Spectrum for a small inclusion of negative material

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Coll. with X.  $Claeys^2$  and S.A.  $Nazarov^3$ .

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Scattering by a negative material in electromagnetism in time-harmonic

Zoom on a metamaterial: practical realizations of metamaterials are achieved by a periodic assembly of small resonators.



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▶ Recently, artificial metamaterials have been realized which can be modelled (at some frequency of interest) by  $\varepsilon < 0$  and  $\mu < 0$ .

▶ In this talk, we investigate a Dirichlet spectral problem for a small inclusion of negative material in a bounded domain.

• Let  $\Omega$ ,  $\omega$  be smooth domains of  $\mathbb{R}^3$  such that  $O \in \omega$ ,  $\overline{\omega} \subset \Omega$ . For  $\delta \in (0, 1]$ , we consider the problem

 $\left| \begin{array}{l} \text{Find } (\lambda^{\delta}, u^{\delta}) \in \mathbb{C} \times (\mathrm{H}^{1}_{0}(\Omega) \setminus \{0\}) \text{ s.t.:} \\ -\mathrm{div}(\sigma^{\delta} \nabla u^{\delta}) = \lambda^{\delta} u^{\delta} \quad \text{in } \Omega, \text{ with,} \end{array} \right.$ 



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$$\mathrm{H}^1_0(\Omega) := \{ u \in \mathrm{H}^1(\Omega) \mid u = 0 \text{ on } \partial \Omega \}$$

• 
$$\sigma^{\delta} = \begin{vmatrix} \sigma_1 > 0 & \text{in} & \Omega_1^{\delta} := \Omega \setminus \overline{\delta \omega} \\ \sigma_2 < 0 & \text{in} & \Omega_2^{\delta} := \delta \omega. \end{vmatrix}$$



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• We define the operator  $A^{\delta} : D(A^{\delta}) \to L^2(\Omega)$  such that

$$\begin{split} D(\mathbf{A}^{\delta}) &= \{ u \in \mathbf{H}_0^1(\Omega) \, | \, \mathrm{div}(\sigma^{\delta} \nabla u) \in \mathbf{L}^2(\Omega) \} \\ \mathbf{A}^{\delta} u &= -\mathrm{div}(\sigma^{\delta} \nabla u). \end{split}$$

### Introduction: main question of the talk

▶ Using boundary integral equations (see Costabel and Stephan 85, Dauge and Texier 97) or the T-coercivity approach (see Bonnet-Ben Dhia *et al.* 99,10,12,13), we can prove the :

PROPOSITION. Assume that  $\sigma_2/\sigma_1 \neq -1$ . For  $\delta > 0$ , the operator  $A^{\delta}$  is selfadjoint and has compact resolvent. Its spectrum  $\mathfrak{S}(A^{\delta})$  consists in two sequences of isolated eigenvalues:

$$-\infty \underset{n \to +\infty}{\leftarrow} \dots \lambda_{-n}^{\delta} \leq \dots \leq \lambda_{-1}^{\delta} < 0 \leq \lambda_1^{\delta} \leq \lambda_2^{\delta} \leq \dots \leq \lambda_n^{\delta} \dots \xrightarrow[n \to +\infty]{} +\infty.$$

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What happens to the negative spectrum when  $\delta$  tends to zero?

#### 1 Limit operators

We introduce the two natural limit operators which appear when  $\delta \to 0$ .

#### 2 Results

We state the main results concerning the asymptotic behaviour of the eigenvalues when  $\delta \to 0$ .

#### **3** Numerical experiments

We illustrate the theoretical results with numerical experiments.





**3** Numerical experiments













• As  $\delta \to 0$ , the small inclusion of negative material disappears.



• We introduce the far field operator  $A^0$  such that

$$D(\mathbf{A}^{\mathbf{0}}) = \{ v \in \mathbf{H}_{0}^{1}(\Omega) \mid \Delta v \in \mathbf{L}^{2}(\Omega) \}$$
$$\mathbf{A}^{\mathbf{0}}v = -\sigma_{1}\Delta v.$$

PROPOSITION. There holds  $\mathfrak{S}(\mathbf{A}^0) = {\{\mu_n\}_{n \ge 1} \text{ with } 0 < \mu_1 < \mu_2 \le \cdots \le \mu_n \dots \xrightarrow[n \to +\infty]{} +\infty.$ 

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• Define the near field operator  $B^{\infty}$  such that

$$D(\mathbf{B}^{\infty}) = \{ w \in \mathrm{H}^{1}(\mathbb{R}^{3}) \mid \operatorname{div} (\sigma^{\infty} \nabla w) \in \mathrm{L}^{2}(\mathbb{R}^{3}) \}$$
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PROPOSITION. Assume that  $\sigma_2/\sigma_1 \neq -1$ . The continuous spectrum of  $\mathbf{B}^{\infty}$  is equal to  $[0; +\infty)$  while its discrete spectrum is a sequence of eigenvalues:  $\mathfrak{S}(\mathbf{B}^{\infty}) \setminus \overline{\mathbb{R}_+} = \{\mu_{-n}\}_{n \geq 1}$  with  $0 > \mu_{-1} \geq \cdots \geq \mu_{-n} \cdots \xrightarrow[n \to +\infty]{n \to +\infty} -\infty$ .





