

An explicit algorithm for solving the acoustic tomography problem for a moving fluid

Alexey Agaltsov

agaltsov@cmap.polytechnique.fr

Moscow

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Acoustic tomography of moving fluid

- A moving fluid in a bounded domain $D \subset \mathbb{R}^d$, $d \geq 2$, is characterized by sound speed $c = c(x)$, density $\rho = \rho(x)$, velocity $\mathbf{v} = \mathbf{v}(x)$ and absorption $\alpha = \omega^{\zeta(x)}\alpha_0(x)$
- There are acoustic transducers on ∂D . A transducer produces time-harmonic acoustic waves which are scattered by the fluid. Scattered acoustic waves are recorded by other transducers.

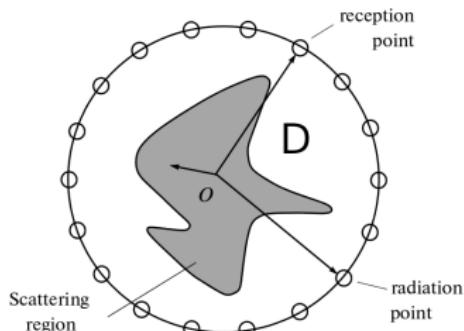


image: (Burov et al. '13)

Acoustic tomography problem. Given this data, recover fluid parameters.

Main applications in ocean tomography (*determine the ocean temperature and heat transferring currents*) and in medical diagnostics (*determine scalar inhomogeneities and the blood flow*)

Acoustic tomography of moving fluid

$$L_\omega = -\Delta - 2i \left(\frac{\omega \mathbf{v}}{c^2} + \frac{i}{2} \nabla \ln \rho \right) \cdot \nabla - \frac{\omega^2}{c^2} - 2i\omega \frac{\alpha}{c} \quad (\text{AC})$$

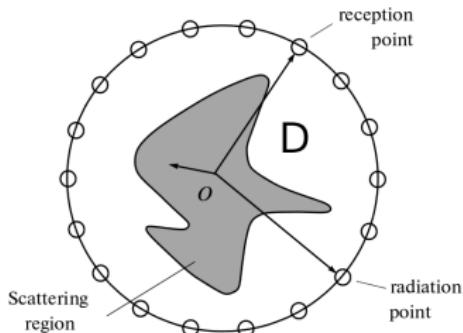


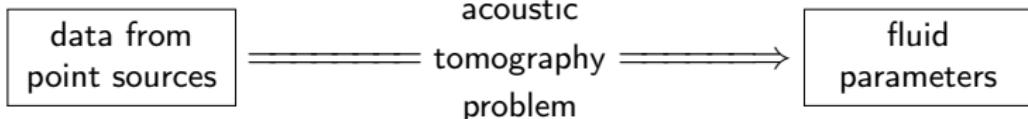
image: (Burov et al. '13)

Data from point sources: $G_\omega|_{X \times Y}$, $\omega \in \Omega$, where $X, Y \subset \partial D$, $\Omega \subset \mathbb{R}_{\geq 0}$,

$$\begin{cases} L_\omega G_\omega(x, y) = -\delta_y(x), & x \in \mathbb{R}^d, \\ G_\omega(\cdot, y) \text{ radiates at } \infty \end{cases}$$

Acoustic tomography problem

Given $G_\omega|_{X \times Y}$ for $\omega \in \Omega$ and c_0 , find c , \mathbf{v} , $\nabla \rho$ and α in D



- We consider the following operator with smooth coefficients:

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + i A_j(x) \right)^2 + Q(x), \quad (\text{OP})$$

where $x = (x_1, \dots, x_d) \in D$,

$A = (A_1, \dots, A_d)$, $A_j(x) \in M_n(\mathbb{C})$, $Q(x) \in M_n(\mathbb{C})$,

D is an open bounded domain in \mathbb{R}^d with boundary ∂D

- $L_{A,Q}$ acts on \mathbb{C}^n -valued functions in D

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + i A_j(x) \right)^2 + Q(x), \quad (\text{OP})$$

- Suppose that $E \in \mathbb{C}$ is not a DE for $L_{A,Q}$ in D :

$$\begin{cases} L_{A,Q}\psi = E\psi & \text{in } D, \\ \psi|_{\partial D} = f, \end{cases}$$

is uniquely solvable for any sufficiently regular f on ∂D .

