# OPTIMAL DESIGN OF STRUCTURES (MAP 562) 

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TOPOLOGY OPTIMIZATION
BY THE HOMOGENIZATION METHOD

## Brief review of the preceding course

1. Topology optimization versus geometric optimization
2. Homogenization method in the periodic case (two-scale asymptotic expansions)
3. An explicit class of composite materials: sequential laminates.

What remains to be done:
To characterize the set $G_{\theta}$ of all composites materials
Towards this goal, prove bounds on $A^{*}$.
Application to shape optimization
To build numerical algorithms for topology optimization

From now on, we assume that the microscopic tensor $A(y)$ is symmetric.
Then $A^{*}$ is symmetric too.
Furthermore, $A^{*}$ is characterized by the variational principle

$$
A^{*} \xi \cdot \xi=\min _{w \in H_{\#}^{1}(Y) / \mathbf{R}} \int_{Y} A(y)(\xi+\nabla w) \cdot(\xi+\nabla w) d y
$$

Indeed, if $w_{\xi}$ is the minimizer, then it satisfies the Euler optimality condition

$$
\begin{cases}-\operatorname{div}\left(A(y)\left(\xi+\nabla w_{\xi}(y)\right)\right)=0 & \text { in } Y \\ y \rightarrow w_{\xi}(y) & Y \text {-periodic. }\end{cases}
$$

By linearity, we have $w_{\xi}=\sum_{i=1}^{N} \xi_{i} w_{i}$ and thus

$$
\int_{Y} A(y)\left(\xi+\nabla w_{\xi}\right) \cdot\left(\xi+\nabla w_{\xi}\right) d y=\sum_{i, j=1}^{N} \xi_{i} \xi_{j} A_{i j}^{*}=A^{*} \xi \cdot \xi
$$

## Arithmetic and harmonic mean bounds

Taking $w=0$ in the variational principle, we deduce the arithmetic mean bound

$$
A^{*} \xi \cdot \xi \leq\left(\int_{Y} A(y) d y\right) \xi \cdot \xi
$$

Enlarging the minimization space, we obtain the harmonic mean bound

$$
\left(\int_{Y} A^{-1}(y) d y\right)^{-1} \xi \cdot \xi \leq A^{*} \xi \cdot \xi
$$

These bounds can be improved for two-phase composites !

Indeed, since $\int_{Y} \nabla w d y=0$, we enlarge the minimization space by replacing $\nabla w$ with any vector field $\zeta(y)$ with zero-average on $Y$

$$
A^{*} \xi \cdot \xi \geq \min _{\zeta \in L_{\#}^{2}(Y)^{N}, \int_{Y} \zeta d y=0} \int_{Y} A(y)(\xi+\zeta(y)) \cdot(\xi+\zeta(y)) d y
$$

The Euler equation for the minimizer $\zeta_{\xi}(y)$ of this convex problem is

$$
A(y)\left(\xi+\zeta_{\xi}(y)\right)=\lambda
$$

where $\lambda \in \mathbb{R}$ is the Lagrange multiplier for the constraint $\int_{Y} \zeta d y=0$. We deduce

$$
\xi=\left(\int_{Y} A(y)^{-1} d y\right) \lambda
$$

and thus

$$
\int_{Y} A(y)\left(\xi+\zeta_{\xi}(y)\right) \cdot\left(\xi+\zeta_{\xi}(y)\right) d y=\left(\int_{Y} A(y)^{-1} d y\right)^{-1} \xi \cdot \xi
$$

### 7.3.5 Characterization of $G_{\theta}$

We consider two isotropic phases $A=\alpha \operatorname{Id}$ and $B=\beta$ Id with $0<\alpha<\beta$.
Theorem 7.17. The set $G_{\theta}$ of all homogenized tensors obtained by mixing $\alpha$ and $\beta$ in proportions $\theta$ and $(1-\theta)$ is the set of all symmetric matrices $A^{*}$ with eigenvalues $\lambda_{1}, \ldots, \lambda_{N}$ such that

$$
\begin{aligned}
\left(\frac{\theta}{\alpha}+\frac{1-\theta}{\beta}\right)^{-1} & =\lambda_{\theta}^{-} \leq \lambda_{i} \leq \lambda_{\theta}^{+}=\theta \alpha+(1-\theta) \beta \quad 1 \leq i \leq N \\
& \sum_{i=1}^{N} \frac{1}{\lambda_{i}-\alpha} \leq \frac{1}{\lambda_{\theta}^{-}-\alpha}+\frac{N-1}{\lambda_{\theta}^{+}-\alpha} \\
& \sum_{i=1}^{N} \frac{1}{\beta-\lambda_{i}} \leq \frac{1}{\beta-\lambda_{\theta}^{-}}+\frac{N-1}{\beta-\lambda_{\theta}^{+}}
\end{aligned}
$$

Furthermore, these so-called Hashin and Shtrikman bounds are optimal and attained by rank- $N$ sequential laminates.

