

# The two dimensional inverse conductivity problem

Dedicated to Gennadi

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IMJ

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- $(M, \sigma)$  is a **2-dimensional conductivity structure** when  $M$  is an abstract 2 dimensional manifold with boundary ( $M \cap bM = \emptyset$ ) and  $\sigma : T^*M \rightarrow T^*M$  is a tensor such that

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- **Dirichlet operator**  $D_\sigma$ . For  $u \in C^0(bM, \mathbb{R})$ ,  $D_\sigma u \in C^0(\overline{M})$  is defined by

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- **Neumann-Dirichlet operator**  $N_\sigma$ . For  $u : bM \rightarrow \mathbb{R}$  sufficiently smooth,  $N_\sigma u$  is defined by

$$N_\sigma u = \frac{\partial}{\partial \nu} D_\sigma u : bM \rightarrow \mathbb{R}$$

where  $\nu \in T_{bM}\overline{M}$  is the outer unit normal vector field of  $bM$ .

- Using *isothermal coordinates*, one find out that

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- If  $M$  is a submanifold of  $\mathbb{R}^3$ ,  $\sigma$  is **isotropic** when  $\mathcal{C}_{\sigma}$  is induced by the standard euclidean metric of  $\mathbb{R}^3$ . Likewise,  $\sigma$  is said isotropic relatively to a complex structure  $\mathcal{C}$  on  $M$  if  $*_{\sigma}$  is the Hodge operator of  $(M, \mathcal{C})$ .

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- Dirichlet problem for  $(\overline{M}, \sigma)$ . For  $u \in C^0(bM)$ , seek  $U$  such that

$$dsd^{\sigma}U = 0 \quad \& \quad U|_{bM} = u$$

where  $s = \det \sigma$ ,  $d^{\sigma} = i(\overline{\partial}^{\sigma} - \partial^{\sigma})$ ,  $\overline{\partial}^{\sigma}$  is the standard Cauchy-Riemann operator associated to the Riemann surface  $(M, \mathcal{C}_{\sigma})$  and  $\partial^{\sigma} = d - \overline{\partial}^{\sigma}$ .

## Inverse conductivity problem

**Data** :  $bM, \nu \in T_{bM}\overline{M}, \sigma|_{bM}$  and  $N_\sigma$

**Problem** : reconstruct  $M$  as a Riemann surface equipped with the conductivity tensor  $\sigma$ .

**Remark** : Let  $\varphi : \overline{M} \rightarrow \overline{M}$  be a  $C^1$ -diffeomorphism such that  $\varphi|_{bM} = Id_{bM}$  and  $\tilde{\sigma} = \varphi_*\sigma$ . Then  $N_{\tilde{\sigma}} = N_\sigma$  and  $\tilde{\sigma} \neq \sigma$  but  $(M, \mathcal{C}_{\tilde{\sigma}})$  and  $(M, \mathcal{C}_\sigma)$  represent the same (abstract) Riemann surface.

**Consequence** : non uniqueness up to a diffeomorphism gives different representations of the same Riemann surface.



- non exhaustive list for uniqueness results, and for a domain in  $\mathbb{R}^2$ , reconstructions process by Kohn-Vogelius (1984), Novikov (1988), Lee-Uhlman (1989), Sylvester (1990), Lassas-Uhlman (2001), Belishev (2003), Nachman (1996),
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  - 2 use of H-N to produce the function  $s : S \rightarrow \mathbb{R}_+^*$  such that  $s \cdot *$  is the pushforward of  $\sigma$ .

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  - 3  $(S, s \cdot *)$  is a solution to our inverse conductivity problem.

## Tools for step 1

- 1 Based on Henkin-Michel (2014) : explicit formulas for a Green function of the bordered nodal Riemann surface  $\mathcal{M}$ . This enable to compute for a given  $u : bM \rightarrow \mathbb{R}$  the  $C_\sigma$ -harmonic extension  $\tilde{u}$  ( $dd^\sigma \tilde{u} = 0$ ) of  $u$  from  $N_\sigma$
- 2 Based on Henkin-Michel (2012) : embedding  $S$  of  $M$  in  $\mathbb{C}P_4$  by a generic canonical map  $(\partial\tilde{u}_0 : \partial\tilde{u}_1 : \partial\tilde{u}_2 : \partial\tilde{u}_3)$  ;  $S$  is given as the solution of boundary problem. Then we seek an atlas for  $S$ . For generic data,  $S$  is covered by preimages of regular parts of the images  $Q$  and  $Q'$  of  $S \setminus \{(0 : 0 : 0 : 1)\}$  and  $S \setminus \{(0 : 0 : 1 : 0)\}$  under the projections  $\mathbb{C}P_4 \rightarrow \mathbb{C}P_3$ ,  $(w_0 : w_1 : w_2 : w_3) \mapsto (w_0 : w_1 : w_2)$  and  $(w_0 : w_1 : w_2 : w_3) \mapsto (w_0 : w_1 : w_3)$ . This reduces the problem by one dimension.