- The Dirichlet-to-Neumann map $\Lambda_{A,Q} = \Lambda_{A,Q}(E)$:

$$\Lambda_{A,Q} f = \sum_{j=1}^d \nu_j \left(\frac{\partial}{\partial x_j} + i A_j \right) \psi \Big|_{\partial D}, \quad (\text{DN})$$

where $\nu = (\nu_1, \dots, \nu_d)$ is the unit exterior normal to ∂D .

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + i A_j(x) \right)^2 + Q(x), \quad (\text{OP})$$

$$\Lambda_{A,Q} f = \sum_{j=1}^d \nu_j \left(\frac{\partial}{\partial x_j} + i A_j \right) \psi|_{\partial D}, \quad (\text{DN})$$

- Conjugation of $L_{A,Q}$ by a smooth $GL_n(\mathbb{C})$ -valued function g :

$$\begin{cases} g L_{A,Q} g^{-1} = L_{A^g, Q^g}, \\ A_j^g = g A_j g^{-1} + i \frac{\partial g}{\partial x_j} g^{-1}, \quad j = 1, \dots, d, \\ Q^g = g Q g^{-1}. \end{cases} \quad (\text{GT})$$

- The following formula holds:

$$\Lambda_{A^g, Q^g} = g|_{\partial D} \Lambda_{A,Q}(g|_{\partial D})^{-1}.$$

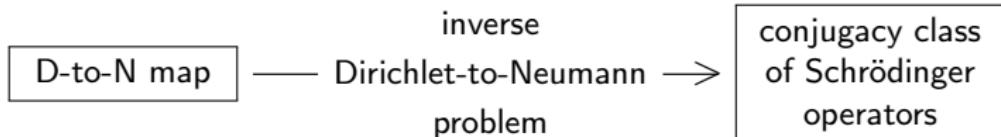
$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + iA_j \right)^2 + Q, \quad (\text{OP})$$

$$\Lambda_{A,Q}(\psi|_{\partial D}) = \sum_{j=1}^d \nu_j \left(\frac{\partial}{\partial x_j} + iA_j \right) \psi|_{\partial D}, \quad L_{A,Q}\psi = E\psi, \quad (\text{DN})$$

$$\begin{cases} gL_{A,Q}g^{-1} = L_{A^g, Q^g}, \\ \Lambda_{A^g, Q^g} = \Lambda_{A,Q}, \\ g \text{ is smooth } GL_n(\mathbb{C})\text{-valued, } g|_{\partial D} = \text{Id} \end{cases} \quad (\text{GT})$$

The inverse Dirichlet-to-Neumann problem

Given $\Lambda_{A,Q}$ at fixed E , find $L_{A,Q}$ modulo (GT).



The IDN problem: scalar case

- A_j, Q are scalar functions, $d \in \{2, 3\}$, $A = (A_1, \dots, A_d)$

$$L_{A,Q} = -(\nabla + iA)^2 + Q, \quad (\text{OP})$$

$$\Lambda_{A,Q} f = (\nu \cdot (\nabla + iA))\psi|_{\partial D}, \quad (\text{DN})$$

$$\begin{cases} e^{i\varphi} L_{A,Q} e^{-i\varphi} = L_{A^\varphi, Q^\varphi}, \\ A^\varphi = A + \nabla \varphi, \\ Q^\varphi = Q \end{cases} \quad (\text{GT})$$

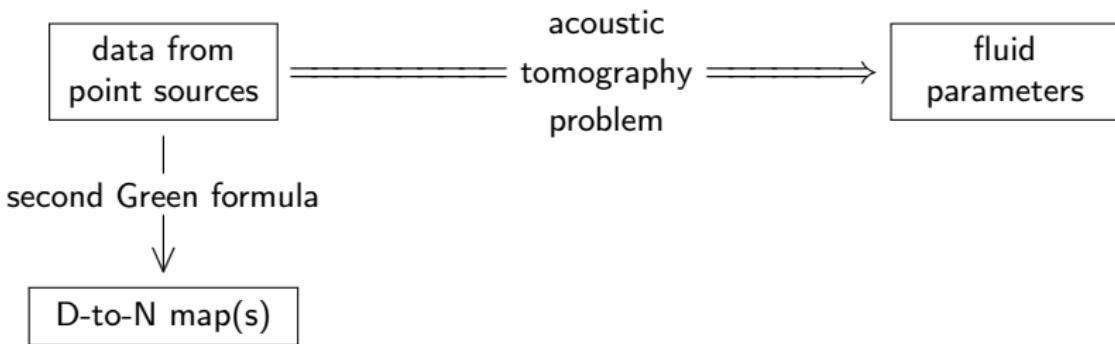
- $F = \operatorname{curl} A$ and Q are gauge invariant and are uniquely determined by $\Lambda_{A,\nu}(E)$, see [10] ($d \geq 3$) and [9] ($d = 2$)
- $(A - (A \cdot \nu)\nu)|_{\partial D}$ is uniquely determined by $\Lambda_{A,\nu}(E)$, see [6]