**3** Numerical experiments

Assume that  $\sigma_2/\sigma_1 \neq -1$  and that  $\mathbf{B}^{\infty}$  is injective. For  $n \in \mathbb{N}^*$ , we denote  $\lambda_{\pm n}^{\delta}, \mu_n^{\delta}, \mu_{-n}^{\delta}$  the eigenvalues of  $\mathbf{A}^{\delta}, \mathbf{A}^0, \mathbf{B}^{\infty}$  as in the previous slides.

THEOREM. (POSITIVE SPECTRUM) For all  $n \in \mathbb{N}^*$ ,  $\varepsilon \in (0; 1)$ , there exist constants  $C, \delta_0 > 0$  depending on  $n, \varepsilon$  but independent of  $\delta$ , such that

 $|\lambda_n^{\delta} - \mu_n| \le C \, \delta^{3/2 - \varepsilon}, \qquad \forall \delta \in (0; \delta_0].$ 

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IDEA OF THE PROOF:

**1** We prove the *a priori* estimate  $||u^{\delta}||_{H_0^1(\Omega)} \leq c ||A^{\delta}u^{\delta}||_{\Omega}$  for  $\delta$  small enough ( $\blacklozenge$  hard part of the proof: weighted Sobolev spaces+overlapping cut-off functions +Nazarov's technique).

**2** If  $(\mu_n, v_n)$  is an eigenpair of  $A^0$ , we construct u such that

$$\|\mathbf{A}^{\delta}u - \mu_n u\|_{\Omega} \le c \,\delta^{\beta} \|u\|_{\Omega}, \qquad \text{for some } \beta > 0.$$

**3** If  $(\lambda_n^{\delta}, u_n^{\delta})$  is an eigenpair of  $A^{\delta}$ , we construct v such that

$$\|\mathbf{A}^0 v - \lambda_n^{\delta} v\|_{\Omega} \le c \,\delta^{\beta} \|v\|_{\Omega}, \qquad \text{for some } \beta > 0.$$

We conclude with a classical lemma on quasi eigenvalues.

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 $|\lambda_{-n}^{\delta} - \delta^{-2} \mu_{-n}| \le C \exp(-\gamma/\delta), \qquad \forall \delta \in (0; \delta_0].$ 



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PROPOSITION. (LOCALIZATION EFFECT) For all  $n \in \mathbb{N}^*$ , let  $u_{-n}^{\delta}$  be an eigenfunction corresponding to the negative eigenvalue  $\lambda_{-n}^{\delta}$ . There exist constants  $C, \gamma, \delta_0 > 0$ , depending on n but independent of  $\delta$ , such that

$$\int_{\Omega} (|u_{-n}^{\delta}|^2 + |\nabla u_{-n}^{\delta}|^2) e^{\gamma x/\delta} d\boldsymbol{x} \le C \, \|u_{-n}^{\delta}\|_{\Omega}, \qquad \forall \delta \in (0; \delta_0].$$







▶ Using FreeFem++, we approximate numerically the spectrum of  $A^{\delta}$  using a usual P1 Finite Element Method. We solve the problem

$$\begin{vmatrix} \text{Find } (\lambda_h^{\delta}, u_h^{\delta}) \in \mathbb{C} \times (\mathcal{V}_h \setminus \{0\}) \text{ s.t.} \\ \int_{\Omega} \sigma_h^{\delta} \nabla u_h^{\delta} \cdot \nabla v_h = \lambda_h^{\delta} \int_{\Omega} u_h^{\delta} v_h, \quad \forall v_h \in \mathcal{V}_h, \end{vmatrix}$$

where  $V_h$  approximates  $H_0^1(\Omega)$  as  $h \to 0$  (*h* is the mesh size).

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• We consider the following 2D geometry:



We display the spectrum as  $\delta \to 0$  (*h* is more or less fixed).



• The positive part of  $\mathfrak{S}(\mathbf{A}^{\delta})$  converges to  $\mathfrak{S}(\mathbf{A}^{0})$  when  $\delta \to 0$ .



The negative part of  $\mathfrak{S}(A^{\delta})$  is asymptotically equivalent to the negative part of  $\delta^{-2}\mathfrak{S}(B^{\infty})$  when  $\delta \to 0$ .

Contrast  $\kappa_{\sigma} = -2.5$ 



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# Localization effect

Eigenfunction associated to the first negative eigenvalue

Eigenfunction associated to the first positive eigenvalue



► The eigenfunctions corresponding to the negative eigenvalues are localized around the small inclusion. Here,  $\sigma_2/\sigma_1 = -2.5$ .

# Thank you for your attention!!!

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