Let  $Q$  be the (possibly singular) nodal complex curve which is the image of  $S$  by the projection  $\mathbb{C}P_4 \rightarrow \mathbb{C}P_3$ ,  $(w_0 : w_1 : w_2 : w_3) \mapsto (w_0 : w_1 : w_2)$ . Some generic assumptions are made on  $\bar{Q}$  such like  $(0 : 1 : 0) \notin \bar{Q}$  and  $bQ \subset \{w_0 w_1 w_2 \neq 0\}$ . Let for  $z = (x, y) \in \mathbb{C}^2$ ,  $L_z$  be the line of equation  $\Lambda_z(w) \stackrel{\text{def}}{=} xw_0 + yw_1 + w_2 = 0$ . Then from Dolbeault-Henkin (1997), for  $z$  near a generic  $z_*$ ,

$$G_{\partial Q, k}(z) \stackrel{\text{def}}{=} \frac{1}{2\pi i} \int_{\partial Q} \left( \frac{w_1}{w_0} \right)^k \frac{d \frac{\Lambda_z(w)}{w_0}}{\frac{\Lambda_z(w)}{w_0}} = \sum_{1 \leq j \leq p} h_j(z)^k + P_k(z) \quad (\text{SWD})$$

where  $P_k \in \mathbb{C}(Y)_k[X]$  and  $h_1, \dots, h_k$  are holomorphic solutions of the **shock wave** equation

$$\frac{\partial h}{\partial y} = h \frac{\partial h}{\partial x}$$

such that

$$Q \cap L_z = \{(1 : h_j(z) : -x - yh_j(z)) ; 1 \leq j \leq p\}$$

- **Difficulty** : If  $K$  is algebraic and

$$K \cap L_z = \left\{ \left( 1 : \varphi_j(z) : -x - y\varphi_j(z) \right) ; 1 \leq j \leq q \right\},$$

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- 4 With Henkin (1995) and Collion (1996), one can prove that  $(g_j) = (h_j)$  and  $\tilde{P} = P$ .

## Building a SWD $G_{\partial Q,1} = \sum_{1 \leq j \leq d} g_j + P$

1. A function  $N$  is the sum of  $d$  different shock wave functions iff there exists  $s_1, \dots, s_d$  such that  $s_1 = -N$ , and

$$-s_d \frac{\partial N}{\partial x} + \frac{\partial s_d}{\partial y} = 0, \quad -s_k \frac{\partial N}{\partial x} + \frac{\partial s_k}{\partial y} = \frac{\partial s_{k+1}}{\partial x}, \quad 1 \leq k \leq d-1, \quad (1)$$

and the discriminant of  $T^d + s_1 T^{d-1} + \dots + s_d \in \mathcal{O}(D)[T]$  is not 0.

2. Let  $N = G_1 - P$  where  $P$  is of the form  $\frac{B'}{B}X + \frac{A}{B}$  with  $A, B \in \mathbb{C}[Y]$ ,  $B(0) = 1$  and  $\deg A < \deg B$ .  $(s_1, \dots, s_d)$  is a solution of (1) iff there exists one variable holomorphic functions  $\mu_1, \dots, \mu_d$  such that

$$s_k = \frac{e^H}{1 \otimes B} \left( \mathcal{E}^0(\mu_k \otimes 1) + \dots + \mathcal{E}^{d-k}(\mu_d \otimes 1) \right) \quad (2)$$

where  $H$  is such that  $\frac{\partial H}{\partial y} = \frac{\partial G_1}{\partial x}$ ,  $\mathcal{E}^j = \mathcal{E}^{j-1} \circ \mathcal{E}$  and  $\mathcal{E} = \int^y e^{-H} \frac{\partial}{\partial x} e^H$ .

