

**ECOLE POLYTECHNIQUE**  
**Applied Mathematics Master Program**  
**MAP 562 Optimal Design of Structures (G. Allaire)**  
**Answers to the written exam of March 4th, 2015.**

## 1 Parametric optimization: 14 points

1. Writing  $v = \langle u'(h), k \rangle$ ,  $\Lambda = \langle \lambda'(h), k \rangle$  and differentiating problem (1) yields

$$\begin{cases} -\operatorname{div}(h\nabla v) - \operatorname{div}(k\nabla u) = \lambda\rho hv + \lambda\rho ku + \Lambda\rho hu & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega. \end{cases}$$

Differentiating the normalization condition  $\int_{\Omega} \rho hu^2 dx = 1$  we obtain

$$\int_{\Omega} \rho u(2hv + ku) dx = 0.$$

2. Multiplying the above problem for  $v$  by  $u$  and integrating by parts leads to

$$\int_{\Omega} (h\nabla v \cdot \nabla u + k\nabla u \cdot \nabla u) dx = \int_{\Omega} (\lambda\rho hv u + \lambda\rho ku^2 + \Lambda\rho hu^2) dx.$$

From the equation for  $u$  we know that

$$\int_{\Omega} h\nabla v \cdot \nabla u dx = \int_{\Omega} \lambda\rho hv u dx.$$

Therefore, from the normalization  $\int_{\Omega} \rho hu^2 dx = 1$ , we deduce

$$\Lambda = \int_{\Omega} k (|\nabla u|^2 - \lambda\rho u^2) dx.$$

In other words, we found  $\lambda'(h) = |\nabla u|^2 - \lambda\rho u^2$  which is a function in  $L^1(\Omega)$ .

3. From the previous question and since, by assumption,  $\nabla u \neq 0$  on  $\partial\Omega$  and  $u = 0$  on  $\partial\Omega$ , we find  $\lambda'(h) > 0$  on  $\partial\Omega$  (and thus by continuity on a neighborhood of  $\partial\Omega$ ). Therefore, for an optimal thickness  $h$ , the lower constraint  $h \geq h_{min}$  must be saturated in this region, i.e.,  $h = h_{min}$  near  $\partial\Omega$ . On the contrary, at the point where  $u$  is maximum, we have  $\nabla u = 0$  and  $u > 0$ , thus  $\lambda'(h) < 0$ . It implies that, in the vicinity of the maximum of  $u$ , we must have  $h = h_{max}$ .

4. The objective function is

$$J(h) = \int_{\Omega} j \left( \frac{u(h)}{\|u(h)\|} \right) dx.$$

It is clear that, since  $j$  is even, i.e.,  $j(-w) = j(w)$ , we have for any  $t \neq 0$

$$j \left( \frac{tu(h)}{\|tu(h)\|} \right) = j \left( \frac{u(h)}{\|u(h)\|} \right),$$

so that the objective function is independent from the normalization of the eigenfunction.

For  $h \in \mathcal{U}_{ad}$ ,  $\hat{\lambda} \in \mathbb{R}$ ,  $\hat{u} \in H_0^1(\Omega)$  and  $\hat{p} \in H_0^1(\Omega)$  we define the Lagrangian

$$\mathcal{L}(h, \hat{\lambda}, \hat{u}, \hat{p}) = \int_{\Omega} j \left( \frac{\hat{u}}{\|\hat{u}\|} \right) dx + \int_{\Omega} (h \nabla \hat{u} \cdot \nabla \hat{p} - \hat{\lambda} \rho h \hat{u} \hat{p}) dx.$$