Set $G_{\theta}$ in dimension $N=2$


Set $G_{\theta}$ in dimension $N=3$


Proof. We first show that all matrices satisfying these inequalities (Hashin-Shtrikman bounds) belong to $G_{\theta}$.

Let us start by showing that the upper bound is attained by sequential laminates. Take a matrix $A^{*}$ such that

$$
\sum_{i=1}^{N} \frac{1}{\beta-\lambda_{i}}=\frac{1}{\beta-\lambda_{\theta}^{-}}+\frac{N-1}{\beta-\lambda_{\theta}^{+}}
$$

Define a rank- $N$ sequentiel laminate $A_{L}^{*}$ of matrix $\beta$ and inclusion $\alpha$, with lamination directions being the (orthogonal) eigenvectors of $A^{*}$
$\theta\left(A_{L}^{*}-\beta \mathrm{Id}\right)^{-1}=\frac{1}{\alpha-\beta} \operatorname{Id}+(1-\theta) \sum_{i=1}^{N} m_{i} \frac{e_{i} \otimes e_{i}}{\beta} \quad$ with $\quad m_{i} \geq 0, \sum_{i=1}^{N} m_{i}=1$.
We have $A^{*}=A_{L}^{*}$ if we can choose the $m_{i}$ 's such that

$$
\frac{\theta}{\lambda_{i}-\beta}=\frac{1}{\alpha-\beta}+\frac{m_{i}(1-\theta)}{\beta} \Leftrightarrow m_{i}=\frac{\beta\left(\lambda_{\theta}^{+}-\lambda_{i}\right)}{(1-\theta)(\beta-\alpha)\left(\beta-\lambda_{i}\right)}
$$

We check that $0<m_{i}<1$ is equivalent to $\lambda_{\theta}^{-}<\lambda_{i}<\lambda_{\theta}^{+}$and that

$$
\sum_{i=1}^{N} m_{i}=1 \Leftrightarrow \sum_{i=1}^{N} \frac{1}{\beta-\lambda_{i}}=\frac{1}{\beta-\lambda_{\theta}^{-}}+\frac{N-1}{\beta-\lambda_{\theta}^{+}}
$$

thus any matrix on the upper bound is a rank- $N$ sequential laminate with matrix $\beta$ and inclusion $\alpha$.

The same proof works for the lower bound upon exchanging the role of $\alpha$ (now the matrix) and $\beta$ (now the inclusions).

Then, the next easy computation shows that the matrices "inside" $G_{\theta}$ are attained by simple lamination of two matrices, one on the upper bound, the other on the lower bound.

## Computation for the interior of $G_{\theta}$

Recall the lamination formula:

$$
\tau\left(A^{*}-B\right)^{-1}=(A-B)^{-1}+\frac{(1-\tau)}{B e_{1} \cdot e_{1}} e_{1} \otimes e_{1}
$$

Particular case: $A, B \in G_{\theta}$ diagonal in the same basis $\left(e_{1}, \ldots, e_{N}\right)$.

$$
A=\operatorname{diag}\left(a_{1}, \ldots, a_{N}\right) \quad B=\operatorname{diag}\left(b_{1}, \ldots, b_{N}\right)
$$

Then, for any $\tau \in[0,1], A^{*} \in G_{\theta}$ and

$$
a_{1}^{*}=\left(\frac{\tau}{a_{1}}+\frac{1-\tau}{b_{1}}\right)^{-1} \quad a_{i}^{*}=\tau a_{i}+(1-\tau) b_{i} \quad 2 \leq i \leq N
$$

Branches of hyperbolas which connect the upper and lower bounds of $G_{\theta}$.

It remains to prove that the lower and upper Hashin-Shtrikman bounds hold true.

To establish the lower bound we introduce the so-called Hashin and Shtrikman variational principle.

Main idea: use Fourier analysis and Plancherel theorem, but, in a first step, eliminate the cubic terms.

