

# CONVERGENCE THEORY

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1. Maximum principle
2. Oscillating test function
3. Two-scale convergence
4. Application to homogenization
5. General theory ( $H$ -convergence)
6. Boundary layers

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# INTRODUCTION

We recall the model problem of diffusion

$$\begin{cases} -\operatorname{div} \left( A \left( \frac{x}{\epsilon} \right) \nabla u_\epsilon \right) = f & \text{in } \Omega \\ u_\epsilon = 0 & \text{on } \partial\Omega \end{cases}$$

where  $\Omega$  is a smooth bounded open set of  $\mathbb{R}^N$ ,  $f \in L^2(\Omega)$  and the coefficient tensor  $A(y)$  is  $Y$ -periodic, uniformly coercive and bounded

$$\alpha |\xi|^2 \leq \sum_{i,j=1}^N A_{ij}(y) \xi_i \xi_j \leq \beta |\xi|^2, \quad \forall \xi \in \mathbb{R}^N, \text{ a.e. } y \in Y \quad (\beta \geq \alpha > 0).$$

**Lemma (a priori estimate).** If  $\Omega$  is bounded, then there exists a constant  $C > 0$  depending only on  $\Omega$  such that

$$\|u_\epsilon\|_{H^1(\Omega)} \leq C \|f\|_{L^2(\Omega)}.$$

Therefore, up to a subsequence, there exists a limit  $u \in H^1(\Omega)$  such that the sequence  $u_\epsilon$  of solutions of the model problem converges to  $u$  weakly in  $H_0^1(\Omega)$  and strongly in  $L^2(\Omega)$ , as  $\epsilon$  goes to 0,

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} |u_\epsilon - u|^2 dx = 0, \quad \lim_{\epsilon \rightarrow 0} \int_{\Omega} \nabla(u_\epsilon - u) \cdot \phi dx = 0 \quad \forall \phi \in L^2(\Omega)^N.$$

**Goal:** find the equation satisfied by  $u$ .

## Proof of the a priori estimate

Variational formulation

$$\int_{\Omega} A\left(x, \frac{x}{\epsilon}\right) \nabla u_{\epsilon}(x) \cdot \nabla \varphi(x) dx = \int_{\Omega} f(x) \varphi(x) dx, \quad \forall \varphi \in H_0^1(\Omega).$$

Take  $\varphi = u_{\epsilon}$  and use coercivity

$$\alpha \|\nabla u_{\epsilon}\|_{L^2(\Omega)}^2 \leq \int_{\Omega} f(x) u_{\epsilon}(x) dx \leq \|f\|_{L^2(\Omega)} \|u_{\epsilon}\|_{L^2(\Omega)}$$

Poincaré inequality in  $\Omega$

$$\|u_{\epsilon}\|_{L^2(\Omega)} \leq C(\Omega) \|\nabla u_{\epsilon}\|_{L^2(\Omega)}$$

Thus

$$\|\nabla u_{\epsilon}\|_{L^2(\Omega)} \leq \frac{C(\Omega) \|f\|_{L^2(\Omega)}}{\alpha}$$

## -I- MAXIMUM PRINCIPLE

Restricted to scalar elliptic equations ! Need some smoothness of the coefficients.

Recall that, at least formally,

$$u_\epsilon(x) \approx u(x) + \epsilon u_1\left(x, \frac{x}{\epsilon}\right) + \epsilon^2 u_2\left(x, \frac{x}{\epsilon}\right)$$

with

$$u_1(x, y) = \sum_{i=1}^N \frac{\partial u}{\partial x_i}(x) w_i(y) + \tilde{u}_1(x)$$

and

$$u_2(x, y) = \sum_{i,j=1}^N \frac{\partial^2 u}{\partial x_i \partial x_j}(x) \chi_{ij}(y) + \sum_{i=1}^N \frac{\partial \tilde{u}_1}{\partial x_i}(x) w_i(y) + \tilde{u}_2(x)$$

Cell problem:

$$\begin{cases} -\operatorname{div}_y A(y) (e_i + \nabla_y w_i(y)) = 0 & \text{in } Y \\ y \rightarrow w_i(y) & Y\text{-periodic} \end{cases}$$

Second order cell problem:

$$\begin{cases} -\operatorname{div}_y A(y) \nabla_y u_2(x, y) = \operatorname{div}_y A(y) \nabla_x u_1 + \operatorname{div}_x A(y) (\nabla_y u_1 + \nabla_x u) \\ \quad \quad \quad \quad \quad \quad \quad - \operatorname{div}_x (A^* \nabla_x u) & \text{in } Y \\ y \rightarrow u_2(x, y) & Y\text{-periodic} \end{cases}$$

which implies the form  $u_2(x, y) = \sum_{i,j=1}^N \frac{\partial^2 u}{\partial x_i \partial x_j}(x) \chi_{ij}(y)$  (up to addition of  $\tilde{u}_1, \tilde{u}_2$ ).

Define the remainder

$$r_\epsilon(x) = u_\epsilon(x) - \left\{ u(x) + \epsilon u_1 \left( x, \frac{x}{\epsilon} \right) + \epsilon^2 u_2 \left( x, \frac{x}{\epsilon} \right) \right\}$$

