

How to find a copy of the Hamiltonian from the scattering map  
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## 1 Statement of the problem

Let  $(M, g)$  be a  $n$ -dimensional smooth compact Riemannian manifold with boundary  $\partial M$  and  $g^t(x, \xi) = (\gamma_{(x, \xi)}(t), \dot{\gamma}_{(x, \xi)}(t))$  is the geodesic flow on  $TM_0 = \{(x, \xi) \in TM \mid \xi \in T_x, \xi \neq 0\}$ , where  $\gamma_{(x, \xi)}(t)$  is the geodesic, defined by initial data  $\gamma_{(x, \xi)}(0) = x$ ,  $\dot{\gamma}_{(x, \xi)}(0) = \xi$ , and  $\dot{\gamma}$  is the velocity vector. We assume that  $(M, g)$  is non-trapping, that is each maximal geodesic is finite. Formulate the following problem:

*To find an isometric copy  $\varphi^*g$ , where  $\varphi : M \rightarrow M$  is a diffeomorphism such that  $\varphi|_{\partial M} = id$ .*

More simple problem we address here is formulated as follows. Let  $(M, H)$  be a smooth  $n$ -dimensional Hamiltonian manifold with hamiltonian  $H(x, \xi) = g^{ij}(x)\xi_i\xi_j/2$ , and  $\varphi^t(x, \xi) = (\gamma_{(x, \xi)}(t), \dot{\gamma}_{(x, \xi)}(t))$  is the hamiltonian flow on  $T^*M_0 = \{(x, \xi) \in T^*M \mid \xi \in T_x^* \neq 0\}$ , where  $(\dot{\gamma}_{(x, \xi)}(t))_i = g_{ij}(\gamma_{(x, \xi)}(t))\dot{\gamma}_{(x, \xi)}^j(t)$ . In this talk we consider the following problem:

*To find a hamiltonian  $H$  up to a symplectomorphism  $f$  such that  $f|_{\partial T^*M_0} = id$ .*

## 2 Hamiltonian flow

Hamiltonian vector field  $\mathcal{H}$  is defined by the equality

$$\mathcal{H} = \frac{d\varphi^t}{dt}\Big|_{t=0}.$$

It is identified with Poisson's bracket  $[\cdot, H]$ :

$$\mathcal{H}u = \frac{d}{dt}(u \circ \varphi^t)\Big|_{t=0} = [u, H] = \sum_{i=1}^n (u_{x^i} H_{\xi_i} - u_{\xi_i} H_{x^i}). \quad (1)$$

Hamiltonian flow satisfies nonlinear Hamilton equation

$$\frac{d\varphi^t}{dt} = \mathcal{H}(\varphi^t),$$

or in coordinates

$$\begin{aligned}\frac{d\gamma^i_{(x,\xi)}(t)}{dt} &= H_{\xi_i}(\gamma_{(x,\xi)}(t), \dot{\gamma}_{(x,\xi)}(t)) \\ \frac{d(\gamma_{(x,\xi)})_i(t)}{dt} &= -H_{x^i}(\gamma_{(x,\xi)}(t), \dot{\gamma}_{(x,\xi)}(t))\end{aligned}$$

and Cauchy data

$$\varphi^0(x, \xi) = (x, \xi).$$

The following simple statement is very important.

**Lemma 1** *Hamiltonian flow satisfies the linear kinetic equation*

$$L\varphi^t := (\partial_t - \mathcal{H})\varphi^t = 0.$$

**Proof.** The hamiltonian flow is one-parameter group,  $\varphi^{t+s} = \varphi^t \circ \varphi^s = \varphi^s \circ \varphi^t$ . Then due to (1) one has

$$\frac{d}{ds}\varphi^s = \frac{d}{dt}(\varphi^{t+s})|_{t=0} = \frac{d}{dt}(\varphi^s \circ \varphi^t)|_{t=0} = [\varphi^s, H] = \mathcal{H}\varphi^s.$$

So, the hamiltonian flow satisfies two equations: nonlinear ODE  $d\varphi^t/dt = \mathcal{H}(\varphi^t)$  and linear PDE  $L\varphi^t = 0$ , and the Cauchy data  $\varphi^0(x, \xi) = (x, \xi)$ . ■

**Remark 2** *If*

$$L\psi = 0, \quad \psi(x, \xi, 0) = (x, \xi), \quad (2)$$

*then by the argument of uniqueness of the solution to the Cauchy problem (2)  $\psi(\cdot, t) = \varphi^t$ .*