By definition of $A^{*}$, for $\xi \in \mathbb{R}^{N}$, we have

$$
A^{*} \xi \cdot \xi=\min _{w(y) \in H_{\#}^{1}(Y)} \int_{Y}(\chi(y) \alpha+(1-\chi(y)) \beta)(\xi+\nabla w) \cdot(\xi+\nabla w) d y
$$

Substracting a reference material $\alpha$

$$
\begin{gathered}
\int_{Y}(\chi \alpha+(1-\chi) \beta)|\xi+\nabla w|^{2} d y= \\
\int_{Y}(1-\chi)(\beta-\alpha)|\xi+\nabla w|^{2} d y+\int_{Y} \alpha|\xi+\nabla w|^{2} d y
\end{gathered}
$$

We use convex duality (or Legendre transform): for any symmetric positive definite matrix $M$

$$
M \zeta \cdot \zeta=\max _{\eta \in \mathbf{R}^{N}}\left(2 \zeta \cdot \eta-M^{-1} \eta \cdot \eta\right) \quad \forall \zeta \in \mathbb{R}^{N}
$$

Since $\beta-\alpha>0$, we apply the above formula at each point in $Y$, and we get

$$
\begin{aligned}
& \quad \int_{Y}(1-\chi)(\beta-\alpha)|\xi+\nabla w|^{2} d y= \\
& \max _{\eta(y) \in L_{\#}^{2}(Y)^{N}} \int_{Y}(1-\chi)\left(2(\xi+\nabla w) \cdot \eta-(\beta-\alpha)^{-1}|\eta|^{2}\right) d y,
\end{aligned}
$$

which becomes an inequality if we restrict the minimization to constant $\eta$ in $Y$

$$
\begin{aligned}
& \int_{Y}(1-\chi)(\beta-\alpha)|\xi+\nabla w|^{2} d y \geq \\
\geq & \max _{\eta} \int_{Y}(1-\chi)\left(2(\xi+\nabla w) \cdot \eta-(\beta-\alpha)^{-1}|\eta|^{2}\right) d y \\
\geq & (1-\theta)\left(2 \xi \cdot \eta-(\beta-\alpha)^{-1}|\eta|^{2}\right)-2 \int_{Y} \chi \nabla w \cdot \eta d y .
\end{aligned}
$$

On the other hand, because of periodicity, $\int_{Y} \nabla w d y=0$ which implies

$$
\int_{Y} \alpha|\xi+\nabla w|^{2} d y=\alpha|\xi|^{2}+\int_{Y} \alpha|\nabla w|^{2} d y
$$

Overall, we obtain, for any $\eta \in \mathbb{R}^{N}$,

$$
A^{*} \xi \cdot \xi \geq \alpha|\xi|^{2}+(1-\theta)\left(2 \xi \cdot \eta-(\beta-\alpha)^{-1}|\eta|^{2}\right)-g(\chi, \eta)
$$

where $g(\chi, \eta)$ is a so-called non-local term, defined by

$$
g(\chi, \eta)=-\min _{w(y) \in H_{\#}^{1}(Y)} \int_{Y}\left(\alpha|\nabla w|^{2}-2 \chi \nabla w \cdot \eta\right) d y
$$

We can now use Fourier analysis to compute $g(\chi, \eta)$.

By periodicity, $\chi$ and the test function $w$ can be written as Fourier series

$$
\chi(y)=\sum_{k \in \mathbb{Z}^{N}} \hat{\chi}(k) e^{2 i \pi k \cdot y}, \quad w(y)=\sum_{k \in \mathbb{Z}^{N}} \hat{w}(k) e^{2 i \pi k \cdot y}
$$

Since $\chi$ and $w$ are real-valued, their Fourier coefficients satisfy

$$
\overline{\hat{\chi}(k)}=\hat{\chi}(-k) \quad \text { and } \quad \overline{\hat{w}(k)}=\hat{w}(-k) .
$$

The gradient of $w$ is

$$
\nabla w(y)=\sum_{k \in \mathbb{Z}^{N}} 2 i \pi e^{2 i \pi k \cdot y} \hat{w}(k) k
$$

Plancherel formula yields

$$
\begin{gathered}
\int_{Y}\left(\alpha|\nabla w|^{2}-2 \chi \nabla w \cdot \eta\right) d y \\
=\sum_{k \in \mathbb{Z}^{N}}\left(4 \pi^{2} \alpha|\hat{w}(k) k|^{2}-4 i \pi \overline{\hat{\chi}(k)} \hat{w}(k) k \cdot \eta\right) \\
=\sum_{k \in \mathbb{Z}^{N}}\left(4 \pi^{2} \alpha|k|^{2}|\hat{w}(k)|^{2}+4 \pi \mathcal{I} m(\overline{\hat{\chi}(k)} \hat{w}(k)) \eta \cdot k\right) .
\end{gathered}
$$