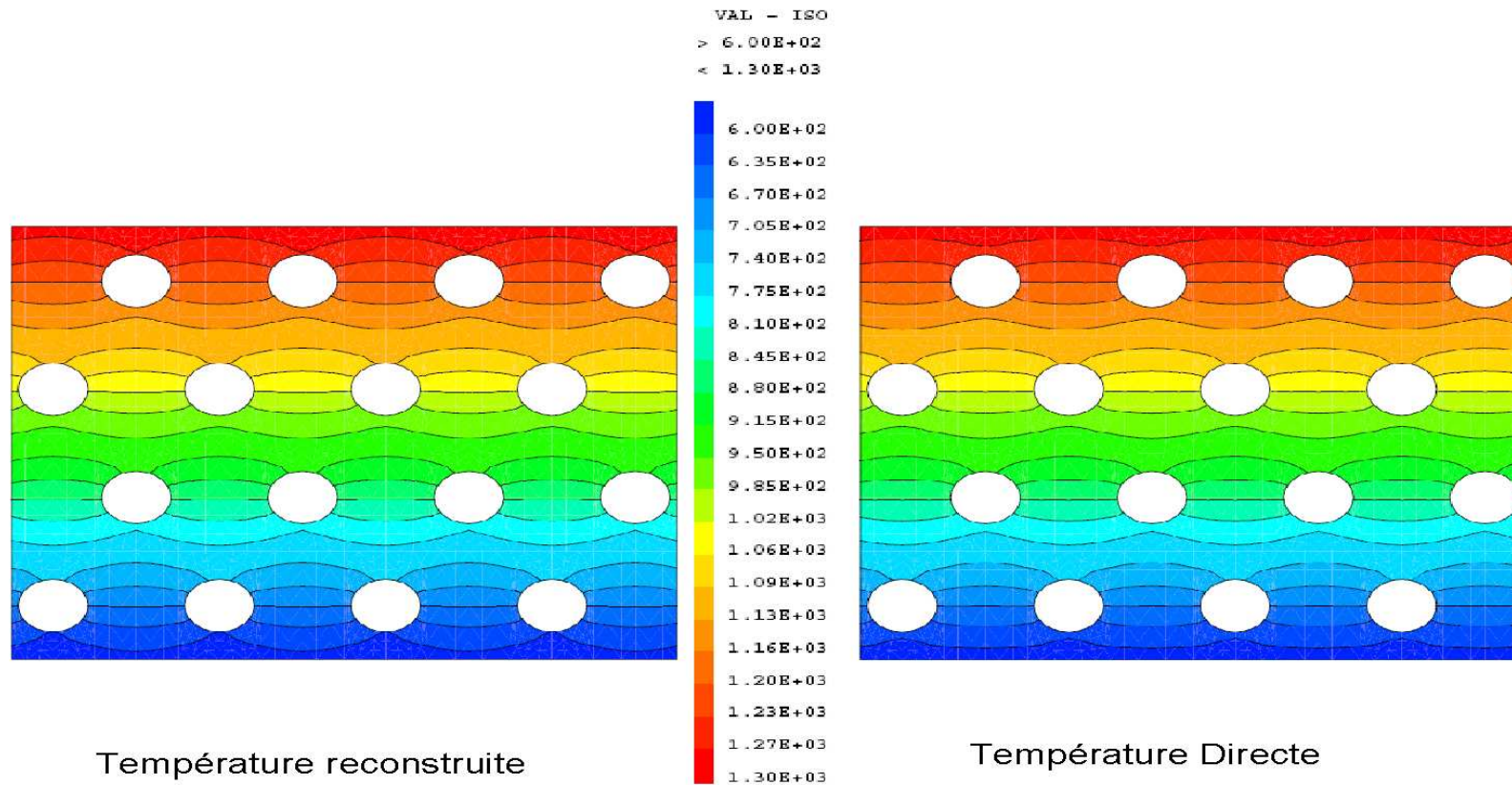
which satisfy

$$\begin{cases} -\operatorname{div} \left( A \left( \frac{x}{\epsilon} \right) \nabla r_\epsilon \right) = -\epsilon [\operatorname{div}_x A(\nabla_x u_1 + \nabla_y u_2) + \operatorname{div}_y A(\nabla_x u_2)] \left( x, \frac{x}{\epsilon} \right) \\ \quad - \epsilon^2 [\operatorname{div}_x A(\nabla_x u_2)] \left( x, \frac{x}{\epsilon} \right) & \text{in } \Omega \\ r_\epsilon = -\epsilon u_1 \left( x, \frac{x}{\epsilon} \right) - \epsilon^2 u_2 \left( x, \frac{x}{\epsilon} \right) & \text{on } \partial\Omega \end{cases}$$

**Theorem.** Assume  $A(y) \in W^{1,\infty}(Y)$  and  $u \in W^{4,\infty}(\Omega)$ . Then

$$\|u_\epsilon(x) - u(x)\|_{L^\infty(\Omega)} \leq C\epsilon$$

**Proof.** Apply the maximum principle to  $r_\epsilon$ .



Comparison of the exact and reconstructed solution  $u + \epsilon u_1$  (K. El Ganaoui)  
 Correctors really matter !



## -II- Oscillating test function method (L. Tartar)

The goal is to **mathematically** prove the following

**Theorem.** The sequence  $u_\epsilon(x)$  of solutions of the model problem converges weakly in  $H_0^1(\Omega)$  and strongly in  $L^2(\Omega)$ , as  $\epsilon$  goes to 0, to a limit  $u(x)$  which is the unique solution of the homogenized problem.

**Naive idea:** pass to the limit in the variational formulation

$$\int_{\Omega} A\left(\frac{x}{\epsilon}\right) \nabla u_\epsilon(x) \cdot \nabla \varphi(x) dx = \int_{\Omega} f(x) \varphi(x) dx, \quad \forall \varphi \in H_0^1(\Omega).$$

**impossible** because  $A\left(\frac{x}{\epsilon}\right)$  and  $\nabla u_\epsilon$  converge only weakly (the limit of the product is not the product of the limits).

## Difference between strong and weak convergence

$u_\epsilon$  converges strongly to  $u$  in  $L^2(\Omega) \Leftrightarrow \lim_{\epsilon \rightarrow 0} \int_{\Omega} |u_\epsilon(x) - u(x)|^2 dx = 0$

$v_\epsilon$  converges weakly to  $v$  in  $L^2(\Omega) \Leftrightarrow \lim_{\epsilon \rightarrow 0} \int_{\Omega} (v_\epsilon - v)\phi dx = 0 \quad \forall \phi \in L^2(\Omega)$

**Lemma.** Let  $G : \mathbb{R} \rightarrow \mathbb{R}$  be a continuous bounded function. Then  $G(u_\epsilon)$  converges strongly to  $G(u)$  in  $L^2(\Omega)$ .

