i)_i$.



Controllability's result in \mathbb{R}^3 [Alouges, DeSimone, Heltai, Lefevbre, Merlet (Preprint)]



 Does the presence of a wall modify the swimmer's reachable set?

Influence of the wall - Main results [Alouges, G]



The 4-spheres swimmer is globally controllable on an dense open set.



- For any initial condition (y_0, θ_0) such that $\theta_0 \neq \frac{\pi}{2}$, the swimmer can reach every (y_G, θ_G) given $(\theta_G \neq \frac{\pi}{2})$.
- If θ₀ = π/2 then the swimmer cannot change its angle and it moves only on a straight line defined by itself. (i.e., the dimension of

 $Lie_{(\xi_1,\xi_2,y,\frac{\pi}{2})}(V_1,V_2)$ is equal to 3.

Outline of the proof

$$\dot{p} = \sum_{i=1}^{M} V_i(p,\xi) \dot{\xi}$$

By studying the dimension of the subspace $Lie_{(p,\xi)}((V_i)_{i=1..M})$ which denotes the set of all tangent vectors $V(p,\xi)$ in $Lie((V_i)_{i=1..M})$.

- By using the limit and the case without wall
- By calculation of Lie Brackets and application of Nagano (1966) Hermann (1963) theorem [Lobry 1970], we show that there are two kinds of orbit :
 - the orbit with a 3 dimensional Lie space (if $\theta_0 = \frac{\pi}{2}$).
 - the others such that the dimension is equal to 4.

Conclusion and outlook

• Influence of the boundary.



Optimal strokes.



イロト イヨト イヨト イ

≣ ▶