### 3 Scattering map

Let  $M$  be a compact manifold with smooth boundary  $\partial M$ . The boundary of tangent space  $\partial T^*M_0$  may be decomposed into the sets of outer (+) and inner (-) covectors:

$$\partial T^*M_0 = \partial_+ T^*M_0 \cup \partial_- T^*M_0,$$

where

$$\begin{aligned}\partial_+ T^*M_0 &= \{(x, \xi) \in T^*M_0 \mid x \in \partial M, \pm(\nu(x), \xi) \geq 0\}, \\ \partial_- T^*M_0 &= \{(x, \xi) \in T^*M_0 \mid x \in \partial M, (\nu(x), \xi) = 0\},\end{aligned}$$

$\nu$  is the outer unit normal to the boundary  $\partial M$ . Denote by  $\tau(x, \xi)$  the length of geodesic ray  $\gamma_{(x,\xi)}(t), t \geq 0$ . Note, that  $\tau$  satisfies equation

$$\mathcal{H}\tau = \frac{d}{dt}(\tau \circ \varphi^t)|_{t=0} = \frac{d}{dt}(\tau - t)|_{t=0} = -1.$$

We put  $\tau(x, \xi) = 0$  for  $(x, \xi) \in \partial_+ T^*M_0$ . Denote by  $X$  the extended phase space

$$X = \{(x, \xi, t) \mid (x, \xi) \in T^*M_0, -\tau(x, -\xi) \leq t \leq \tau(x, \xi)\}.$$

Introduce notations:

$$\begin{aligned} \partial_- X &= \{(x, \xi, t) \mid (x, \xi) \in \partial_- T^*M_0, 0 \leq t \leq \tau(x, \xi)\} \\ \partial_+ X &= \{(x, \xi, t) \mid (x, \xi) \in \partial_+ T^*M_0, -\tau(x, -\xi) \leq t \leq 0\}, \\ \partial_0 X &= \partial_0 T^*M_0. \end{aligned}$$

Introduce the scattering map  $\alpha_H : \partial_{\mp} X \rightarrow \partial_{\pm} X$ . It is defined by the formulas

$$\begin{aligned} \alpha_H(x, \xi, t) &= (\gamma_{(x, \xi)}(\tau(x, \xi)), \dot{\gamma}_{(x, \xi)}(\tau(x, \xi)), t - \tau(x, \xi)), \quad (x, \xi, t) \in \partial_- X \\ \alpha_H(x, \xi, t) &= (\gamma_{(x, -\xi)}(\tau(x, -\xi)), \dot{\gamma}_{(x, -\xi)}(-\tau(x, -\xi)), \tau(x, -\xi) + t), \quad (x, \xi, t) \in \partial_+ X. \end{aligned}$$

Clearly,  $\alpha_H|_{\partial_0 X} = id$ , and  $\alpha_H$  is involution, that is  $\alpha_H^2 = id$ .

The solution to the Cauchy problem

$$\begin{aligned} Lu &= 0, \quad \text{in } X, \\ u|_{t=0} &= u_0, \quad \text{in } T^*M_0 \end{aligned}$$

is given by the formula

$$u(x, \xi, t) = u_0(\varphi^t(x, \xi)), \quad -\tau(x, -\xi) \leq t \leq \tau(x, \xi)$$

and its trace  $u' = u|_Y$  is even function w.r.t. involution  $\alpha_H$  :

$$u' = u' \circ \alpha_H,$$

where  $Y = \partial_- X \cup \partial_+ X$ . And other way around, any even function (w.r.t.  $\alpha_H$ )  $u'$  on  $Y$  may be extended onto  $X$  as a solution  $u$  of the kinetic equation:

$$Lu = 0, \quad u|_Y = u'.$$

For such extension we use notation  $u = \tilde{u}'$ . We use also notation

$$u_0 = u|_{t=0}.$$

## 4 Copy

The map  $\phi : T^*M_0 \rightarrow X$ ,  $\phi(x, \xi, t) = \varphi^t(x, \xi)$  satisfies kinetic equation

$$L\phi = 0, \quad \phi(x, \xi, 0) = (x, \xi).$$

For any  $t$  the map  $\varphi^t = (\gamma_{(\cdot)}(t), \dot{\gamma}_{(\cdot)}(t))$  is symplectomorphism, that is, for any local coordinates on  $T^*M_0$

$$[\gamma^i, \gamma^j] = 0, \quad [\gamma^i, \dot{\gamma}_j] = \delta_j^i, \quad [\dot{\gamma}_i, \dot{\gamma}_j] = 0, \quad i, j = 1, \dots, n.$$

Further, it is easy to check that the equality

$$L[u, v] = [Lu, v] + [u, Lv] \quad (3)$$

holds.