This is false for weak convergence ! Counter-example:  $v_\epsilon(x) = \sin\left(\frac{x}{\epsilon}\right)$ .

## Oscillating test function method (ctd.)

The **key idea** is to replace  $\varphi$  by

$$\varphi_\epsilon(x) = \varphi(x) + \epsilon \sum_{i=1}^N \frac{\partial \varphi}{\partial x_i}(x) w_i \left( \frac{x}{\epsilon} \right)$$

when  $A$  is symmetric. Then, one can pass to the limit in the variational formulation **using the cell equation satisfied by  $w_i(y)$** . Remark that the difference between the gradients of  $\varphi_\epsilon$  and  $\varphi$  is not small

$$\nabla \varphi_\epsilon(x) = \nabla \varphi(x) + \sum_{i=1}^N \frac{\partial \varphi}{\partial x_i}(x) (\nabla_y w_i) \left( \frac{x}{\epsilon} \right) + \mathcal{O}(\epsilon).$$

This idea can be used in **multiscale numerical computations**: replace the usual piecewise affine finite element basis by **oscillating** functions.

## -III- TWO-SCALE CONVERGENCE METHOD

This is a simpler method to prove the convergence theorem.

**Definition.** A sequence of functions  $u_\epsilon$  in  $L^2(\Omega)$  is said to **two-scale converge** to a limit  $u_0(x, y)$  belonging to  $L^2(\Omega \times Y)$  if, for any  $Y$ -periodic smooth function  $\varphi(x, y)$ , it satisfies

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} u_\epsilon(x) \varphi\left(x, \frac{x}{\epsilon}\right) dx = \int_{\Omega} \int_Y u_0(x, y) \varphi(x, y) dx dy.$$

**Theorem 1.** From each bounded sequence  $u_\epsilon$  in  $L^2(\Omega)$  one can extract a subsequence, and there exists a limit  $u_0(x, y) \in L^2(\Omega \times Y)$  such that this subsequence two-scale converges to  $u_0$ .

**Theorem 2.** Let  $u_\epsilon$  be a bounded sequence in  $H^1(\Omega)$ . Then, up to a subsequence,  $u_\epsilon$  two-scale converges to a limit  $u_0(x, y) \equiv u(x) \in H^1(\Omega)$ , and  $\nabla u_\epsilon$  two-scale converges to  $\nabla_x u(x) + \nabla_y u_1(x, y)$  with  $u_1 \in L^2(\Omega; H^1_{\#}(Y))$ .

**Lemma 3.** For a bounded open set  $\Omega$ , let  $B = C(\bar{\Omega}; C_{\#}(Y))$  be the space of continuous functions  $\varphi(x, y)$  on  $\bar{\Omega} \times Y$  which are  $Y$ -periodic in  $y$ . Then,  $B$  is a separable Banach space (i.e. it contains a dense countable family), is dense in  $L^2(\Omega \times Y)$ , and there exists  $C > 0$  such that

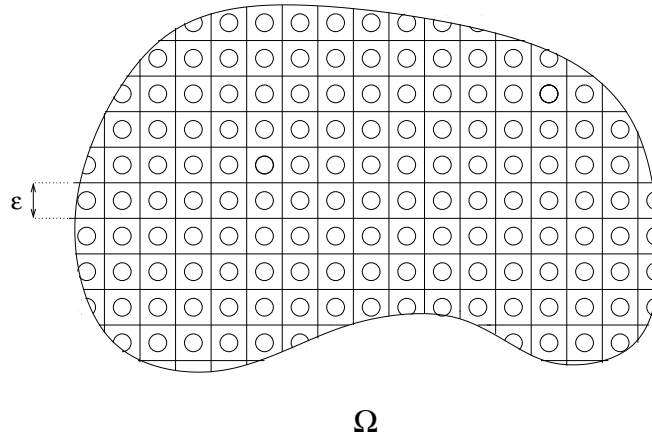
$$\int_{\Omega} \left| \varphi \left( x, \frac{x}{\epsilon} \right) \right|^2 dx \leq C \|\varphi\|_B^2,$$

and

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} \left| \varphi \left( x, \frac{x}{\epsilon} \right) \right|^2 dx = \int_{\Omega} \int_Y |\varphi(x, y)|^2 dx dy,$$

for any  $\varphi(x, y) \in B$ .

**Remark.** The same works with  $B = L^2(\Omega; C_{\#}(Y))$  and  $\Omega$  not necessarily bounded.



**Proof of Lemma 3.** We mesh  $\Omega$  with cubes of the type  $(0, \epsilon)^N$

$$\Omega = \cup_{1 \leq i \leq n(\epsilon)} Y_i^\epsilon \quad \text{with } Y_i^\epsilon = x_i^\epsilon + (0, \epsilon)^N.$$

$$\begin{aligned} \int_{\Omega} \left| \varphi \left( x, \frac{x}{\epsilon} \right) \right|^2 dx &= \sum_{i=1}^{n(\epsilon)} \int_{Y_i^\epsilon} \left| \varphi \left( x, \frac{x}{\epsilon} \right) \right|^2 dx = \sum_{i=1}^{n(\epsilon)} \int_{Y_i^\epsilon} \left| \varphi \left( x_i^\epsilon, \frac{x}{\epsilon} \right) \right|^2 dx + o(1) \\ &= \sum_{i=1}^{n(\epsilon)} \epsilon^N \int_Y \left| \varphi \left( x_i^\epsilon, y \right) \right|^2 dy + o(1) = \int_{\Omega} \int_Y \left| \varphi \left( x, y \right) \right|^2 dx dy + o(1) \end{aligned}$$

**Proof of Theorem 1.** By Schwarz inequality, we have

$$\left| \int_{\Omega} u_{\epsilon}(x) \varphi \left( x, \frac{x}{\epsilon} \right) dx \right| \leq C \left| \int_{\Omega} \varphi \left( x, \frac{x}{\epsilon} \right) dx \right|^{\frac{1}{2}} \leq C \|\varphi\|_B.$$