To minimiz in $w(y) \in H_{\#}^{1}(Y) \Leftrightarrow$ to minimize in $\hat{w}(k) \in \mathbb{C}$.
For $k \neq 0$ the minimum is achieved by

$$
\hat{w}(k)=-\frac{i \hat{\chi}(k)}{2 \pi \alpha|k|^{2}} \eta \cdot k
$$

and we deduce

$$
g(\chi, \eta)=\left(\alpha^{-1} \sum_{k \in \mathbb{Z}^{N}, k \neq 0}|\hat{\chi}(k)|^{2} \frac{k}{|k|} \otimes \frac{k}{|k|}\right) \eta \cdot \eta=\alpha^{-1} \theta(1-\theta) M \eta \cdot \eta,
$$

where $M$ is a symmetric non-negative matrix. Since, by Plancherel theorem, we have

$$
\sum_{k \in \mathbb{Z}^{N}, k \neq 0}|\hat{\chi}(k)|^{2}=\int_{Y}|\chi(y)-\theta|^{2} d y=\theta(1-\theta)
$$

we deduce that the trace of $M$ is equal to 1 .

Regrouping terms yields, for any $\xi, \eta \in \mathbb{R}^{N}$,

$$
A^{*} \xi \cdot \xi \geq \alpha|\xi|^{2}+(1-\theta)\left(2 \xi \cdot \eta-(\beta-\alpha)^{-1}|\eta|^{2}\right)-\alpha^{-1} \theta(1-\theta) M \eta \cdot \eta .
$$

The minimum (in $\xi$ ) of this inequality is obtained when

$$
\xi=(1-\theta)\left(A^{*}-\alpha\right)^{-1} \eta
$$

We deduce

$$
\begin{gathered}
(1-\theta)\left(A^{*}-\alpha\right)^{-1} \eta \cdot \eta \leq(\beta-\alpha)^{-1}|\eta|^{2}+\alpha^{-1} \theta M \eta \cdot \eta \quad \forall \eta \in \mathbb{R}^{N} . \\
\Leftrightarrow(1-\theta)\left(A^{*}-\alpha\right)^{-1} \leq(\beta-\alpha)^{-1} \mathrm{Id}+\alpha^{-1} \theta M
\end{gathered}
$$

Taking the trace of this matrix inequality, and since $\operatorname{Tr} M=1$, we obtain the lower Hashin-Shtrikman bound.

The proof of the upper bound is similar.

### 7.4 Homogenized formulation of shape optimization

The relaxed or homogenized optimization problem is

$$
\min _{\left(\theta, A^{*}\right) \in \mathcal{U}_{a d}^{*}} J\left(\theta, A^{*}\right)
$$

with an objective function

$$
J\left(\theta, A^{*}\right)=\int_{\Omega} f u d x, \quad \text { or } \quad J\left(\theta, A^{*}\right)=\int_{\Omega}\left|u-u_{0}\right|^{2} d x
$$

and an homogenized admissible set given by
$\mathcal{U}_{a d}^{*}=\left\{\left(\theta, A^{*}\right) \in L^{\infty}\left(\Omega ;[0,1] \times \mathbb{R}^{N^{2}}\right), A^{*}(x) \in G_{\theta(x)}\right.$ in $\left.\Omega, \int_{\Omega} \theta(x) d x=V_{\alpha}\right\}$,
where $G_{\theta}$ is explicitly characterized.
The homogenized state equation is

$$
\begin{cases}-\operatorname{div}\left(A^{*} \nabla u\right)=f & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

Theorem 7.19 (admitted). The homogenized formulation is actually a relaxation of the original shape optimization problem in the sense that:
there exists, at least, one optimal composite shape $\left(\theta, A^{*}\right)$,
any minimizing sequence of classical shapes $\chi$ converges, in the sense of homogenization, to a composite optimal solution $\left(\theta, A^{*}\right)$,
any composite optimal solution $\left(\theta, A^{*}\right)$ is the limit of a minimizing sequence of classical shapes.
The minima of the original and homogenized objective functions coincide

$$
\inf _{\chi \in \mathcal{U}_{a d}} J(\chi)=\min _{\left(\theta, A^{*}\right) \in \mathcal{U}_{a d}^{*}} J\left(\theta, A^{*}\right)
$$

## Remark.

The shape optimization problem is thus not changed by relaxation.
Close to any optimal composite shape, we are sure to find a quasi-optimal classical shape.

