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A spectral method to detect invisibility in waveguides

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General setting

• We are interested in the propagation of waves in acoustic waveguides.



• In this talk, we study questions of invisibility.

Can we find situations where waves go through like if there were no defect

• One can wish to have good energy transmission through the structure.

• One can wish to hide objects.

Scattering in time-harmonic regime of an incident wave in the acoustic waveguide Ω coinciding with $\{(x, y) \in \mathbb{R} \times (0; 1)\}$ outside a compact region.



Find $v = v_i + v_s$ s. t. $\Delta v + k^2 v = 0$ in Ω , $\partial_n v = 0$ on $\partial \Omega$, v_s is outgoing.

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 v_s is outgoing.

For this problem with $k \in ((N-1)\pi; N\pi)$, $N \in \mathbb{N}^*$, the modes are Propagating $\begin{vmatrix} w_n^{\pm}(x,y) = e^{\pm i\beta_n x} \cos(n\pi y), \ \beta_n = \sqrt{k^2 - n^2 \pi^2}, \ n \in \llbracket 0, N-1 \rrbracket$ Evanescent $\begin{vmatrix} w_n^{\pm}(x,y) = e^{\mp \beta_n x} \cos(n\pi y), \ \beta_n = \sqrt{n^2 \pi^2 - k^2}, \ n \ge N.$

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• Set $v_i = \sum_{n=0}^{N-1} \alpha_n w_n^+$ (propagating) for some given $(\alpha_n)_{n=0}^{N-1} \in \mathbb{C}^N$.

Scattering in time-harmonic regime of an incident wave in the acoustic waveguide Ω coinciding with $\{(x, y) \in \mathbb{R} \times (0; 1)\}$ outside a compact region.



$$\begin{array}{lll} \mbox{Find} \ v = v_i + v_s \ {\rm s.} \ {\rm t.} \\ \Delta v + k^2 v &= 0 & \mbox{in} \ \Omega, \\ \partial_n v &= 0 & \mbox{on} \ \partial \Omega, \\ v_s \ {\rm is \ outgoing.} \end{array}$$

For this problem with $k \in ((N-1)\pi; N\pi)$, $N \in \mathbb{N}^*$, the modes are

 $\begin{array}{l} \text{Propagating} & \left| \begin{array}{l} w_n^{\pm}(x,y) = e^{\pm i\beta_n x} \cos(n\pi y), \ \beta_n = \sqrt{k^2 - n^2 \pi^2}, \ n \in \llbracket 0, N - 1 \rrbracket \\ \text{Evanescent} & \left| \begin{array}{l} w_n^{\pm}(x,y) = e^{\mp \beta_n x} \cos(n\pi y), \ \beta_n = \sqrt{n^2 \pi^2 - k^2}, \ n \ge N. \end{array} \right. \end{array}$

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$$v_s \text{ is outgoing } \Leftrightarrow \quad v_s = \sum_{n=0}^{+\infty} \gamma_n^{\pm} w_n^{\pm} \quad \text{ for } \pm x \ge L, \text{ with } (\gamma_n^{\pm}) \in \mathbb{C}^{\mathbb{N}}.$$

Goal of the talk

DEFINITION: v is a non reflecting mode if v_s is expo. decaying for $x \leq -L$ $\Leftrightarrow \quad \gamma_n^- = 0, \ n \in [\![0, N-1]\!] \quad \Leftrightarrow \quad \text{energy is completely transmitted.}$



For a given geometry, we present a method to find values of k such that there is a non reflecting mode v.

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GOAL

For a given geometry, we present a method to find values of k such that there is a non reflecting mode v.

 \rightarrow Note that non reflection occurs for **particular** v_i to be computed.

1 Introduction

Classical complex scaling

We recall how to use classical complex scaling to compute trapped modes and complex resonances.

3 Conjugated complex scaling

We explain how to use conjugated complex scaling to compute non reflecting modes.