Thus, the l.h.s. is a **continuous linear form on  $B$**  which can be identified to a duality product  $\langle \mu_{\epsilon}, \varphi \rangle_{B', B}$  for some bounded sequence of measures  $\mu_{\epsilon}$ . **Since  $B$  is separable**, by the Banach-Alaoglu theorem, one can extract a subsequence and there exists a limit  $\mu_0$  such  $\mu_{\epsilon}$  converges to  $\mu_0$  in the weak \* topology of  $B'$  (the dual of  $B$ ). On the other hand, Lemma 3 allows us to pass to the limit in the middle term above. It yields

$$|\langle \mu_0, \varphi \rangle_{B', B}| \leq C \left| \int_{\Omega} \int_Y |\varphi(x, y)|^2 dx dy \right|^{\frac{1}{2}}.$$

Therefore  $\mu_0$  is actually a **continuous linear form on  $L^2(\Omega \times Y)$** , by density of  $B$  in this space. Thus, there exists  $u_0(x, y) \in L^2(\Omega \times Y)$  such that

$$\langle \mu_0, \varphi \rangle_{B', B} = \int_{\Omega} \int_Y u_0(x, y) \varphi(x, y) dx dy.$$

**Proof of Theorem 2.** Since  $u_\epsilon$  and  $\nabla u_\epsilon$  are bounded in  $L^2(\Omega)$ , up to a subsequence, they two-scale converge to limits  $u_0(x, y) \in L^2(\Omega \times Y)$  and  $\xi_0(x, y) \in L^2(\Omega \times Y)^N$ . Thus

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} \nabla u_\epsilon(x) \cdot \vec{\psi} \left( x, \frac{x}{\epsilon} \right) dx = \int_{\Omega} \int_Y \xi_0(x, y) \cdot \vec{\psi}(x, y) dx dy \quad \forall \vec{\psi} \in \mathcal{D}(\Omega; C_{\#}^\infty(Y)^N).$$

Integrating by parts the left hand side gives

$$\epsilon \int_{\Omega} \nabla u_\epsilon(x) \cdot \vec{\psi} \left( x, \frac{x}{\epsilon} \right) dx = - \int_{\Omega} u_\epsilon(x) \left( \operatorname{div}_Y \vec{\psi} \left( x, \frac{x}{\epsilon} \right) + \epsilon \operatorname{div}_X \vec{\psi} \left( x, \frac{x}{\epsilon} \right) \right) dx.$$

Passing to the limit yields

$$0 = - \int_{\Omega} \int_Y u_0(x, y) \operatorname{div}_Y \vec{\psi}(x, y) dx dy \quad \Rightarrow \quad u_0(x, y) \equiv u(x) \in L^2(\Omega).$$

Next, we choose  $\vec{\psi}$  such that  $\operatorname{div}_Y \vec{\psi}(x, y) = 0$ . We obtain

$$\int_{\Omega} \nabla u_\epsilon(x) \cdot \vec{\psi} \left( x, \frac{x}{\epsilon} \right) dx = - \int_{\Omega} u_\epsilon(x) \operatorname{div}_X \vec{\psi} \left( x, \frac{x}{\epsilon} \right) dx.$$



Passing to the two-scale limit

$$\int_{\Omega} \int_Y \xi_0(x, y) \cdot \vec{\psi}(x, y) dx dy = - \int_{\Omega} \int_Y u(x) \operatorname{div}_x \vec{\psi}(x, y) dx dy.$$

If  $\vec{\psi}$  does not depend on  $y$ , it proves that  $u(x) \in H^1(\Omega)$ . Furthermore,

$$\int_{\Omega} \int_Y (\xi_0(x, y) - \nabla u(x)) \cdot \vec{\psi}(x, y) dx dy = 0 \quad \forall \vec{\psi} \text{ with } \operatorname{div}_y \vec{\psi} = 0.$$

The orthogonal of divergence-free functions are exactly the gradients. Thus, there exists a unique function  $u_1(x, y)$  in  $L^2(\Omega; H_{\#}^1(Y)/\mathbb{R})$  such that

$$\xi_0(x, y) = \nabla u(x) + \nabla_y u_1(x, y).$$

**Theorem 4.** Let  $u_\epsilon \in L^2(\Omega)$  two-scale converge to  $u_0(x, y) \in L^2(\Omega \times Y)$ .

1. Then,  $u_\epsilon$  converges weakly in  $L^2(\Omega)$  to  $u(x) = \int_Y u_0(x, y) dy$ , and we have

$$\lim_{\epsilon \rightarrow 0} \|u_\epsilon\|_{L^2(\Omega)}^2 \geq \|u_0\|_{L^2(\Omega \times Y)}^2 \geq \|u\|_{L^2(\Omega)}^2.$$

2. Assume further that  $u_0(x, y)$  is smooth and that

$$\lim_{\epsilon \rightarrow 0} \|u_\epsilon\|_{L^2(\Omega)}^2 = \|u_0\|_{L^2(\Omega \times Y)}^2.$$

Then, we have

$$\left\| u_\epsilon(x) - u_0\left(x, \frac{x}{\epsilon}\right) \right\|_{L^2(\Omega)}^2 \rightarrow 0.$$

**Remark.** In the last case we say that  $u_\epsilon$  two-scale converges **strongly** to  $u_0$ .

**Proof of Theorem 4.** Take a test function depending only on  $x$

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} u_{\epsilon}(x) \varphi(x) dx = \int_{\Omega} \int_Y u_0(x, y) \varphi(x) dx dy = \int_{\Omega} u(x) \varphi(x) dx.$$