This theorem is at the root of new numerical algorithms.

### 7.4.2 Optimality conditions

We now compute the gradient of the following objective function

$$
J\left(\theta, A^{*}\right)=\int_{\Omega}\left|u-u_{0}\right|^{2} d x
$$

where $u_{0} \in L^{2}(\Omega)$. We introduce the adjoint state $p$, unique solution in $H_{0}^{1}(\Omega)$ of

$$
\begin{cases}-\operatorname{div}\left(A^{*} \nabla p\right)=-2\left(u-u_{0}\right) & \text { in } \Omega \\ p=0 & \text { on } \partial \Omega\end{cases}
$$

Proposition 7.20. Let $\alpha>0$ and $\mathcal{M}_{\alpha}$ be the set of symmetric positive definite matrices $M$ such that $M \geq \alpha$ Id. The functional $J$ is differentiable with respect to $A^{*}$ in $L^{\infty}\left(\Omega ; \mathcal{M}_{\alpha}\right)$, and its derivative is

$$
\nabla_{A^{*}} J\left(\theta, A^{*}\right)=\nabla u \otimes \nabla p
$$

Remark. The partial derivative with respect to $\theta$ vanishes because $\theta$ appears only in the constraint $A^{*} \in G_{\theta}$.

## Proof of Proposition 7.20

It is standard! It became a parametric (sizing) shape optimization problem where $A^{*}$ plays the role of a thickness.

We introduce the Lagrangian

$$
\mathcal{L}\left(A^{*}, v, q\right)=\int_{\Omega}\left|v-u_{0}\right|^{2} d x+\int_{\Omega} A^{*} \nabla v \cdot \nabla q d x-\int_{\Omega} f q d x
$$

Its partial derivative with respect to $q$ yields the state.
Its partial derivative with respect to $v$ yields the adjoint.
Its partial derivative with respect to $A^{*}$ yields the gradient

$$
\nabla_{A^{*}} J\left(\theta, A^{*}\right)=\frac{\partial \mathcal{L}}{\partial A^{*}}\left(A^{*}, u, p\right)=\nabla u \otimes \nabla p
$$

## Essential consequence

Theorem 7.21. Let $\left(\theta, A^{*}\right)$ be a global minimizer of $J$ in $\mathcal{U}_{a d}^{*}$ which admits $u$ and $p$ as state and adjoint. There exists $\left(\tilde{\theta}, \tilde{A}^{*}\right)$, another global minimizer of $J$ in $\mathcal{U}_{a d}^{*}$, which admits the same state and adjoint $u$ and $p$, and such that $\tilde{A}^{*}$ is a rank-1 simple laminate.

Simplification: in the definition of $\mathcal{U}_{a d}^{*}$ the set $G_{\theta}$ can be replaced by its simpler subset of rank-1 simple laminates.

## Remark.

Optimality condition $\Rightarrow$ simplification of the problem.
We actually use this simplification in the numerical algorithms.
Simplification which holds true for other objective functions, but not for multiple loads optimization.

Proof. We fix $\theta$ and makes variations on $A^{*}$ only. Remarking that $G_{\theta}$ is convex (not obvious), the optimality condition is an Euler inequality which is

$$
\int_{\Omega}\left(A^{0}-A^{*}\right) \nabla u \cdot \nabla p d x \geq 0
$$

for any $A^{0} \in G_{\theta}$, which is equivalent to

$$
A^{*} \nabla u \cdot \nabla p=\min _{A^{0} \in G_{\theta}}\left(A^{0} \nabla u \cdot \nabla p\right) \quad \forall x \in \Omega .
$$

If $\nabla u$ or $\nabla p$ vanishes, then any $A^{*}$ is optimal. Otherwise, we define

$$
e=\frac{\nabla u}{|\nabla u|} \quad \text { and } \quad e^{\prime}=\frac{\nabla p}{|\nabla p|},
$$

and we look for minimizers $A^{*}$ of

$$
\min _{A^{0} \in G_{\theta}} 4 A^{0} e \cdot e^{\prime}=A^{0}\left(e+e^{\prime}\right) \cdot\left(e+e^{\prime}\right)-A^{0}\left(e-e^{\prime}\right) \cdot\left(e-e^{\prime}\right)
$$