- For a given  $\mu = (\mu_1, \dots, \mu_d)$ ,  $\left( (-1)^j s_j \right)_{1 \leq j \leq d}$  defined by (2) are the symmetric functions of a  $d$ -uple of shock wave functions iff  $s_1 = -N$  which turns out to be equivalent to a linear differential system

$$\forall n \in \mathbb{Z}, \quad \sum_{0 \leq m < j \leq d} c_{j,m}^{0,n} \mu_j^{(m)} = K_n^0(B, A, .) \quad (3)$$

where  $\left( c_{j,m}^{0,n} \right)$  depends only on  $G_1$  and  $K_n^0(B, A, .)$  vanishes for  $n \geq d$ , is linear with respect to  $(A, B)$  and has coefficients depending only on  $G_1$ .

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- The compatibility system has at least a solution. This solution gives birth to a solution  $\mu$  for (3) and thus to a SWD.

## Process

- If  $G_1$  is affine in  $x$ , solving the compatibility system on  $(A, B)$  and then (3) gives, after reduction, a parametrization of  $Q$  by its intersection with the lines  $L_z$ .



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- If  $G_1$  is affine in  $x$ ,  $Q$  is a "special" domain in an algebraic curve  $K$ . This reverts to the 1st case by choosing coordinates such that at least one line  $L_z$  hits  $Q$  and  $K \setminus Q$ .

## Green formula for singular complex curves

**Smooth case** symmetric dunction  $g : (\overline{M} \times \overline{M}) \setminus \Delta_{\overline{M}} \rightarrow \mathbb{R}$  such that for all  $q \in M$ ,

- $g_q = g(q, \cdot)$  is harmonic on  $M \setminus \{q\}$  and continuous on  $\overline{M} \setminus \{q\}$
- $g_q = -\frac{1}{2\pi} \ln |z|$  extends harmonically around  $q$  ( $z$  holomorphic coordinate  $z$  centered at  $q$ )

**Singular case** A Green function for a curve  $\mathcal{Y}$  in  $\mathbb{C}^2$  is a symmetric function  $g : (\text{Reg } \overline{\mathcal{Y}} \times \text{Reg } \overline{\mathcal{Y}}) \setminus \Delta_{\text{Reg } \overline{\mathcal{Y}}} \rightarrow \mathbb{R}$  s.t. for all  $q \in \text{Reg } \overline{\mathcal{Y}}$ ,  $g_q = g(q, \cdot)$  satisfies

$$i\partial\bar{\partial}g_q = \delta_q dV$$

in the sense of currents on  $\mathcal{Y}$ ,  $\delta_q$  being the Dirac measure supported by  $\{q\}$  and  $dV = i\partial\bar{\partial}|\cdot|^2$  - this implies in particular that  $\partial g_q$  is a weakly holomorphic  $(1, 0)$ -form on  $\mathcal{Y} \setminus \{q\}$  in the sense of Rosenlicht

### Nodal case

We demands that  $\mathcal{Y}$  extends as a usual continuous function along any branch except for one branch passing through  $q$  where it has an isolated logarithmic singularity.

Let us now detail a formula of Henkin-Michel (2014) establishing the existence of Green functions for a 1-paramter family of complex curves whose possible singularities are arbitrary. Let us consider a complex curve  $\mathcal{Y}$  in an open subset of  $\mathbb{C}^2$ ,  $\Omega$  a Stein neighborhood of  $\mathcal{Y}$  in  $\mathbb{C}^2$ ,  $\Phi$  a holomorphic function on  $\Omega$  such that  $\mathcal{Y} = \{\Phi = 0\}$  and  $d\Phi|_{\mathcal{Y}} \neq 0$  then a strictly pseudoconvex domain  $\Omega^*$  of  $\mathbb{C}^2$  verifying

$$\mathcal{Y}_0 = \mathcal{Y} \cap \Omega^* \subset \Omega,$$

and lastly a symmetric function  $\Psi \in \mathcal{O}(\Omega \times \Omega, \mathbb{C}^2)$  such that for all  $(z, z') \in \mathbb{C}^2$ ,

$$\Phi(z') - \Phi(z) = \langle \Psi(z', z), z' - z \rangle$$

where  $\langle v, w \rangle = v_1 w_1 + v_2 w_2$  when  $v, w \in \mathbb{C}^2$ . We define on  $\text{Reg}\mathcal{Y}$  a  $(1, 0)$ -form  $\omega$  by setting