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$$\mathcal{I}_{\theta}(x) = \begin{vmatrix} -L + (x+L) e^{i\theta} & \text{for } x \leq -L \\ x & \text{for } |x| < L \\ +L + (x-L) e^{i\theta} & \text{for } x \geq L. \end{vmatrix}$$



• For $\theta \in (0; \pi/2)$, consider the complex change of variables

$$\mathcal{I}_{\theta}(x) = \begin{vmatrix} -L + (x+L) e^{i\theta} & \text{for } x \leq -L \\ x & \text{for } |x| < L \\ +L + (x-L) e^{i\theta} & \text{for } x \geq L. \end{vmatrix}$$

• Set $v_{\theta} := v_s \circ (\mathcal{I}_{\theta}(x), y)$.



$$v_{\theta} = \sum_{n=0} \tilde{\gamma}_{n}^{\pm} e^{\pm i\beta_{n}x} \cos(n\pi y) + \sum_{n=N} \tilde{\gamma}_{n}^{\pm} e^{\pm\beta_{n}x} \cos(n\pi y), \quad \pm x \ge L \quad \tilde{\beta}_{n} = \beta_{n} e^{i\theta}$$

$$\text{Set } v_{\theta} := v_{s} \circ (\mathcal{I}_{\theta}(x), y).$$

$$1) v_{\theta} = v_{s} \text{ for } |x| < L.$$

$$2) w_{\theta} \text{ is arm decaying at infinity.}$$

2) v_{θ} is exp. decaying at infinity.



 \triangleright v_{θ} solves

(*)
$$\left| \begin{array}{c} \alpha_{\theta} \frac{\partial}{\partial x} \left(\alpha_{\theta} \frac{\partial v_{\theta}}{\partial x} \right) + \frac{\partial^2 v_{\theta}}{\partial y^2} + k^2 v_{\theta} = 0 \quad \text{in } \Omega \\ \partial_n v_{\theta} = -\partial_n v_i \quad \text{on } \partial\Omega. \end{array} \right.$$

2/2

 \triangleright v_{θ} solves

$$\mathbf{s}\left[(\ast) \middle| \begin{array}{c} \alpha_{\theta} \frac{\partial}{\partial x} \left(\alpha_{\theta} \frac{\partial v_{\theta}}{\partial x} \right) + \frac{\partial^{2} v_{\theta}}{\partial y^{2}} + k^{2} v_{\theta} = 0 \quad \text{in } \Omega \\ \partial_{n} v_{\theta} = -\partial_{n} v_{i} \quad \text{on } \partial\Omega. \end{array} \right.$$
$$\alpha_{\theta}(x) = 1 \text{ for } |x| < L \qquad \alpha_{\theta}(x) = e^{-i\theta} \text{ for } |x| \ge L$$

•
$$v_{\theta}$$
 solves $\left| \begin{array}{c} (*) \\ \alpha_{\theta} \frac{\partial}{\partial x} \left(\alpha_{\theta} \frac{\partial v_{\theta}}{\partial x} \right) + \frac{\partial^2 v_{\theta}}{\partial y^2} + k^2 v_{\theta} = 0 \quad \text{in } \Omega \\ \partial_n v_{\theta} = -\partial_n v_i \quad \text{on } \partial\Omega. \end{array} \right|$
 $\alpha_{\theta}(x) = 1 \text{ for } |x| < L \qquad \alpha_{\theta}(x) = e^{-i\theta} \text{ for } |x| \ge L$

• Numerically we solve (*) in the truncated domain



 \Rightarrow We obtain a good approximation of v_s for |x| < L.

• This is the method of Perfectly Matched Layers (PMLs).

Spectral analysis

• Define the operators A, A_{θ} of $L^{2}(\Omega)$ such that

$$Av = -\Delta v, \qquad A_{\theta}v = -\left(\alpha_{\theta}\frac{\partial}{\partial x}\left(\alpha_{\theta}\frac{\partial v}{\partial x}\right) + \frac{\partial^2 v}{\partial y^2}\right) \qquad + \partial_n v = 0 \text{ on } \partial\Omega.$$





•
$$A_{\theta}$$
 is not selfadjoint. $\sigma(A_{\theta}) \subset \{\rho e^{i\gamma}, \rho \ge 0, \gamma \in [-2\theta; 0]\}.$
• $\sigma_{\text{ess}}(A_{\theta}) = \bigcup_{n \in \mathbb{N}} \{n^2 \pi^2 + t e^{-2i\theta}, t \ge 0\}.$

• real eigenvalues of A_{θ} = real eigenvalues of A.







• Discretized spectrum of A_{θ} in k (not in k^2). We take $\theta = \pi/4$.



10 / 22

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10 / 22

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We explain how to use conjugated complex scaling to compute non reflecting modes.

• Usual complex scaling selects scattered fields which are

outgoing at $-\infty$ and outgoing at $+\infty$.