A lower bound is easily obtained

$$
\begin{aligned}
\min _{A^{0} \in G_{\theta}} 4 A^{0} e \cdot e^{\prime} & \geq \min _{A^{0} \in G_{\theta}} A^{0}\left(e+e^{\prime}\right) \cdot\left(e+e^{\prime}\right)-\max _{A^{0} \in G_{\theta}} A^{0}\left(e-e^{\prime}\right) \cdot\left(e-e^{\prime}\right) \\
& =\lambda_{\theta}^{-}\left|e+e^{\prime}\right|^{2}-\lambda_{\theta}^{+}\left|e-e^{\prime}\right|^{2}
\end{aligned}
$$

This lower bound is actually the precise minimal value.
Indeed, choosing $A^{0}=A^{1}$ which is a rank-1 simple laminate in the direction $e+e^{\prime}$, orthogonal to $e-e^{\prime}$, we get

$$
A^{1}\left(e+e^{\prime}\right)=\lambda_{\theta}^{-}\left(e+e^{\prime}\right) \quad \text { and } \quad A^{1}\left(e-e^{\prime}\right)=\lambda_{\theta}^{+}\left(e-e^{\prime}\right)
$$

and an easy computation shows that

$$
4 A^{1} e \cdot e^{\prime}=\lambda_{\theta}^{-}\left|e+e^{\prime}\right|^{2}-\lambda_{\theta}^{+}\left|e-e^{\prime}\right|^{2}
$$

Thus

$$
\min _{A^{0} \in G_{\theta}} 4 A^{0} e \cdot e^{\prime}=\lambda_{\theta}^{-}\left|e+e^{\prime}\right|^{2}-\lambda_{\theta}^{+}\left|e-e^{\prime}\right|^{2}
$$

If now $A^{*}$ is any optimal tensor, then, as a rank- 1 laminate, it satisfies

$$
\begin{equation*}
A^{*}\left(e+e^{\prime}\right)=\lambda_{\theta}^{-}\left(e+e^{\prime}\right) \quad \text { and } \quad A^{*}\left(e-e^{\prime}\right)=\lambda_{\theta}^{+}\left(e-e^{\prime}\right) \tag{1}
\end{equation*}
$$

Indeed, if (1) does not hold true, one of the arithmetic and harmonic bounds would give a strict inequality

$$
4 A^{*} e \cdot e^{\prime}=A^{*}\left(e+e^{\prime}\right) \cdot\left(e+e^{\prime}\right)-A^{*}\left(e-e^{\prime}\right) \cdot\left(e-e^{\prime}\right)>\lambda_{\theta}^{-}\left|e+e^{\prime}\right|^{2}-\lambda_{\theta}^{+}\left|e-e^{\prime}\right|^{2}
$$

which is a contradiction with the optimal character of $A^{*}$.

We deduce that any optimal $A^{*}$ satisfies, like the rank-1 simple laminate $A^{1}$,

$$
\begin{aligned}
& 2 A^{*} \nabla u=2 A^{1} \nabla u=\left(\lambda_{\theta}^{+}+\lambda_{\theta}^{-}\right) \nabla u+\left(\lambda_{\theta}^{+}-\lambda_{\theta}^{-}\right) \frac{|\nabla u|}{|\nabla p|} \nabla p \\
& 2 A^{*} \nabla p=2 A^{1} \nabla p=\left(\lambda_{\theta}^{+}+\lambda_{\theta}^{-}\right) \nabla p+\left(\lambda_{\theta}^{+}-\lambda_{\theta}^{-}\right) \frac{|\nabla p|}{|\nabla u|} \nabla u,
\end{aligned}
$$

Therefore any optimal tensor $A^{*}$ can be replaced by this rank- 1 simple laminate $A^{1}$ without changing $u$ and $p$.

$$
\begin{gathered}
-\operatorname{div}\left(A^{*} \nabla u\right)=-\operatorname{div}\left(A^{1} \nabla u\right)=f \\
-\operatorname{div}\left(A^{*} \nabla p\right)=-\operatorname{div}\left(A^{1} \nabla p\right)=-2\left(u-u_{0}\right)
\end{gathered}
$$

## Parametrization of rank-1 simple laminates

In space dimension $N=2$ (to simplify) a rank-1 laminate is defined by

$$
A^{*}(\theta, \phi)=\left(\begin{array}{rr}
\cos \phi & \sin \phi \\
-\sin \phi & \cos \phi
\end{array}\right)\left(\begin{array}{ll}
\lambda_{\theta}^{+} & 0 \\
0 & \lambda_{\theta}^{-}
\end{array}\right)\left(\begin{array}{rr}
\cos \phi & -\sin \phi \\
\sin \phi & \cos \phi
\end{array}\right) \quad \phi \in[0, \pi]
$$