Thus,  $u_{\epsilon}$  converges weakly to  $u$  in  $L^2(\Omega)$ . Then, developing the inequality

$$\int_{\Omega} \left| u_{\epsilon}(x) - \varphi\left(x, \frac{x}{\epsilon}\right) \right|^2 dx \geq 0$$

$$\int_{\Omega} |u_{\epsilon}(x)|^2 dx - 2 \int_{\Omega} u_{\epsilon}(x) \varphi\left(x, \frac{x}{\epsilon}\right) dx + \int_{\Omega} \left| \varphi\left(x, \frac{x}{\epsilon}\right) \right|^2 dx \geq 0$$

$$\liminf_{\epsilon \rightarrow 0} \int_{\Omega} |u_{\epsilon}(x)|^2 dx - 2 \int_{\Omega} \int_Y u_0(x, y) \varphi(x, y) dx dy + \int_{\Omega} \int_Y |\varphi(x, y)|^2 dx dy \geq 0$$

Take  $\varphi = u_0$  to get

$$\lim_{\epsilon \rightarrow 0} \|u_{\epsilon}\|_{L^2(\Omega)}^2 \geq \|u_0\|_{L^2(\Omega \times Y)}^2.$$

**Proof of Theorem 4 (continued).** If we assume

$$\lim_{\epsilon \rightarrow 0} \|u_\epsilon\|_{L^2(\Omega)}^2 = \|u_0\|_{L^2(\Omega \times Y)}^2,$$

the same computation yields

$$\liminf_{\epsilon \rightarrow 0} \int_{\Omega} \left| u_\epsilon(x) - \varphi\left(x, \frac{x}{\epsilon}\right) \right|^2 dx = \int_{\Omega} \int_Y |u_0(x, y) - \varphi(x, y)|^2 dx dy$$

If  $u_0$  is smooth enough to be a test function  $\varphi$  (Carathéodory function), it gives the desired result

$$\left\| u_\epsilon(x) - u_0\left(x, \frac{x}{\epsilon}\right) \right\|_{L^2(\Omega)}^2 \rightarrow 0.$$

### Theorem 5.

1. Let  $u_\epsilon$  be a bounded sequence in  $L^2(\Omega)$  such that  $\epsilon \nabla u_\epsilon$  is also bounded in  $L^2(\Omega)^N$ . Then, there exists a two-scale limit  $u_0(x, y) \in L^2(\Omega; H_{\#}^1(Y)/\mathbb{R})$  such that, up to a subsequence,  $u_\epsilon$  two-scale converges to  $u_0(x, y)$ , and  $\epsilon \nabla u_\epsilon$  to  $\nabla_y u_0(x, y)$ .
2. Let  $u_\epsilon$  be a bounded sequence in  $L^2(\Omega)^N$  such that  $\operatorname{div} u_\epsilon$  is also bounded in  $L^2(\Omega)$ . Then, there exists a two-scale limit  $u_0(x, y) \in L^2(\Omega \times Y)^N$  with  $\operatorname{div}_y u_0 = 0$  and  $\operatorname{div}_x u_0 \in L^2(\Omega \times Y)$  such that, up to a subsequence,  $u_\epsilon$  two-scale converges to  $u_0(x, y)$ , and  $\operatorname{div} u_\epsilon$  to  $\operatorname{div}_x u_0(x, y)$ .

**Proof of Theorem 5 (first part).** We have

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} u_{\epsilon}(x) \varphi \left( x, \frac{x}{\epsilon} \right) dx = \int_{\Omega} \int_Y u_0(x, y) \varphi(x, y) dx dy$$

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} \epsilon \nabla u_{\epsilon}(x) \cdot \vec{\psi} \left( x, \frac{x}{\epsilon} \right) dx = \int_{\Omega} \int_Y \xi_0(x, y) \cdot \vec{\psi}(x, y) dx dy$$

By integration by parts

$$\int_{\Omega} \epsilon \nabla u_{\epsilon}(x) \cdot \vec{\psi} \left( x, \frac{x}{\epsilon} \right) dx = - \int_{\Omega} u_{\epsilon}(x) \left( \operatorname{div}_y \vec{\psi} + \epsilon \operatorname{div}_x \vec{\psi} \right) \left( x, \frac{x}{\epsilon} \right) dx$$

Passing to the two-scale limit

$$\int_{\Omega} \int_Y \xi_0(x, y) \cdot \vec{\psi}(x, y) dx dy = - \int_{\Omega} \int_Y u_0(x, y) \operatorname{div}_y \vec{\psi}(x, y) dx dy$$

which implies that  $\xi_0(x, y) = \nabla_y u_0(x, y)$ .

## -IV- APPLICATION TO HOMOGENIZATION

Conductivity or diffusion equation

$$\begin{cases} -\operatorname{div} \left( A \left( x, \frac{x}{\epsilon} \right) \nabla u_\epsilon \right) = f & \text{in } \Omega \\ u_\epsilon = 0 & \text{on } \partial\Omega \end{cases}$$

with a coefficient tensor  $A(x, y)$  which is  $Y$ -periodic, uniformly coercive and bounded

$$\alpha|\xi|^2 \leq \sum_{i,j=1}^N A_{ij}(x, y)\xi_i\xi_j \leq \beta|\xi|^2, \quad \forall \xi \in \mathbb{R}^N, \forall y \in Y, \forall x \in \Omega \quad (\beta \geq \alpha > 0).$$

**A priori estimate.** If  $\Omega$  is bounded, then

$$\|u_\epsilon\|_{H^1(\Omega)} \leq C\|f\|_{L^2(\Omega)}.$$

## Two-scale convergence method

**First step.** We deduce from the a priori estimates [the precise form](#) of the two-scale limit of the sequence  $u_\epsilon$ .

By application of Theorem 2, there exist two functions,  $u(x) \in H_0^1(\Omega)$  and  $u_1(x, y) \in L^2(\Omega; H_{\#}^1(Y)/\mathbb{R})$ , such that, up to a subsequence,  $u_\epsilon$  two-scale converges to  $u(x)$ , and  $\nabla u_\epsilon$  two-scale converges to  $\nabla_x u(x) + \nabla_y u_1(x, y)$ .

In view of these limits,  $u_\epsilon$  is expected to behave as  $u(x) + \epsilon u_1(x, \frac{x}{\epsilon})$ .