5. For  $\hat{u} \in L^2(\Omega)$  we define

$$F(\hat{u}) = \int_{\Omega} j \left( \frac{\hat{u}}{\|\hat{u}\|} \right) dx.$$

By the chain rule lemma, since the derivative of  $\|\hat{u}\|$  in the direction of  $\phi$  is  $\langle \hat{u}, \phi \rangle / \|\hat{u}\|$  with the notation  $\langle \hat{u}, \phi \rangle = \int_{\Omega} \hat{u} \phi dx$ , we obtain

$$\langle F'(\hat{u}), \phi \rangle = \int_{\Omega} F'(\hat{u}) \phi dx = \int_{\Omega} j' \left( \frac{\hat{u}}{\|\hat{u}\|} \right) \left( \frac{\phi}{\|\hat{u}\|} - \frac{\langle \hat{u}, \phi \rangle}{\|\hat{u}\|^3} \hat{u} \right) dx.$$

Clearly we find  $\langle F'(\hat{u}), \hat{u} \rangle = 0$ . Equivalently,

$$\int_{\Omega} F'(\hat{u}) \phi dx = \int_{\Omega} j' \left( \frac{\hat{u}}{\|\hat{u}\|} \right) \frac{\phi}{\|\hat{u}\|} dx - \left( \int_{\Omega} j' \left( \frac{\hat{u}}{\|\hat{u}\|} \right) \frac{\hat{u}}{\|\hat{u}\|^3} dx \right) \int_{\Omega} \hat{u} \phi dx$$

which implies

$$F'(\hat{u}) = \frac{1}{\|\hat{u}\|} j' \left( \frac{\hat{u}}{\|\hat{u}\|} \right) - \left( \int_{\Omega} j' \left( \frac{\hat{u}}{\|\hat{u}\|} \right) \frac{\hat{u}}{\|\hat{u}\|^3} dx \right) \hat{u}.$$

6. The variational formulation of the adjoint equation is, by definition,

$$\left\langle \frac{\mathcal{L}}{\partial \hat{u}}(h, \lambda, u, p), \phi \right\rangle = 0 \quad \forall \phi \in H_0^1(\Omega).$$

We compute

$$\left\langle \frac{\mathcal{L}}{\partial \hat{u}}(h, \lambda, u, p), \phi \right\rangle = \int_{\Omega} F'(u) \phi dx + \int_{\Omega} (h \nabla \phi \cdot \nabla p - \lambda \rho h \phi p) dx.$$

7. By integration by parts we find that the adjoint  $p$  is a solution of

$$\begin{cases} -\operatorname{div}(h\nabla p) - \lambda \rho h p = -\frac{1}{\|u\|} j' \left( \frac{u}{\|u\|} \right) + \alpha u & \text{in } \Omega, \\ p = 0 & \text{on } \partial\Omega, \end{cases}$$

with

$$\alpha = \left( \int_{\Omega} j' \left( \frac{u}{\|u\|} \right) \frac{u}{\|u\|^3} dx \right)$$

We check that the right hand side is orthogonal to  $u$  since

$$\int_{\Omega} \frac{u}{\|u\|} j' \left( \frac{u}{\|u\|} \right) dx = \alpha \int_{\Omega} u^2 dx = \alpha \|u\|^2.$$

Clearly, if  $p$  is a solution, then  $p + Cu$  is another possible solution. To determine the value of the constant  $C$  we use

$$\frac{\mathcal{L}}{\partial \hat{\lambda}}(h, \lambda, u, p) = - \int_{\Omega} \rho h u p dx = 0$$

which implies that  $C = - \int_{\Omega} \rho h u p dx$ .

8. The derivative satisfies

$$\langle J'(h), k \rangle = \left\langle \frac{\partial \mathcal{L}}{\partial h}(h, \lambda, u, p), k \right\rangle.$$

We compute

$$\left\langle \frac{\partial \mathcal{L}}{\partial h}(h, \lambda, u, p), k \right\rangle = \int_{\Omega} (k \nabla u \cdot \nabla p - \lambda \rho k u p) dx,$$

which implies

$$J'(h) = \nabla u \cdot \nabla p - \lambda \rho u p.$$

## 2 Geometric optimization: 6 points

1. To define the Lagrangian we introduce two Lagrange multipliers  $q \in H^1(\mathbb{R}^N)$  and  $\mu \in H^1(\mathbb{R}^N)$ , which, together with  $v \in H^1(\mathbb{R}^N)$ , are the arguments of  $\mathcal{L}$

$$\mathcal{L}(\Omega, v, q, \mu) = \int_{\Omega} j(v) dx + \int_{\Omega} (\Delta v + f) q dx + \int_{\partial\Omega} (v - g) \mu ds.$$