The admissible set is thus simply

$$
\mathcal{U}_{a d}^{L}=\left\{(\theta, \phi) \in L^{\infty}(\Omega ;[0,1] \times[0, \pi]), \int_{\Omega} \theta(x) d x=V_{\alpha}\right\}
$$

Proposition 7.23. The objective function $J(\theta, \phi)$ is differentiable with respect to $(\theta, \phi)$ in $\mathcal{U}_{a d}^{L}$, and its derivative is

$$
\nabla_{\phi} J(\theta, \phi)=\frac{\partial A^{*}}{\partial \phi} \nabla u \cdot \nabla p \quad \text { and } \quad \nabla_{\theta} J(\theta, \phi)=\frac{\partial A^{*}}{\partial \theta} \nabla u \cdot \nabla p
$$

### 7.4.3 Numerical algorithm

Projected gradient algorithm for the minimization of $J(\theta, \phi)$.

1. We initialize the design parameters $\theta_{0}$ and $\phi_{0}$ (for example, equal to constants).
2. Until convergence, for $k \geq 0$ we iterate by computing the state $u_{k}$ and adjoint $p_{k}$, solutions with the previous design parameters $\left(\theta_{k}, \phi_{k}\right)$, then we update these parameters by

$$
\begin{aligned}
\theta_{k+1} & =\max \left(0, \min \left(1, \theta_{k}-t_{k}\left(\ell_{k}+\frac{\partial A^{*}}{\partial \theta}\left(\theta_{k}, \phi_{k}\right) \nabla u_{k} \cdot \nabla p_{k}\right)\right)\right) \\
\phi_{k+1} & =\phi_{k}-t_{k} \frac{\partial A^{*}}{\partial \phi}\left(\theta_{k}, \phi_{k}\right) \nabla u_{k} \cdot \nabla p_{k}
\end{aligned}
$$

with $\ell_{k}$ a Lagrange multiplier for the volume constraint (iteratively enforced), and $t_{k}>0$ a descent step such that $J\left(\theta_{k+1}, \phi_{k+1}\right)<J\left(\theta_{k}, \phi_{k}\right)$.

## The self-adjoint case

A first example: maximization of torsional rigidity (maximization of compliance).

$$
\min _{\left(\theta, A^{*}\right) \in \mathcal{U}_{a d}^{L}}\left\{J\left(\theta, A^{*}\right)=-\int_{\Omega} u(x) d x\right\},
$$

where $u$ is the solution of

$$
\begin{cases}-\operatorname{div}\left(A^{*} \nabla u\right)=1 & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

and the adjoint state is just $p=u$.
We solve in the domain $\Omega=(0,1)^{2}$ with the phases $\alpha=1$ and $\beta=2$. We fix a $50 \%$ volume constraint of $\alpha$. We initialize with a constant value of $\theta=0.5$ and a constant zero lamination angle. We perform 30 iterations.

Self-adjoint case $p=u$.

$$
\nabla_{A^{*}} J\left(\theta, A^{*}\right)=\nabla u \otimes \nabla u \geq 0 .
$$

To minimize $J$ we have to decrease $A^{*}$.
Any optimal $A^{*}$ satisfies

$$
A^{*} \nabla u=\lambda_{\theta}^{-} \nabla u
$$

thus the optimal composite is the worst possible conductor.
Consequence. We can eliminate the angle $\phi$ and it remains to optimize with respect to $\theta$ only !

## Convexity

We rewrite the optimization problem thanks to the primal energy

$$
-\int_{\Omega} u d x=-\int_{\Omega} \lambda_{\theta}^{-}|\nabla u|^{2} d x=\min _{v \in H_{0}^{1}(\Omega)} \int_{\Omega} \lambda_{\theta}^{-}|\nabla v|^{2} d x-2 \int_{\Omega} v d x
$$

Thus, we obtain a double minimization

$$
\min _{\theta, A^{*}=\lambda_{\theta}^{-}} J\left(\theta, A^{*}\right)=\min _{\theta, v} \int_{\Omega} \lambda_{\theta}^{-}|\nabla v|^{2} d x-2 \int_{\Omega} v d x
$$

Remember: the function $(\theta, v) \rightarrow \lambda_{\theta}^{-}|\nabla v|^{2}$ is convex.
Consequence. There are only global minima!
Numerically, we use an algorithm based on alternate direction minimization (see chapter 5).

Convergence history:
objective function in terms of the iteration number.