## Two-scale convergence method

**Second step.** We multiply the p.d.e. by a test function similar to the limit of  $u_\epsilon$ , namely  $\varphi(x) + \epsilon\varphi_1\left(x, \frac{x}{\epsilon}\right)$ , where  $\varphi(x) \in \mathcal{D}(\Omega)$  and  $\varphi_1(x, y) \in \mathcal{D}(\Omega; C_\#^\infty(Y))$ . This yields

$$\begin{aligned} \int_{\Omega} A\left(x, \frac{x}{\epsilon}\right) \nabla u_\epsilon \cdot \left( \nabla \varphi(x) + \nabla_y \varphi_1\left(x, \frac{x}{\epsilon}\right) + \epsilon \nabla_x \varphi_1\left(x, \frac{x}{\epsilon}\right) \right) dx \\ = \int_{\Omega} f(x) \left( \varphi(x) + \epsilon \varphi_1\left(x, \frac{x}{\epsilon}\right) \right) dx. \end{aligned}$$

Regarding  $A^t\left(x, \frac{x}{\epsilon}\right) \left( \nabla \varphi(x) + \nabla_y \varphi_1\left(x, \frac{x}{\epsilon}\right) \right)$  as a test function for the two-scale convergence, we pass to the two-scale limit

$$\int_{\Omega} \int_Y A(x, y) (\nabla u(x) + \nabla_y u_1(x, y)) \cdot (\nabla \varphi(x) + \nabla_y \varphi_1(x, y)) dx dy = \int_{\Omega} f(x) \varphi(x) dx.$$

## Two-scale convergence method

**Third step.** We read off a variational formulation for  $(u, u_1)$ . Take  $(\varphi, \varphi_1)$  in the Hilbert space  $H_0^1(\Omega) \times L^2\left(\Omega; H_{\#}^1(Y)/\mathbb{R}\right)$  endowed with the norm

$$\sqrt{(\|\nabla u(x)\|_{L^2(\Omega)}^2 + \|\nabla_y u_1(x, y)\|_{L^2(\Omega \times Y)}^2)}$$

The assumptions of the Lax-Milgram lemma are easily checked. The main point is the coercivity of the bilinear form defined by the left hand side

$$\begin{aligned} & \int_{\Omega} \int_Y A(x, y) (\nabla \varphi(x) + \nabla_y \varphi_1(x, y)) \cdot (\nabla \varphi(x) + \nabla_y \varphi_1(x, y)) \, dx dy \geq \\ & \alpha \int_{\Omega} \int_Y |\nabla \varphi(x) + \nabla_y \varphi_1(x, y)|^2 \, dx dy = \alpha \int_{\Omega} |\nabla \varphi(x)|^2 \, dx + \alpha \int_{\Omega} \int_Y |\nabla_y \varphi_1(x, y)|^2 \, dx dy. \end{aligned}$$

$$\int_{\Omega} \int_Y A(x, y) (\nabla u(x) + \nabla_y u_1(x, y)) \cdot (\nabla \varphi(x) + \nabla_y \varphi_1(x, y)) dx dy = \int_{\Omega} f(x) \varphi(x) dx.$$

By application of the Lax-Milgram lemma, there exists a unique solution  $(u, u_1) \in H_0^1(\Omega) \times L^2\left(\Omega; H_{\#}^1(Y)/\mathbb{R}\right)$ . Consequently, the **entire sequences**  $u_\epsilon$  and  $\nabla u_\epsilon$  converge to  $u(x)$  and  $\nabla u(x) + \nabla_y u_1(x, y)$ .

An easy integration by parts shows that the associated p.d.e.'s are the so-called “two-scale homogenized problem”,

$$\left\{ \begin{array}{ll} -\operatorname{div}_y (A(x, y) (\nabla u(x) + \nabla_y u_1(x, y))) = 0 & \text{in } \Omega \times Y \\ -\operatorname{div}_x \left( \int_Y A(x, y) (\nabla u(x) + \nabla_y u_1(x, y)) dy \right) = f(x) & \text{in } \Omega \\ y \rightarrow u_1(x, y) & Y\text{-periodic} \\ u = 0 & \text{on } \partial\Omega. \end{array} \right.$$

## Two-scale convergence method

**Fourth (and optional) step.** Eliminate the  $y$  variable and the  $u_1$  unknown

$$u_1(x, y) = \sum_{i=1}^N \frac{\partial u}{\partial x_i}(x) w_i(x, y),$$

where  $w_i(x, y)$  are the unique solutions in  $H_{\#}^1(Y)/\mathbb{R}$  of the cell problems

$$\begin{cases} -\operatorname{div}_y (A(x, y) (\vec{e}_i + \nabla_y w_i(x, y))) = 0 & \text{in } Y \\ y \rightarrow w_i(x, y) & Y\text{-periodic,} \end{cases}$$

at each point  $x \in \Omega$ , and

$$\begin{cases} -\operatorname{div}_x (A^*(x) \nabla u(x)) = f(x) & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

with  $A_{ij}^*(x) = \int_Y A(x, y) (\vec{e}_i + \nabla_y w_i(x, y)) \cdot (\vec{e}_j + \nabla_y w_j(x, y)) dy$ .

## Conclusion

We have just proved.

**Theorem.** The sequence of solution  $u_\epsilon$  converges weakly to  $u$  in  $H_0^1(\Omega)$  and the sequence  $\nabla u_\epsilon$  two-scale converges to  $\nabla_x u(x) + \nabla_y u_1(x, y)$  where  $u$  is the solution of the homogenized problem and  $u_1(x, y) = \sum_{i=1}^N \frac{\partial u}{\partial x_i}(x) w_i(x, y)$ .