Acoustic scattering: reduction to the IDN problem

Use the second Green formula (Nachman '88):

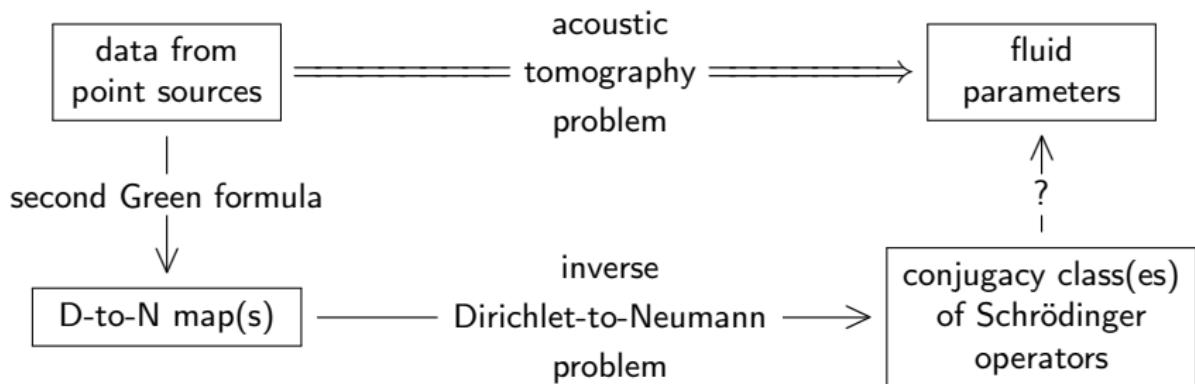
$$G_\omega(x, y) - G_\omega^0(x, y) = \int_{\partial D} \int_{\partial D} G_\omega^0(x, z)(\Lambda_\omega - \Lambda_\omega^0)(z, w) G_\omega(w, y) dy dw$$

where G_ω^0 , Λ_ω^0 correspond to $\mathbf{v} = 0$, $\nabla \rho = 0$, $c = c_0$, $\alpha = 0$.



Gauge fixing

- *Question.* Suppose that we know how to solve the IDN problem. How to complete the following diagram?



Gauge fixing

$$L_\omega = -\Delta - 2i \left(\frac{\omega \mathbf{v}}{c^2} + \frac{i}{2} \nabla \ln \rho \right) \cdot \nabla - \frac{\omega^2}{c^2} - 2i\omega \frac{\alpha}{c} \quad (\text{AC})$$

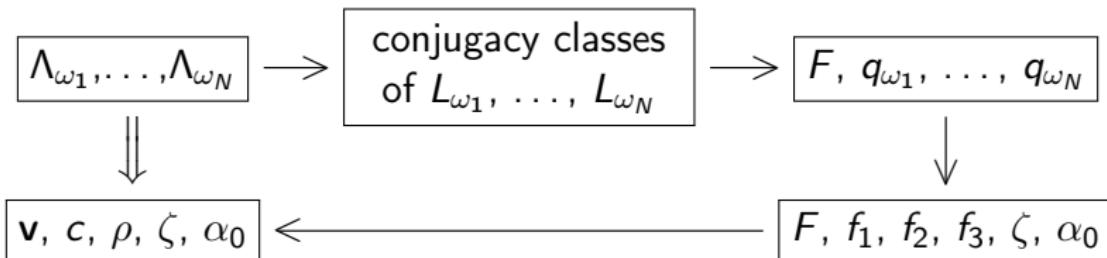
- functions F and q_ω are invariants of the conjugacy class:

$$F = \operatorname{curl} \frac{\mathbf{v}}{c^2},$$

$$q_\omega = f_1 - \omega^2 f_2 + i\omega f_3 - 2i\omega^{1+\zeta} \alpha_0,$$

$$f_1 = \rho^{\frac{1}{2}} \Delta \rho^{-\frac{1}{2}}, \quad f_2 = \frac{1}{c^2} + \frac{\mathbf{v}}{c^2} \frac{\mathbf{v}}{c^2}, \quad f_3 = \nabla \cdot \left(\frac{\mathbf{v}}{c^2} \right) - \frac{\mathbf{v}}{c^2} \cdot \nabla \ln \rho$$

- The fluid parameters can be recovered as follows:



$$L_\omega = -\Delta - 2i \left(\frac{\omega \mathbf{v}}{c^2} + \frac{i}{2} \nabla \ln \rho \right) \cdot \nabla - \frac{\omega^2}{c^2} - 2i\omega^{1+\zeta} \frac{\alpha_0}{c}, \quad (\text{AC})$$
$$\Lambda_\omega(\psi|_{\partial D}) = \frac{\partial \psi}{\partial \nu}|_{\partial D}, \quad L_\omega \psi = 0.$$

- $\rho \equiv \rho_0, \alpha_0 \equiv 0 \implies \Lambda_\omega$ at fixed ω determines \mathbf{v}, c
- $\alpha_0 \equiv 0 \implies \Lambda_\omega$ at 2 ω 's determines \mathbf{v}, c, ρ
- $\zeta \neq 0 \implies \Lambda_\omega$ at 3 ω 's determines $\mathbf{v}, c, \rho, \zeta, \alpha_0$
- Explicit examples of non-uniqueness when $\zeta \equiv 0$

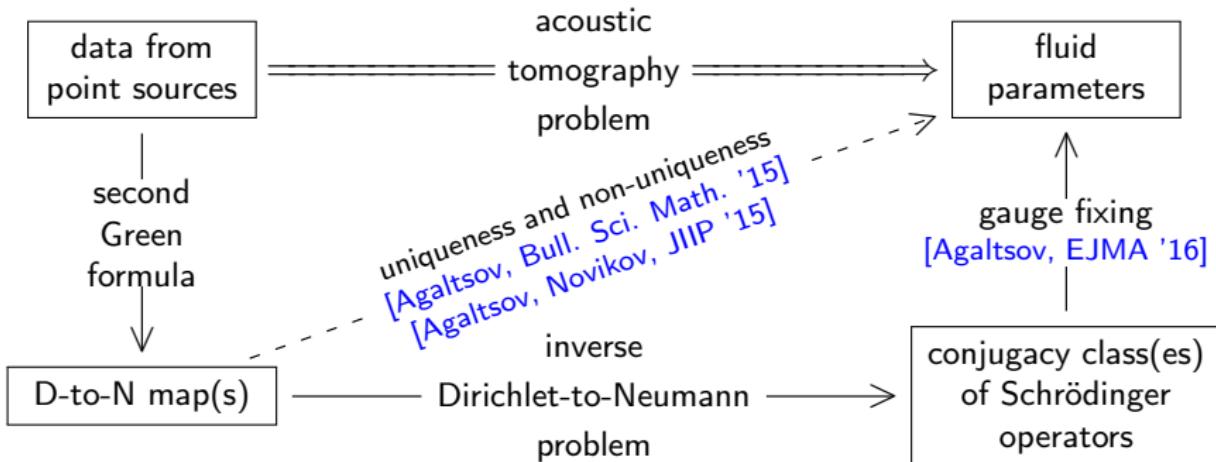
[Agaltsov, Bull. Sci. Math. '15]: uniqueness