IMPORTANT REMARK: general v decompose as

$$v = v_i + \sum_{n=0}^{N-1} \gamma_n^- w_n^- + \sum_{n=N}^{+\infty} \gamma_n^- w_n^- \quad x \le -L, \quad v = \sum_{n=0}^{+\infty} \gamma_n^+ w_n^+ \quad x \ge L.$$

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IMPORTANT REMARK: **non reflecting** v decompose as

$$v = v_i + \sum_{n=0}^{\infty-1} w_n w_n^- + \sum_{n=N}^{+\infty} \gamma_n^- w_n^- \quad x \le -L, \quad v = \sum_{n=0}^{+\infty} \gamma_n^+ w_n^+ \quad x \ge L.$$

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IMPORTANT REMARK: **non reflecting** v decompose as

$$v = \sum_{n=0}^{N-1} \alpha_n w_n^+ + \sum_{n=N}^{+\infty} \gamma_n^- w_n^- \quad x \le -L, \quad v = \sum_{n=0}^{+\infty} \gamma_n^+ w_n^+ \quad x \ge L.$$

In other words, **non reflecting** v are

ingoing at $-\infty$ and outgoing at $+\infty$.

• Usual complex scaling selects scattered fields which are

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IMPORTANT REMARK: **non reflecting** v decompose as

$$v = \sum_{n=0}^{N-1} \alpha_n w_n^+ + \sum_{n=N}^{+\infty} \gamma_n^- w_n^- \quad x \le -L, \quad v = \sum_{n=0}^{+\infty} \gamma_n^+ w_n^+ \quad x \ge L.$$

 \blacktriangleright In other words, **non reflecting** v are

ingoing at $-\infty$ and outgoing at $+\infty$.



Let us **change the sign** of the complex scaling at $-\infty$!

$$\mathcal{J}_{\theta}(x) = \begin{vmatrix} -L + (x+L) & e^{-i\theta} & \text{for } x \leq -L \\ x & \text{for } |x| < L \\ +L + (x-L) & e^{+i\theta} & \text{for } x \geq L. \end{vmatrix}$$

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Set $u_{\theta} := v \circ (\mathcal{J}_{\theta}(x), y)$.

$$1) \quad u_{\theta} = v \text{ for } |x| < L.$$

$$2) \quad u_{\theta} \text{ is exp. decaying at infinity.}$$

$$-i\beta_{1}$$

$$-\theta$$

$$-i\beta_{1}$$

$$\theta_{2}$$

$$\theta_{3}$$

$$exp. \text{ growing exp. decaying } \theta_{3}$$
Modal exponents for v ($x \leq -L$)
Modal exponents for u_{θ} ($x \leq -L$)

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Set $u_{\theta} := v \circ (\mathcal{J}_{\theta}(x), y)$.
$$\begin{array}{c} 1) & u_{\theta} = v \text{ for } |x| < L. \\ 2) & u_{\theta} \text{ is exp. decaying at infinity.} \end{aligned}$$

$$\begin{array}{c} \bullet & -i\beta_{0} \\ \bullet & -i\beta_{1} \\ \theta_{2} & \theta_{3} \\ exp. \text{ growing} & exp. \text{ decaying} \\ \end{array}$$
Modal exponents for $v \ (x \leq -L) \end{array}$
Modal exponents for $u_{\theta} \ (x \leq -L)$

$$u_{\theta} \text{ solves} \left[(*) \middle| \begin{array}{c} \beta_{\theta} \frac{\partial}{\partial x} \left(\beta_{\theta} \frac{\partial u_{\theta}}{\partial x} \right) + \frac{\partial^{2} u_{\theta}}{\partial y^{2}} + k^{2} u_{\theta} = 0 & \text{in } \Omega \\ \partial_{n} u_{\theta} = 0 & \text{on } \partial \Omega. \\ \end{array} \right]$$

For $\theta \in (0; \pi/2)$, consider the complex change of variables

$$\mathcal{J}_{\theta}(x) = \begin{vmatrix} -L + (x+L) & e^{-i\theta} & \text{for } x \leq -L \\ x & \text{for } |x| < L \\ +L + (x-L) & e^{+i\theta} & \text{for } x \geq L. \end{vmatrix}$$

$$\text{Set } u_{\theta} := v \circ (\mathcal{J}_{\theta}(x), y) \text{.} \qquad 1) \quad u_{\theta} = v \text{ for } |x| < L. \\ 2) \quad u_{\theta} \text{ is exp. decaying at infinity.} \end{vmatrix}$$

$$\overset{\bullet^{-i\beta_0}}{\overset{\bullet^{-i\beta_1}}}{\overset{\bullet^{-i\beta_1}}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}}{\overset{\bullet^{-i\beta_1}}{\overset{\bullet^{-i\beta_1}}}}}}}}}}} \\ \mathcal{I}$$