**Remark.**  $u_\epsilon$  converges strongly to  $u$  in  $L^2(\Omega)$  but its gradient does not converge strongly !

## Corrector and strong two-scale convergence

**Proposition.** Assume that  $A(x, y) \in C(\bar{\Omega}; L^\infty_{\#}(Y))$ . Then  $u_1(x, y)$  is indeed a corrector in the sense that

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} \left| \nabla u_{\epsilon}(x) - \nabla u(x) - \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right|^2 dx = 0$$

**Key ingredient:** strong two-scale convergence. If  $v_{\epsilon} \in L^2(\Omega)$  two-scale converge to a smooth  $v_0(x, y)$  and

$$\lim_{\epsilon \rightarrow 0} \|v_{\epsilon}\|_{L^2(\Omega)}^2 = \|v_0\|_{L^2(\Omega \times Y)}^2,$$

then

$$\lim_{\epsilon \rightarrow 0} \left\| v_{\epsilon}(x) - v_0 \left( x, \frac{x}{\epsilon} \right) \right\|_{L^2(\Omega)}^2 = 0.$$

## Proof

We develop

$$\begin{aligned}
& \int_{\Omega} A \left( x, \frac{x}{\epsilon} \right) \left( \nabla u_{\epsilon} - \nabla u - \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) \cdot \left( \nabla u_{\epsilon} - \nabla u - \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) dx = \\
& \int_{\Omega} A \left( x, \frac{x}{\epsilon} \right) \left( \nabla u(x) + \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) \cdot \left( \nabla u(x) + \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) dx \\
& \quad - 2 \int_{\Omega} A^{sym} \left( x, \frac{x}{\epsilon} \right) \nabla u_{\epsilon} \cdot \left( \nabla u(x) + \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) dx \\
& \quad + \int_{\Omega} A \left( x, \frac{x}{\epsilon} \right) \nabla u_{\epsilon} \cdot \nabla u_{\epsilon} dx
\end{aligned}$$

The last term is equal to  $\int_{\Omega} f u_{\epsilon} dx \rightarrow \int_{\Omega} f u dx$ .

## Proof (ctd.)

Passing to the limit

$$\begin{aligned}
 & \lim_{\epsilon \rightarrow 0} \alpha \int_{\Omega} \left| \nabla u_{\epsilon}(x) - \nabla u(x) - \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right|^2 dx \leq \\
 & \lim_{\epsilon \rightarrow 0} \int_{\Omega} A \left( x, \frac{x}{\epsilon} \right) \left( \nabla u_{\epsilon} - \nabla u - \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) \cdot \left( \nabla u_{\epsilon} - \nabla u - \nabla_y u_1 \left( x, \frac{x}{\epsilon} \right) \right) dx = \\
 & \qquad \qquad \qquad \int_{\Omega} f u dx \\
 & -2 \int_{\Omega} \int_Y A^{sym}(x, y) (\nabla u(x) + \nabla_y u_1(x, y)) \cdot (\nabla u(x) + \nabla_y u_1(x, y)) dx dy \\
 & + \int_{\Omega} \int_Y A(x, y) (\nabla u(x) + \nabla_y u_1(x, y)) \cdot (\nabla u(x) + \nabla_y u_1(x, y)) dx dy \\
 & \qquad \qquad \qquad = 0
 \end{aligned}$$

because of the variational formulation of the two-scale limit problem !



## Error estimate

**Proposition.** Assume that  $A(x, y) \in C(\bar{\Omega}; L^\infty_\#(Y))$  and  $u \in W^{2,\infty}(\Omega)$ . Then

$$\left\| u_\epsilon(x) - u(x) - \epsilon u_1\left(x, \frac{x}{\epsilon}\right) \right\|_{H^1(\Omega)} \leq C\sqrt{\epsilon}$$

## -V- GENERAL THEORY (H-convergence)

Let  $\Omega$  be a bounded open set in  $\mathbb{R}^N$ , and let  $0 < \alpha \leq \beta$ . We introduce the set  $\mathcal{M}(\alpha, \beta, \Omega)$  of all possible matrices  $A(x)$  such that, a.e. in  $\Omega$

$$\alpha|\xi|^2 \leq A(x)\xi \cdot \xi \quad \text{and} \quad \beta|\xi|^2 \leq A^{-1}(x)\xi \cdot \xi \quad \forall \xi \in \mathbb{R}^N.$$

We consider a sequence  $A_\epsilon(x)$  in  $\mathcal{M}(\alpha, \beta, \Omega)$ , indexed by  $\epsilon$ , going to 0, which is not associated to any specific lengthscale or statistical property of the elastic medium. In other words, **no special assumptions** (like periodicity or stationarity) are placed on the sequence  $A_\epsilon$ .

For  $f(x) \in L^2(\Omega)$ , there exists a unique solution  $u_\epsilon \in H_0^1(\Omega)$  of

$$\begin{cases} -\operatorname{div}(A_\epsilon(x)\nabla u_\epsilon) = f(x) & \text{in } \Omega \\ u_\epsilon = 0 & \text{on } \partial\Omega. \end{cases}$$

**Definition.** The sequence  $A_\epsilon(x) \in \mathcal{M}(\alpha, \beta, \Omega)$  is said to  $H$ -converge to a limit  $A^*(x)$ , as  $\epsilon$  goes to 0, if, for any  $f \in L^2(\Omega)$ , the sequence  $u_\epsilon$  satisfies

$$\begin{cases} u_\epsilon \rightharpoonup u \text{ weakly in } H_0^1(\Omega) \\ A^\epsilon \nabla u_\epsilon \rightharpoonup A^* \nabla u \text{ weakly in } L^2(\Omega)^N, \end{cases}$$

where  $u$  is the solution of the homogenized equation associated to  $A^*$

$$\begin{cases} -\operatorname{div}(A^*(x)\nabla u) = f(x) & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

Remark that, by definition, the homogenized tensor  $A^*$  is independent of the source term  $f$ . We shall see that it is also independent of the boundary condition and of the domain.

**Theorem (compactness).** For any sequence  $A_\epsilon$  in  $\mathcal{M}(\alpha, \beta, \Omega)$ , there exist a subsequence (still denoted by  $\epsilon$ ) and a homogenized limit  $A^*$ , belonging to  $\mathcal{M}(\alpha, \beta, \Omega)$ , such that  $A_\epsilon$   $H$ -converges to  $A^*$ .