[Agaltsov, Novikov, JIIP '15] : uniqueness and invisible fluids

[Agaltsov, EJMA '16]: algorithms

Solving the acoustic tomography problem

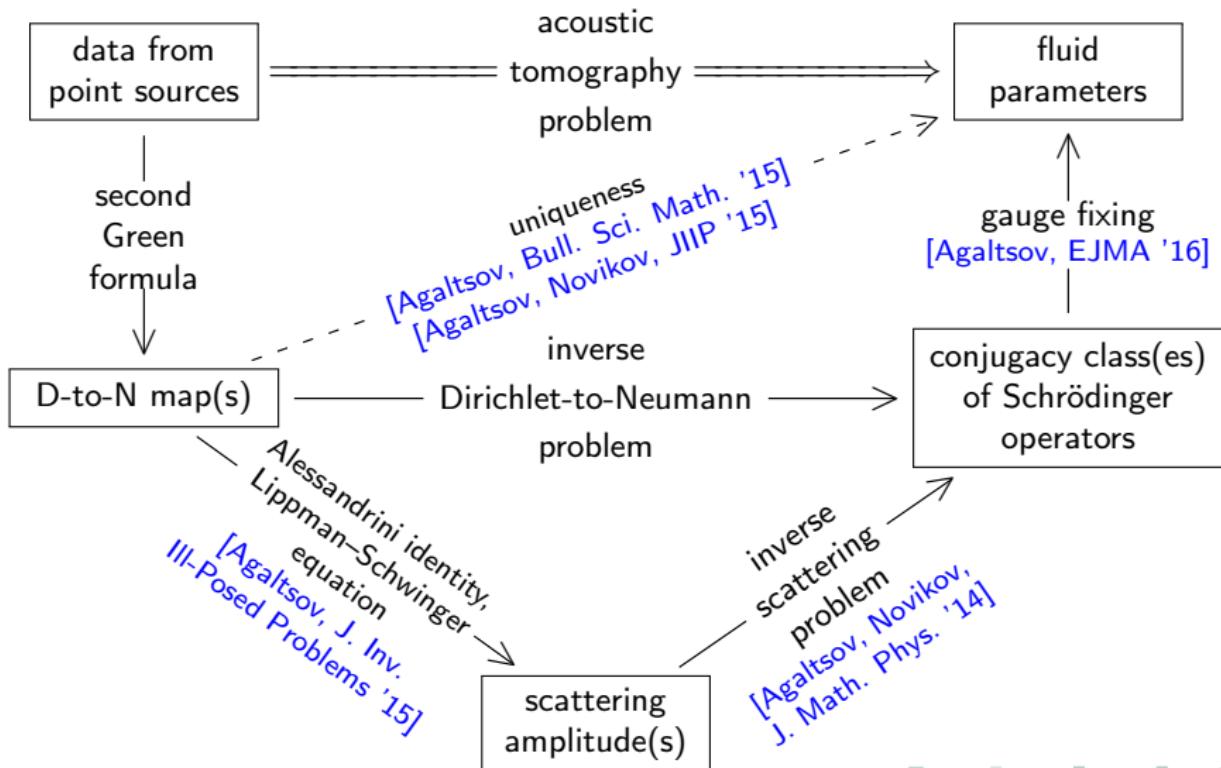
- So far we have the following scheme with vertical arrows being explicit algorithms:



- Question.* How to solve constructively the inverse Dirichlet-to-Neumann problem?

Solving the acoustic tomography problem

A common project with Moscow University Acoustical Physics group (Burov et al.)



The inverse scattering problem

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + iA_j \right)^2 + Q, \quad (\text{OP})$$

A_j, Q are smooth $M_n(\mathbb{C})$ -valued in D

- Set A, Q equal to zero outside of D
- Consider functions $\psi^+(\cdot, k)$, $k \in S_{\sqrt{E}}^{d-1} = \{\kappa \in \mathbb{R}^d \mid \kappa^2 = E\}$:

$$\begin{cases} L_{A,Q}\psi^+(x, k) = E\psi^+(x, k), & x \in \mathbb{R}^d, \\ \psi^+(x, k) = e^{ikx} \text{Id}_n + \psi_{\text{sc}}^+(x, k), \\ \psi_{\text{sc}}^+ \text{ radiates at } \infty \end{cases}$$

- The scattering amplitude $f_{A,Q}$ on $\mathcal{M}_E = S_{\sqrt{E}}^{d-1} \times S_{\sqrt{E}}^{d-1}$:

$$\psi_{\text{sc}}^+(x, k) = \begin{pmatrix} \text{universal} \\ \text{spherical} \\ \text{wave} \end{pmatrix} \cdot f_{A,Q} \left(k, \frac{|k|}{|x|} x \right) (1 + o(1)), \quad |x| \rightarrow \infty.$$

Direct scattering problem

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + iA_j \right)^2 + Q \quad (\text{OP})$$

Direct scattering problem

Given $L_{A,Q}$, find $f_{A,Q}$.