$$\omega = \frac{-dz_1}{\partial\Phi/\partial z_2} \text{ on } \mathcal{Y}^1 = \mathcal{Y} \cap \{\partial\Phi/\partial z_2 \neq 0\}$$

$$\omega = \frac{+dz_2}{\partial\Phi/\partial z_1} \text{ on } \mathcal{Y}^2 = \mathcal{Y} \cap \{\partial\Phi/\partial z_1 \neq 0\}$$

and we consider

$$k(z', z) = \det \left[ \frac{\bar{z}' - \bar{z}}{|z' - z|^2}, \Psi(z', z) \right].$$

When  $q_* \in \text{Reg } \mathcal{Y}_0$ , H-M (2014) proves that the formula

$$\tilde{g}_{q_*}(q) = \frac{1}{4\pi^2} \int_{q' \in \mathcal{Y}_0} k(q', q) k(q_*, q') i\omega(q') \wedge \bar{\omega}(q'). \quad (4)$$

defines for  $\mathcal{Y}_0$  a Green function in the above sense and that if  $q_* \in \text{Reg } \mathcal{Y}_0$

$$\partial g_{q_*} = \tilde{k}_{q_*} \omega$$

where  $\tilde{k}_{q_*} = \frac{1}{2\pi} k(\cdot, q_*)$ .

## Proposition

*Suppose  $\mathcal{Y}$  has only nodal singularities. In this case, when  $q_* \in \text{Reg } \mathcal{Y}_0$ ,  $g_{q_*}$  extends as usual harmonic function along the branches of  $\mathcal{Y}_0 \setminus \{q_*\}$ ; in other words,  $\partial g_{q_*}$  extends as a standard holomorphic  $(1, 0)$ -form along the branches of  $\mathcal{Y}_0 \setminus \{q_*\}$ .*

## Corollary

*Suppose that  $\mathcal{Y}$  is an open nodal Riemann surface of  $\mathbb{C}^2$  and  $g$  is defined by (4). Then  $g$  is a simple Green function for  $\mathcal{Y}$ .*

## Corollary

Let  $(M, \sigma)$  be a two dimensional conductivity structure. We select, which is always possible, a two dimensional conductivity structure  $(\tilde{M}, \tilde{\sigma})$  extending plainly  $(M, \sigma)$ , which means that  $M \subset \subset \tilde{M}$ ,  $\tilde{\sigma}|_M = \sigma$  and  $\tilde{\sigma}|_p = \text{Id}_{T_p^* \tilde{M}}$  for all  $p \in b\tilde{M}$ . One denote then by  $F : \tilde{M} \rightarrow \mathbb{C}^2$  the map obtained by applying H-S theorem to  $(\tilde{M}, \tilde{\sigma})$ , we set  $\mathcal{Y} = F(\tilde{M})$  and fix a Stein neighborhood  $\Omega$  of  $\mathcal{Y}$  in  $\mathbb{C}^2$ . Lastly,  $\mathcal{M} = F(M)$  being relatively compact in  $\mathcal{Y}$ , we can pick up in  $\mathbb{C}^2$  a strictly pseudoconvex domain  $\Omega^*$  s.t.  $\mathcal{M} \subset \subset \mathcal{Y}_0 = \mathcal{Y} \cap \Omega^* \subset \Omega$ . We note  $g$  the function defined by (4). Then,  $F^* g \Big|_{\overline{M \times M} \setminus \Delta_M}$  is a Green function for  $(M, c_\sigma)$ .

## Corollary

$\mathcal{M}$  admits a principal Green function and if  $g$  is such a function, for all  $u \in C^\infty(b\mathcal{M})$ ,  $F^*\theta^{\mathcal{M}}f_*u$  is given by the formula

$$F^*\theta^{\mathcal{M}}f_*u = \left( F^*\widehat{\partial f_*u} \right) |_{b\mathcal{M}}, \widehat{f_*u} : \text{Reg } \overline{\mathcal{M}} \ni q \mapsto \begin{cases} \frac{i}{2} \int_{\partial\mathcal{M}} (f_*u) \bar{\partial}g_q & \text{si } q \in \mathcal{M} \\ f_*u(q) & \text{si } q \in b\mathcal{M} \end{cases}$$