**Proposition (locality of  $H$ -convergence).** Let  $A^\epsilon(x)$  and  $B^\epsilon(x)$  be two sequences in  $\mathcal{M}(\alpha, \beta, \Omega)$ , which  $H$ -converge to  $A^*(x)$  and  $B^*(x)$ , respectively. Let  $\omega$  be an open subset compactly embedded in  $\Omega$ , i.e.,  $\bar{\omega} \subset \Omega$ . If  $A^\epsilon(x) = B^\epsilon(x)$  in  $\omega$ , then  $A^*(x) = B^*(x)$  in  $\omega$ .

**Proposition (energy convergence).** Let  $A^\epsilon(x)$  be a sequence in  $\mathcal{M}(\alpha, \beta, \Omega)$  that  $H$ -converges to  $A^*(x)$ . For any  $f \in L^2(\Omega)$ , the sequence  $u_\epsilon$  satisfies

$$A^\epsilon \nabla u_\epsilon \cdot \nabla u_\epsilon \rightharpoonup A^* \nabla u \cdot \nabla u \text{ in } L^1(\Omega)$$

and

$$\int_{\Omega} A^\epsilon \nabla u_\epsilon \cdot \nabla u_\epsilon dx \rightarrow \int_{\Omega} A^* \nabla u \cdot \nabla u dx,$$

## G-convergence

When restricted to symmetric matrices,  $H$  convergence simplifies and is called  $G$ -convergence.

**Proposition ( $G$ -convergence).** Let  $A^\epsilon(x)$  be a sequence of **symmetric** matrices in  $\mathcal{M}(\alpha, \beta, \Omega)$  that  $H$ -converges to  $A^*(x)$ . Then  $A^*$  is symmetric too. Furthermore,  $H$ -convergence is equivalent to so-called  $G$ -convergence which is defined by: for any  $f \in L^2(\Omega)$ ,  $u_\epsilon \rightharpoonup u$  weakly in  $H_0^1(\Omega)$  where  $u$  is the solution of the homogenized equation.

**Remark.** The convergence of the flux  $A^\epsilon \nabla u_\epsilon$  **is not required** in  $G$ -convergence and is rather an automatic property (due to the symmetry).

**Bibliography.**  $H$ -convergence was introduced by Murat and Tartar.  $G$ -convergence was older and due to De Giorgi and Spagnolo.

## -VI- BOUNDARY LAYERS

$$u_\epsilon(x) \approx u(x) + \epsilon u_1\left(x, \frac{x}{\epsilon}\right)$$

but  $u_1$  does not satisfy the Dirichlet boundary condition.

We correct it by introducing a **boundary layer** solution of

$$\begin{cases} -\operatorname{div}\left(A\left(\frac{x}{\epsilon}\right)\nabla u_\epsilon^{bl}\right) = 0 & \text{in } \Omega \\ u_\epsilon^{bl} = -u_1\left(x, \frac{x}{\epsilon}\right) & \text{on } \partial\Omega \end{cases}$$

which cannot be computed explicitly except in special cases. In particular

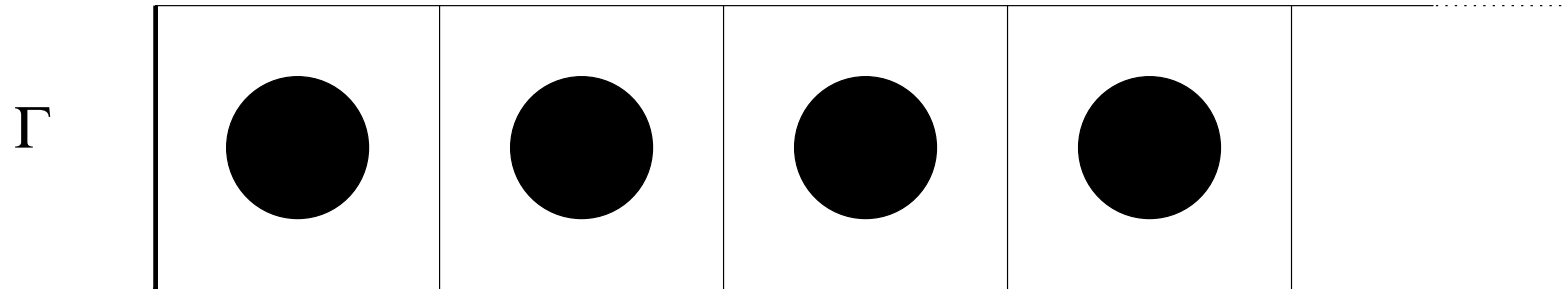
$$\|u_\epsilon^{bl}\|_{H^1(\Omega)} = \mathcal{O}\left(\frac{1}{\sqrt{\epsilon}}\right)$$

**Proposition.** Assume that  $u \in W^{2,\infty}(\Omega)$ . Then

$$\left\| u_\epsilon(x) - u(x) - \epsilon u_1\left(x, \frac{x}{\epsilon}\right) - \epsilon u_\epsilon^{bl}(x) \right\|_{H^1(\Omega)} \leq C\epsilon$$

Consider a rectangular domain  $\Omega$ . In such a case we can approximate the boundary layer as follows.

Introduce the semi-infinite band  $G$  (in the direction  $y_N$ )



$$\begin{cases} -\operatorname{div}_y (A(y) \nabla_y w_i^{bl}(y)) = 0 & \text{in } G \\ w_i^{bl} = -w_i & \text{on } \Gamma \\ y' \rightarrow w_i^{bl}(y', y_N) & Y\text{-periodic} \end{cases}$$

Approximate the boundary layer (on just one side) by

$$u_\epsilon^{bl}(x) \approx \frac{\partial u}{\partial x_N}(x) w_N^{bl}\left(\frac{x}{\epsilon}\right)$$