- $\psi^+(\cdot, k)$ satisfies the Lippmann-Schwinger equation:

$$\psi^+(x, k) = e^{ikx} \text{Id}_n + \int_D G^+(x - y, k)(L_{A,Q} - L_{0,0})\psi^+(y, k) dy, \quad (\text{LS})$$

$$G^+(x, k) = -(2\pi)^{-d} \int_{\mathbb{R}^d} \frac{e^{i\xi x} d\xi}{\xi^2 - k^2 - i0} \underset{\sim}{\simeq} \frac{1}{E - L_{0,0}}$$

- The scattering amplitude $f_{A,Q}$ can be found from:

$$f_{A,Q}(k, l) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{-ilx} (L_{A,Q} - L_{0,0})\psi^+(x, k) dx \quad (\text{SA})$$

Gauge invariance

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + i A_j \right)^2 + Q, \quad (\text{OP})$$

$$f_{A,Q}(k, l) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{-ilx} (L_{A,Q} - L_{0,0}) \psi^+(x, k) dx \quad (\text{SA})$$

- Conjugation of $L_{A,Q}$ by a $GL_n(\mathbb{C})$ -valued function g :

$$\begin{cases} g L_{A,Q} g^{-1} = L_{A^g, Q^g}, \\ A_j^g = g A_j g^{-1} + i \frac{\partial g}{\partial x_j} g^{-1}, \quad j = 1, \dots, d, \\ Q^g = g Q g^{-1} \end{cases} \quad (\text{GT})$$

- The amplitude is gauge invariant:

$$f_{A^g, V^g}(k, l) = f_{A, V}(k, l),$$

if $g \rightarrow \text{Id}$ at ∞ sufficiently fast

The inverse scattering problem

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + i A_j \right)^2 + Q, \quad (\text{OP})$$

$$f_{A,Q}(k, l) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{-ilx} (L_{A,Q} - L_{0,0}) \psi^+(x, k) dx \quad (\text{SA})$$

$$\begin{cases} g L_{A,Q} g^{-1} = L_{A^g, Q^g}, \\ f_{A^g, Q^g} = f_{A,Q}, \\ g \text{ is smooth } GL_n(\mathbb{C})\text{-valued}, \\ g \rightarrow \text{Id at } \infty \text{ sufficiently fast} \end{cases} \quad (\text{GT})$$

The inverse scattering problem

Given $f_{A,Q}$ at fixed E , find $L_{A,Q}$ modulo (GT).

[Agaltsov, Novikov, J. Math. Phys. '14]: algorithm

The inverse scattering problem

- Born approx. ($A = 0$, $Q \in C_c^n$), see (Faddeev, '56; R. Novikov '15)

$$\widehat{Q}(k - I) = f(k, I) + O(E^{-\frac{1}{2}}) \implies Q = Q^B + O(E^{-\frac{n-2}{2n}})$$

- Generalized scatt. sol. $\psi(x, k) = e^{ikx} \mu(x, k)$ (Faddeev '65), where $k = k(\lambda)$, $\lambda \in \mathbb{C} \setminus S^1$ и $k \in \mathbb{C}^2 \setminus \mathbb{R}^2$, $k^2 = E$
- For $\lambda \in \mathbb{C} \setminus S^1 \cup 0$ (Grinevich, Manakov '86; R. Novikov '99)