This motivates the following way to get a copy of hamiltonian  $H$ . Take arbitrary smooth map  $\sigma : Y \rightarrow T^*M_0$  such that

$$\sigma \circ \alpha_H = \sigma, \quad \sigma|_{\partial_0 X} = id,$$

and locally  $\sigma$  is determined by functions  $u^i, v_i$ ,  $i = 1, \dots, n$  on  $Y$  which satisfy symplectic conditions

$$[u^i, u^j] = 0, \quad [u^i, v_j] = \delta_j^i, \quad [v_i, v_j] = 0$$

and initial data

$$u(x, \xi, 0) = x, \quad v(x, \xi, 0) = \xi, \quad (x, \xi) \in \partial_0 X.$$

Then functions  $\tilde{u}^i, \tilde{v}^i$  (extensions) due to 3 also satisfy (for any  $t$ )

$$[\tilde{u}^i, \tilde{u}^j] = 0, \quad [\tilde{u}^i, \tilde{v}_j] = \delta_j^i, \quad [\tilde{v}_i, \tilde{v}_j] = 0.$$

So, by construction, the map  $\tilde{\psi} : X \rightarrow T^*M_0$ ,  $\tilde{\psi} = (\tilde{u}, \tilde{v})$  satisfies kinetic equation  $L\tilde{\psi} = 0$ . At  $t = 0$  according to our notation  $\tilde{\psi}|_{t=0} = (\tilde{u}_0, \tilde{v}_0)$ . The map

$$f : T^*M_0 \rightarrow T^*M_0, \\ (x, \xi) \xrightarrow{f} (\tilde{u}_0(x, \xi), \tilde{v}_0(x, \xi))$$

is symplectomorphism by construction. Then for  $\psi : X \rightarrow T^*M_0$ , defined by the equality

$$\tilde{\psi}(x, \xi, t) = \psi(f(x, \xi), t)$$

we have

$$\psi_t - [\psi, \tilde{H}] = 0, \quad \psi|_{t=0} = id,$$

where  $\tilde{H} = H \circ f$  is a copy of  $H$ . By construction this copy has the same scattering map,  $\alpha_H = \alpha_{\tilde{H}}$ . One can find this copy by the following way. Due to

Remark after Lemma 1 we have  $\psi(x, \xi, t) = \tilde{\varphi}^t(x, \xi)$ , where  $\tilde{\varphi}^t$  is the Hamiltonian flow on  $(T^*M_0, \tilde{H})$ , and therefore for any  $(x, \xi, t) \in X$   $\psi$  satisfies the Hamiltonian system on the whole extended phase space  $X$  :

$$\begin{aligned}\frac{d}{dt}\psi(x, \xi, t) &= \tilde{\mathcal{H}}(\psi(x, \xi, t)), \\ \psi(x, \xi, 0) &= (x, \xi).\end{aligned}$$

In particular we have on the set  $Y$  equalities

$$\begin{aligned}\frac{d\psi}{dt}(x, \xi, t) &= \tilde{\mathcal{H}}(\psi(x, \xi, t)), \quad (x, \xi, t) \in Y, \\ \psi(x, \xi, 0) &= (x, \xi) \in \partial_0 X.\end{aligned}$$

Since  $\psi|_Y = (u, v)$  is given one can obtain the function  $\tilde{H}(u, v)$  (up to a constant) from these equalities.

Thus, a copy  $\tilde{H}$  may be obtained by the following way.

- 1). Take any even w.r.t.  $\alpha_H$  functions  $u^i, v_i, i = 1, \dots, n$  on  $Y$  which satisfy a) symplectic conditions

$$[u^i, u^j] = 0, \quad [u^i, v_j] = \delta_j^i, \quad [v_i, v_j] = 0,$$

- b) the map  $\sigma : (x, \xi, t) \rightarrow (u(x, \xi, t), v(x, \xi, t))$  is the diffeomorphism from  $Y$  onto  $T^*M_0$ .

- 2) The copy  $\tilde{H}$  of hamiltonian  $H$  is determined by choosing  $(u, v)$  and may be found out from the hamiltonian system

$$\begin{aligned}\frac{du^i}{dt}(x, \xi, t) &= \tilde{H}_{v_i}(u(x, \xi, t), v(x, \xi, t)), \\ \frac{dv^i}{dt}(x, \xi, t) &= -\tilde{H}_{u_i}(u(x, \xi, t), v(x, \xi, t)), \quad (x, \xi, t) \in Y. \\ (u, v)|_{t=0} &= (x, \xi).\end{aligned}$$

Notice, that we nowhere used the special form of hamiltonian  $H(x, \xi) = g^{ij}(x)\xi_i\xi_j/2$ . And by construction all copies of unknown hamiltonian  $H$  generate the same scattering map  $\alpha_H$ .