## Volume fraction $\theta$ (iterations 1,5 , and 30)







## A second self-adjoint example

Compliance minimization.

$$
\min _{\left(\theta, A^{*}\right) \in \mathcal{U}_{a d}^{L}}\left\{J\left(\theta, A^{*}\right)=\int_{\Omega} u(x) d x\right\}
$$

where $u$ is the solution of

$$
\begin{cases}-\operatorname{div}\left(A^{*} \nabla u\right)=1 & \text { in } \Omega \\ u=0 & \text { on } \partial \Omega\end{cases}
$$

and the adjoint state is just $p=-u$.
We solve in the domain $\Omega=(0,1)^{2}$ with the phases $\alpha=1$ and $\beta=2$. We fix a $50 \%$ volume constraint of $\alpha$. We initialize with a constant value of $\theta=0.5$ and a constant zero lamination angle. We perform 30 iterations.

Self-adjoint case $p=-u$.

$$
\nabla_{A^{*}} J\left(\theta, A^{*}\right)=-\nabla u \otimes \nabla u \leq 0
$$

To minimize $J$ we have to increase $A^{*}$.
Any optimal $A^{*}$ satisfies

$$
A^{*} \nabla u=\lambda_{\theta}^{+} \nabla u
$$

thus the optimal composite is the best possible conductor.
Consequence. We can eliminate the angle $\phi$ and it remains to optimize with respect to $\theta$ only !

## Convexity

We rewrite the optimization problem thanks to the dual energy

$$
\int_{\Omega} u d x=\min _{\substack{\tau \in L^{2}(\Omega) \\-d_{i v}^{N} \\ \text { div } \tau=\text { in } \Omega}} \int_{\Omega}\left(\lambda_{\theta}^{+}\right)^{-1}|\tau|^{2} d x .
$$

Thus, we obtain a double minimization

$$
\min _{\theta, A^{*}=\lambda_{\theta}^{+}} J\left(\theta, A^{*}\right)=\min _{\theta, \tau} \int_{\Omega}\left(\lambda_{\theta}^{+}\right)^{-1}|\tau|^{2} d x
$$

Remember: the function $(\theta, \tau) \rightarrow \frac{|\tau|^{2}}{\lambda_{\theta}^{+}}$is convex.
Consequence. There are only global minima!
Numerically, we use an algorithm based on alternate direction minimization (see chapter 5).

Minimal compliance membrane (iterations 1, 10, and 30)


## Remarks

Convergence to a global minimum.

1. Numerical experiments with various initializations.
2. Convexity properties.

Shape optimization rather than two-phase optimization.

1. Numerically, holes can be mimicked by a very weak phase $\alpha\left(\approx 10^{-3} \beta\right)$.
2. Mathematically, when $\alpha \rightarrow 0$ we obtain Neumann boundary conditions on the holes boundaries.

## Penalization

The previous algorithm compute composite shapes while we are rather interested by classical shapes.

Therefore we use a penalization process to force the density to take values close to 0 or 1 .

Possible algorithms: after convergence to a composite shape,

1. either we add a penalization term to the objective function

$$
J\left(\theta, A^{*}\right)+c_{p e n} \int_{\Omega} \theta(1-\theta) d x
$$

2. either we continue the previous algorithm with a modified "penalized" density

$$
\theta_{p e n}=\frac{1-\cos \left(\pi \theta_{o p t}\right)}{2}
$$

If $0<\theta_{o p t}<1 / 2$, then $\theta_{\text {pen }}<\theta_{o p t}$, while, if $1 / 2<\theta_{o p t}<1$, then $\theta_{\text {pen }}>\theta_{\text {opt }}$.

## Example

Optimal radiator.

$$
\begin{cases}-\operatorname{div}\left(A^{*} \nabla u\right)=0 & \text { in } \Omega \\ A^{*} \nabla u \cdot n=1 & \text { on } \Gamma_{N} \\ A^{*} \nabla u \cdot n=0 & \text { on } \Gamma \\ u=0 & \text { on } \Gamma_{D}\end{cases}
$$

We minimize the temperature where heating takes place

$$
\min _{\left(\theta, A^{*}\right) \in \mathcal{U}_{a d}^{L}}\left\{J\left(\theta, A^{*}\right)=\int_{\Gamma_{N}} u d s\right\} .
$$

This is precisely the compliance! Thus the problem is self-adjoint with $p=-u$.

Isotropic materials with conductivity $\alpha=0.01$ and $\beta=1$, in proportions $50,50 \%$, in the domain $\Omega=(0,1)^{2}$.

## Optimal radiator