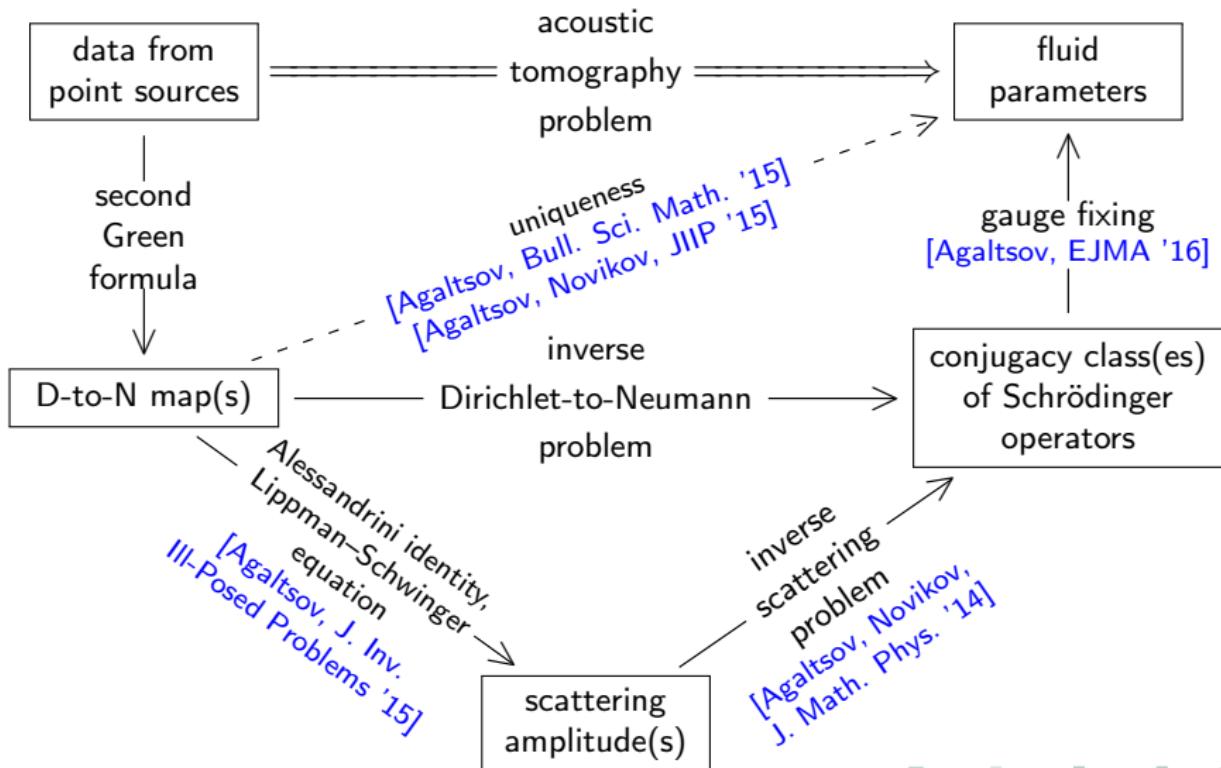
$$\begin{aligned}\overline{\partial}_\lambda \mu(x, \lambda) &= r(x, \lambda) \mu(x, -\frac{1}{\bar{\lambda}}) \approx 0, \quad \text{since } r(x, \cdot) = O(E^{-\frac{n}{2}}), \\ \mu_+(x, \lambda) &= \mu_-(x, \lambda) + \int_{S^1} \rho(x, \lambda, \lambda') \mu_-(x, \lambda') d\lambda', \quad \lambda \in S^1, \\ \mu(x, \infty) &= 1, \quad \mu(x, \cdot)|_{S^1 \pm 0} = \mu_\mp,\end{aligned}$$

If $r \equiv 0$, we get the nonlocal Riemann-Hilbert problem. Study of such problems goes back to (Manakov '81). Finding ρ from f (R. Novikov '86, '86). Solution for $\rho(x, \lambda, \lambda') = \overline{\rho(x, \lambda', \lambda)}$ (Grinevich, R. Novikov '85, '86)

- $\psi(x, k(\lambda)) = e^{ikx} (\mu_0^\pm + \mu_1^\pm \lambda^{\pm 1} + \dots)$ collect coefficients of λ^0 , $\lambda^{\pm 1}$ in $L_{A, Q} \psi = E \psi$, for obtaining expressions for A , Q in terms of μ_k^\pm (Grinevich, R. Novikov '85)
- Solution for $A = 0$ and estimate $Q = \tilde{Q} + O(E^{-\frac{n-2}{2}})$ (R. Novikov '98, '99).
General solution [Agaltsov, Novikov, JMP '14]

Solving the acoustic tomography problem

A common project with Moscow University Acoustical Physics group (Burov et al.)