12 / 22

Spectral analysis

• Define the operator B_{θ} of $L^2(\Omega)$ such that

$$B_{\theta}v = -\left(\beta_{\theta}\frac{\partial}{\partial x}\left(\beta_{\theta}\frac{\partial v}{\partial x}\right) + \frac{\partial^2 v}{\partial y^2}\right) \qquad + \partial_n v = 0 \text{ on } \partial\Omega.$$

B_θ is not selfadjoint. σ(B_θ) ⊂ {ρe^{iγ}, ρ ≥ 0, γ ∈ [-2θ; 2θ]}.
σ_{ess}(B_θ) = ∪_{n∈N}{n²π² + t e^{-2iθ}, t ≥ 0} ∪ {n²π² + t e^{2iθ}, t ≥ 0}.
real eigenvalues of B_θ = real eigenvalues of A+non reflecting k².





• It is not simple to prove that $\sigma(B_{\theta}) \setminus \sigma_{\text{ess}}(B_{\theta})$ is discrete.



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\rightarrow Not true in general!

 $e^{ikx} \circ \mathcal{J}_{\theta}$ is an eigenfunction for all $k \in \mathscr{R}$.



► It is not simple to prove that $\sigma(B_{\theta}) \setminus \sigma_{ess}(B_{\theta})$ is discrete. → Not true in general! $e^{ikx} \circ \mathcal{J}_{\theta}$ is an eigenfunction for all $k \in \mathscr{R}$.









• Define the operators \mathcal{P} (Parity), \mathcal{T} (Time reversal) such that

$$\mathcal{P}v(x,y) = v(-x,y)$$
 and $\mathcal{T}v(x,y) = \overline{v(x,y)}$.

PROP.: For symmetric $\Omega = \{(-x, y) | (x, y) \in \Omega\}, B_{\theta} \text{ is } \mathcal{PT} \text{ symmetric:}$

 $\mathcal{PT}B_{\theta}\mathcal{PT} = B_{\theta}.$

As a consequence, $\sigma(B_{\theta}) = \overline{\sigma(B_{\theta})}$.

 \Rightarrow If λ is an "isolated" eigenvalue located close to the real axis, then $\lambda \in \mathbb{R}$!

• Discretized spectrum in k (not in k^2). We take $\theta = \pi/4$.



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15 / 22

• We display the eigenmodes for the ten first real eigenvalues in the whole computational domain (including PMLs).



• Let us focus on the eigenmodes such that $0 < k < \pi$.



First trapped mode k = 1.2355...



Second trapped mode k = 2.3897...



First non reflecting mode k = 1.4513...



Second non reflecting mode k = 2.8896...

• To check our results, we compute $k \mapsto |R(k)|$ for $0 < k < \pi$.



First non reflecting mode k = 1.4513...

Second non reflecting mode k = 2.8896...

• To check our results, we compute $k \mapsto |R(k)|$ for $0 < k < \pi$.



• Now the geometry is not symmetric in x nor in y:



- The operator B_{θ} is no longer \mathcal{PT} -symmetric and we expect:
 - No trapped modes
 - No invariance of the spectrum by complex conjugation.

• Discretized spectrum of B_{θ} in k (not in k^2). We take $\theta = \pi/4$.



• We compute $k \mapsto |R(k)|$ for $0 < k < \pi$.



• We compute $k \mapsto |R(k)|$ for $0 < k < \pi$.



Complex eigenvalues also contain information on almost no reflection.

17 / 22

Spectra for a changing geometry

▶ Two series of computations: one with PMLs with different sign, one with classical PMLs. We compute the spectra for $h \in (1.3; 8)$.



The magenta marks on the real axis correspond to $k = \pi/\ell \& k = 2\pi/\ell$. For $k = 2\pi/\ell$, trapped modes and T = 1 should occur for certain h.

• We zoom at the region
$$0 < \Re e k < \pi$$
.

* PMLs with different signs

+ Classical PMLs

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Future work

- 1) How to justify the numerics? Absence of spectral pollution?
- 2) Can we find a spectral approach to compute completely reflecting or completely invisible k for a given geometry?
- 3) Can we find a spectral approach to identify modal conversion?
- 4) Can we prove existence of non reflecting k for the \mathcal{PT} -symmetric pb?

21

Thank you for your attention!