2. To get the adjoint problem we differentiate the Lagrangian with respect to  $v$  and set this partial derivative equal to 0. Before that we perform two successive integration by parts

$$\mathcal{L}(\Omega, v, q, \mu) = \int_{\Omega} j(v) dx + \int_{\Omega} (f q - \nabla q \cdot \nabla v) dx + \int_{\partial\Omega} \left( \frac{\partial v}{\partial n} q + (v - g) \mu \right) ds$$

$$\mathcal{L}(\Omega, v, q, \mu) = \int_{\Omega} j(v) dx + \int_{\Omega} (fq + \Delta qv) dx + \int_{\partial\Omega} \left( \frac{\partial v}{\partial n} q - \frac{\partial q}{\partial n} v + (v - g)\mu \right) ds$$

For any  $\phi \in H^1(\mathbb{R}^N)$  we have

$$\begin{aligned} \left\langle \frac{\partial \mathcal{L}}{\partial v}(\Omega, v, q, \mu), \phi \right\rangle &= \int_{\Omega} j'(v)\phi dx + \int_{\Omega} \Delta q\phi dx \\ &+ \int_{\partial\Omega} \left( \frac{\partial \phi}{\partial n} q - \frac{\partial q}{\partial n} \phi + \phi\mu \right) ds \end{aligned} \quad (1)$$

We first take a test function  $\phi$  with compact support in  $\Omega$ , so we deduce that the optimal value of  $q$ , the adjoint  $p$ , satisfies

$$-\Delta p = j'(u) \quad \text{in } \Omega.$$

Then we take  $\phi = 0$  on  $\partial\Omega$  but with no restriction on the value of  $\frac{\partial \phi}{\partial n}$  on  $\partial\Omega$ , so that

$$p = 0 \quad \text{on } \partial\Omega.$$

This yields the adjoint problem

$$\begin{cases} -\Delta p = j'(u) & \text{in } \Omega, \\ p = 0 & \text{on } \partial\Omega. \end{cases}$$

Eventually, varying the trace of  $\phi$  on  $\partial\Omega$  gives the optimal value of the Lagrange multiplier

$$\lambda = \frac{\partial p}{\partial n} \quad \text{on } \partial\Omega.$$

3. Formally we know that the shape derivative is given by

$$J'(\Omega)(\theta) = \frac{\partial \mathcal{L}}{\partial \Omega}(\Omega, u, p)(\theta).$$

We compute the partial derivative

$$\frac{\partial \mathcal{L}}{\partial \Omega}(\Omega, v, q, \mu)(\theta) = \int_{\partial\Omega} \left( j(v) + fq - \nabla q \cdot \nabla v \right) \theta \cdot n ds + \int_{\partial\Omega} \left( \frac{\partial h}{\partial n} + Hh \right) \theta \cdot n ds,$$

with  $h = \frac{\partial v}{\partial n} q + (v - g)\mu$ . Taking into account  $p = 0$  and  $u = g$  on  $\partial\Omega$ , we deduce

$$J'(\Omega)(\theta) = \int_{\partial\Omega} \left( j(u) - \nabla p \cdot \nabla u \right) \theta \cdot n ds + \int_{\partial\Omega} \left( \frac{\partial u}{\partial n} \frac{\partial p}{\partial n} + \mu \frac{\partial(u - g)}{\partial n} \right) \theta \cdot n ds.$$

Since  $\lambda = \frac{\partial p}{\partial n}$  on  $\partial\Omega$  and  $\nabla_t p = 0$  on  $\partial\Omega$ , it leads to

$$J'(\Omega)(\theta) = \int_{\partial\Omega} \left( j(u) + \frac{\partial p}{\partial n} \frac{\partial(u - g)}{\partial n} \right) \theta \cdot n ds.$$