From D-to-N map to scattering amplitude

$$L_{A,Q} = - \sum_{j=1}^d \left(\frac{\partial}{\partial x_j} + iA_j \right)^2 + Q, \quad (\text{OP})$$

$$\Lambda_{A,Q}(\psi|_{\partial D}) = \sum_{j=1}^d \nu_j \left(\frac{\partial}{\partial x_j} + iA_j \right) \psi|_{\partial D}, \quad L_{A,Q}\psi = E\psi,$$

A_j, Q are $M_n(\mathbb{C})$ -valued with compact support in D (DN)

- Let u_0 satisfy $L_{0,0}u_0 = Eu_0$ and u satisfy $L_{A,Q}u = Eu$ in D .
Then the following formula holds (Alessandrini '88):

$$\int_D u_0(x)(L_{A,Q} - L_{0,0})u(x)dx = \int_{\partial D} u_0(x)(\Lambda_{A,Q} - \Lambda_{0,0})u(x)dx \quad (\text{AI})$$

From D-to-N map to scattering amplitude

$$\psi^+(x, k) = e^{ikx} + \int_D G^+(x - y, k)(L_{A,Q} - L_{0,0})\psi^+(y, k) dy \quad (\text{LS})$$

$$f_{A,Q}(k, l) = (2\pi)^{-d} \int_D e^{-ilx}(L_{A,Q} - L_{0,0})\psi^+(x, k) dx \quad (\text{SA})$$

$$\int_D u_0(x)(L_{A,Q} - L_{0,0})u(x)dx = \int_{\partial D} u_0(x)(\Lambda_{A,Q} - \Lambda_{0,0})u(x)dx \quad (\text{AI})$$

- (LS) + (AI) imply

$$\psi^+(x, k) = e^{ikx} + \int_{\partial D} G^+(x - y, k)(\Lambda_{A,Q} - \Lambda_{0,0})\psi^+(y, k) dy$$

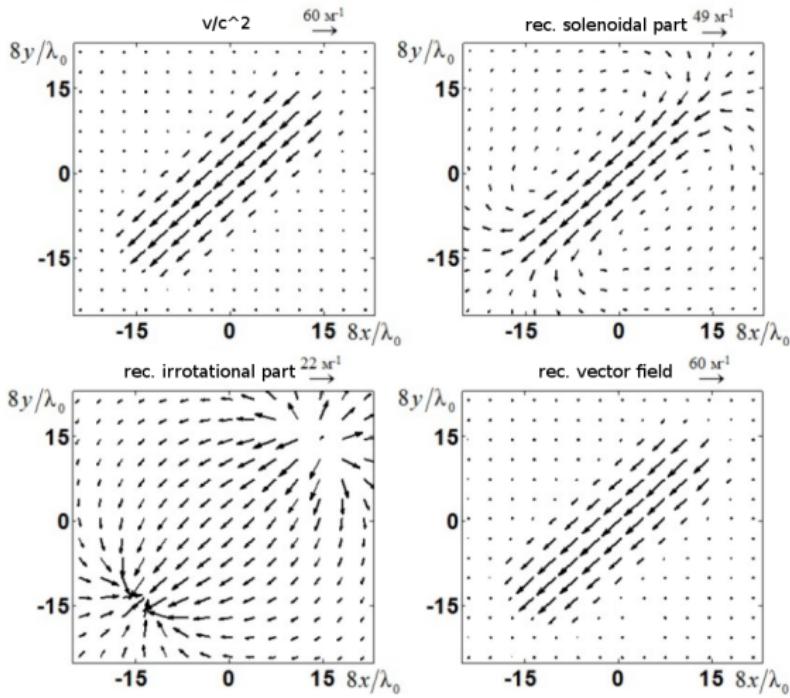
- (SA) + (AI) imply

$$f_{A,Q}(k, l) = (2\pi)^{-d} \int_{\partial D} e^{-ilx}(\Lambda_{A,Q} - \Lambda_{0,0})\psi^+(x, k) dx$$

(R. Novikov '92): $A = 0$, (Eskin-Ralston '97): $A \neq 0$, (R. Novikov '05): background potential, [Agaltsov, JIIP '15]: general case, matrix coefficients

Acoustic scattering: numerical example of [11]

$$L_\omega = -\Delta - 2i\omega \frac{\mathbf{v}}{c^2} \cdot \nabla - \frac{\omega^2}{c^2} - 2i\omega^2 \frac{\alpha_0}{c}, \quad \omega \in \{\omega_1, \omega_2\}$$



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